



Aalto University
School of Electrical
Engineering

Celestial Mechanics and Satellite Orbits

Introduction to Space

Slides: Jaan Praks, Hannu Koskinen, Zainab Saleem

Lecture: Jaan Praks

Assignment

- **Draw Earth, and a satellite orbiting the Earth.**
- **Draw the orbit of the satellite.**
- **Mark rotation direction of the Earth.**
- **Change your picture with your neighbour. Discuss.**



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History

Schema huius præmissæ diuisionis Sphærarum.



Schema huius præmiſſæ diuifionis Sphærarum.





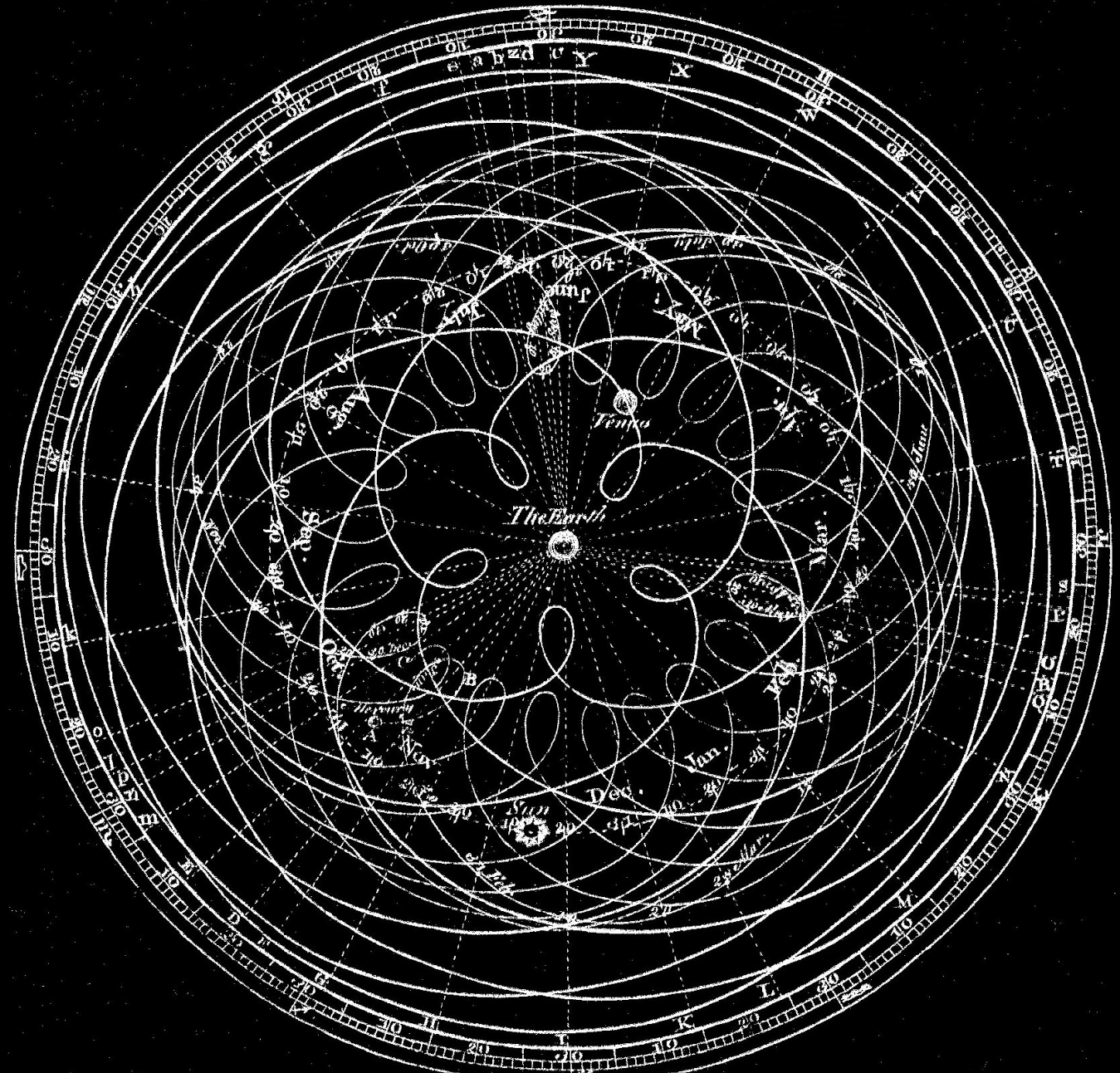
April 28
2008

August 23
2007

(Photograph ©2007–08 Tunç Tezel.)

