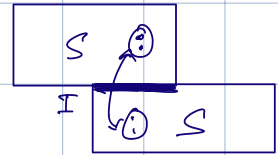


Problem set 4

Problem 1: Charge - phase relation

At $T=0$ and $V=0$, charge can only flow through units of Cooper pairs. We want to describe the behaviour of junction in terms of number of Cooper pairs.



Define \hat{N} which represents number of Cooper pair on side of the junction.

$\hat{N}|n\rangle = n|n\rangle \rightarrow \hat{N}$ tells you number of Cooper pair on one side of junction.

Phase is another quantity that defines a superconductor. We want to relate these

$$|\varphi\rangle = \sum_{n=-\infty}^{\infty} e^{in\varphi} |n\rangle$$

Switching between the basis,

$$|n\rangle = \frac{1}{2\pi} \int_0^{2\pi} d\varphi e^{-in\varphi} |\varphi\rangle$$

Fourier like relation
just like $\hat{x} - \hat{p}$.

$$(a) |n\rangle = \frac{1}{2\pi} \int_0^{2\pi} d\varphi e^{-in\varphi} \sum_{m=-\infty}^{\infty} e^{im\varphi} |m\rangle$$

When evaluating this, use Kronecker delta relation

$$\frac{1}{2\pi} \int_0^{2\pi} dx e^{-i(x-x')k} = \delta_{xx'}$$

$$(b) \langle \varphi | \varphi' \rangle = \sum_{n=-\infty}^{\infty} e^{-in\varphi} \langle n | \sum_{m=-\infty}^{\infty} e^{im\varphi'} |m\rangle$$

For charge states, orthogonality implies $\langle n | m \rangle = \delta_{nm}$

Recognize discrete Dirac delta function

$$\delta(\varphi - \varphi') = \frac{1}{2\pi} \sum_n e^{-in(\varphi - \varphi')}$$

Problem 2: Exponent of a phase operator

$$e^{i\hat{\varphi}} = \frac{1}{2\pi} \int_0^{2\pi} d\varphi' e^{i\varphi'} |\varphi'\rangle \langle \varphi'|$$

(a) Evaluate $e^{i\hat{\varphi}} |\varphi\rangle$

$$\Rightarrow e^{i\hat{\varphi}} |\varphi\rangle = \frac{1}{2\pi} \int_0^{2\pi} d\varphi' e^{i\varphi'} |\varphi'\rangle \langle \varphi' | \varphi \rangle$$

Recognize $\langle \varphi' | \varphi \rangle$ from problem 1 (Dirac delta)

$$(b) e^{i\hat{\phi}}|n\rangle = \frac{1}{2\pi} \int_0^{2\pi} d\varphi' e^{i\varphi'} |\varphi'\rangle \langle \varphi'|n\rangle$$

$$\downarrow \sum_{m=-\infty}^{\infty} e^{-im\varphi'} |m\rangle$$

Use orthogonality δ_{mn} and configure the expression as in eqⁿ (2) of problem 1

$$(c) e^{i\hat{\phi}} = \frac{1}{2\pi} \int_0^{2\pi} d\varphi' e^{i\varphi'} \sum_{m=-\infty}^{\infty} e^{im\varphi'} |m\rangle \sum_{n=-\infty}^{\infty} e^{-in\varphi'} \langle n| = \sum_n |n-\rangle \langle n|$$

1° use the definition of $e^{i\hat{\phi}}$

2° Use definition of $|\varphi'\rangle$ and $\langle \varphi'|$

3° combine the exponential

4° Form Kronecker delta form

5° Apply Kronecker delta to get rid of one sum.

$$e^{-i\hat{\phi}} = \sum_n |n\rangle \langle n-1| \quad \# \text{ Final result!}$$

Problem 3: Josephson Junction

$$H_J = -\frac{E_J}{2} \sum_n |n\rangle \langle n+1| + |n+1\rangle \langle n|$$

$$(a) H_J |n\rangle = -\frac{E_J}{2} \sum_m |m\rangle \langle m+1|n\rangle + |m+1\rangle \langle m|n\rangle$$

Use orthonormality with Dirac delta.

$$(b) H_J = - \frac{E_J}{2} \sum_n |n\rangle \langle n+1| + |n+1\rangle \langle n|$$

Express in terms of $\exp(\pm i\hat{\phi})$. Then, use $\frac{e^{i\phi} + e^{-i\phi}}{2} = \cos \phi$

Problem 4:

$$L = \frac{1}{2} m \dot{\mathbf{x}}^2 + e \dot{\mathbf{x}} \cdot \mathbf{A}(\mathbf{x}, t) - e\phi(\mathbf{x})$$

generalized co-ordinate: $(\bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}) = (x, y, z, \dot{x}, \dot{y}, \dot{z})$

(a) Euler-Lagrange equation [equation of motion]

$$\text{Use } \frac{d}{dt}(\bar{\mathbf{a}} \cdot \bar{\mathbf{b}}) = \frac{d\bar{\mathbf{a}}}{dt} \cdot \bar{\mathbf{b}} + \bar{\mathbf{a}} \cdot \frac{d\bar{\mathbf{b}}}{dt}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$

$$\bar{\mathbf{A}}(\bar{\mathbf{x}}, t) = (A_x(\bar{\mathbf{x}}, t), A_y(\bar{\mathbf{x}}, t), A_z(\bar{\mathbf{x}}, t))$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\mathbf{x}}} \right) - \frac{\partial L}{\partial \mathbf{x}} = 0 \quad \left[L = \frac{1}{2} m (v_x^2 + v_y^2 + v_z^2) + e (v_x A_x + v_y A_y + v_z A_z) - e\phi(\mathbf{x}) \right]$$

1° Consider x -coordinate first.

$$\frac{d}{dt} \left(\frac{\partial L}{\partial v_x} \right) - \frac{\partial L}{\partial x} = 0$$

$$\dot{\bar{\mathbf{x}}} = (v_x, v_y, v_z)$$

↑

$\bar{\mathbf{x}}$ is a 3d vector (x, y, z)

$\dot{\bar{\mathbf{x}}}$ is a derivative of the position x, y, z

First evaluate.

$$\frac{\partial L}{\partial v_x} = \text{# only } v_x \text{ term survives and } \phi(\bar{\mathbf{x}}) \text{ does not depend on } v_x$$

You should get something like this.

$$\frac{d}{dt} \left(\frac{\partial L}{\partial v_x} \right) = \frac{d}{dt}(mv_x) + e \cdot \frac{d}{dt} A_x(\bar{\mathbf{r}}, t) \quad \text{# Use } \frac{d}{dt} \bar{\mathbf{A}}(\bar{\mathbf{r}}, t) = \frac{\partial \bar{\mathbf{A}}}{\partial x} + \frac{\partial \bar{\mathbf{A}}}{\partial y} + \frac{\partial \bar{\mathbf{A}}}{\partial z} + \frac{\partial \bar{\mathbf{A}}}{\partial t}$$

Then evaluate.

Product rule & $\frac{\partial u_i}{\partial i} = 0$. $i \in \{x, y, z\}$

$$\frac{\partial L}{\partial x} = -e \frac{\partial \psi}{\partial x} + e \cdot \frac{\partial}{\partial x} (V_x A_x(r, t) + V_y A_y(r, t) + V_z A_z(r, t))$$

$$\text{Use } \vec{E} = -\left(\frac{\partial \vec{A}}{\partial t} + \nabla \psi\right) \quad \text{and} \quad \vec{B} = \nabla \times \vec{A} \quad \text{to get} \quad \nabla \times \nabla \times \vec{A} = \nabla \times \vec{B}.$$

to recognise the resulting term for x . Similar method can be used to find y , and z component.

$$(b) \quad \vec{p} = \frac{\partial L}{\partial \dot{x}} \quad \# \text{ Evaluate the derivative}$$

$$(c) \quad H = \dot{x} p - L \quad \# \text{ Use result 4(b) in here}$$

express \dot{x} in terms of \vec{p} from 4(b).

Problem 5

$$(a) \quad \text{Use the definition: } |n\rangle = \frac{(\hat{a}^\dagger)^n}{\sqrt{n!}} |0\rangle \quad \text{and use the fact}$$
$$|n+1\rangle = \frac{(\hat{a}^\dagger)^{n+1}}{\sqrt{(n+1)!}} |0\rangle$$

$$(b) \quad \text{Use the complete basis: } \hat{I} = \sum_n |n\rangle \langle n| \quad \text{and the fact that}$$
$$\hat{a}^\dagger = \hat{a}^\dagger \hat{I} \quad \text{and} \quad \hat{a} = \hat{a} \hat{I}. \quad \text{Use the definition of } |n\rangle.$$

(c) Write down the matrix (see lecture notes) and compute the terms.

Problem 6

(a) Use $|z\rangle = \hat{I}_n |z\rangle = \sum_{n=0}^{\infty} c_n |n\rangle$ with

completeness $\hat{I}_n = \sum_{n=0}^{\infty} |n\rangle \langle n|$.

- Use the definition of $|n\rangle = \frac{(\hat{a}^\dagger)^n |0\rangle}{\sqrt{n!}}$. Remember to take the conjugate for $\langle n|$.

- $\hat{a}|z\rangle = z|z\rangle$ (given)

(b) Use $\langle z|z\rangle = 1$

- Plug $\langle z|$, and $|z\rangle$ from (a). Remember to conjugate $\langle z|$.

- Recognize Taylor expansion

$$e^x = \sum_n \frac{x^n}{n!}$$

(c) Find $P(n) = |\langle n|z\rangle|^2$

- Use $|z\rangle$ from (a) with proper normalization constant found in (b).

(d) $1^0 \langle z|\hat{n}|z\rangle = \langle z|\hat{a}^\dagger \hat{a}|z\rangle$ Use $\hat{n} = \hat{a}^\dagger \hat{a}$.

- $\hat{a}|z\rangle = z|z\rangle$ given

- Remember $(\hat{a}|z\rangle)^\dagger = \langle z|\hat{a}^\dagger$.

$$2^0 \langle z|\hat{n}^2|z\rangle = \langle z|\hat{n}\hat{n}|z\rangle = \langle z|\hat{a}^\dagger \hat{a} \hat{a}^\dagger \hat{a}|z\rangle$$

- Use $\hat{a}\hat{a}^\dagger = 1 + \hat{a}^\dagger \hat{a}$. Evaluate $\langle \hat{n} \rangle^2$.