Planets, stars with will of their own...





405

NICOLAI COPERNICI TORINENSIS
DE REVOLUTIONIBVS ORBIVM
CELESTIVM, Libri VI.

Habes in hoc opere iam recens nato, & æditi, studiose lector, Motus stellarum, tam fixarum, quam erraticarum, cum ex ueteribus, tum etiam ex recentibus obseruationibus restitutos: & nouis insuper ac admirabilibus hypothesibus ornatos. Habes etiam Tabulas expeditissimas, ex quibus eisdem ad quoduis tempus quam facillime calculare poteris. Igitur eme, lege, frue.

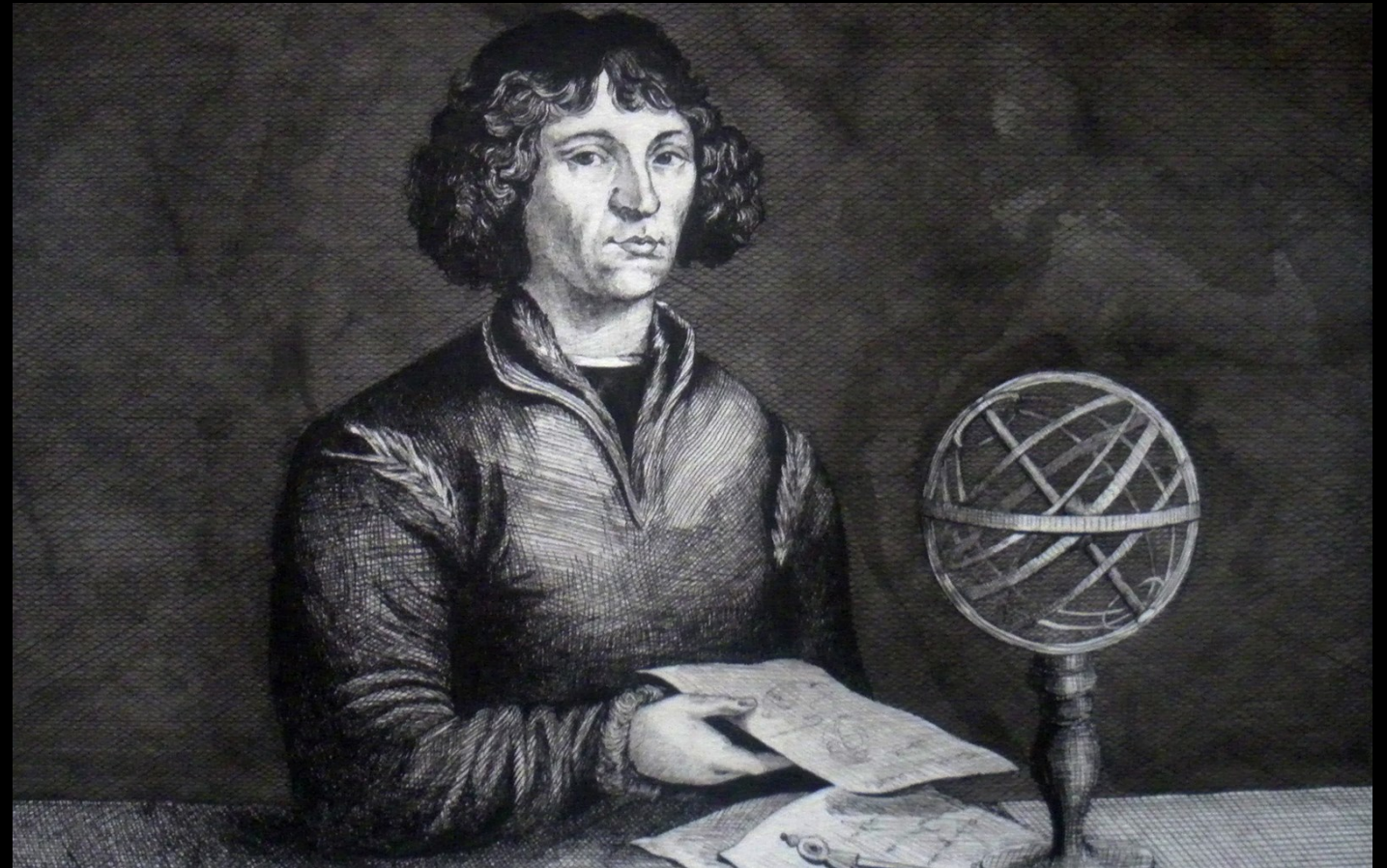


Αγορεύεται εις τας εδρας.

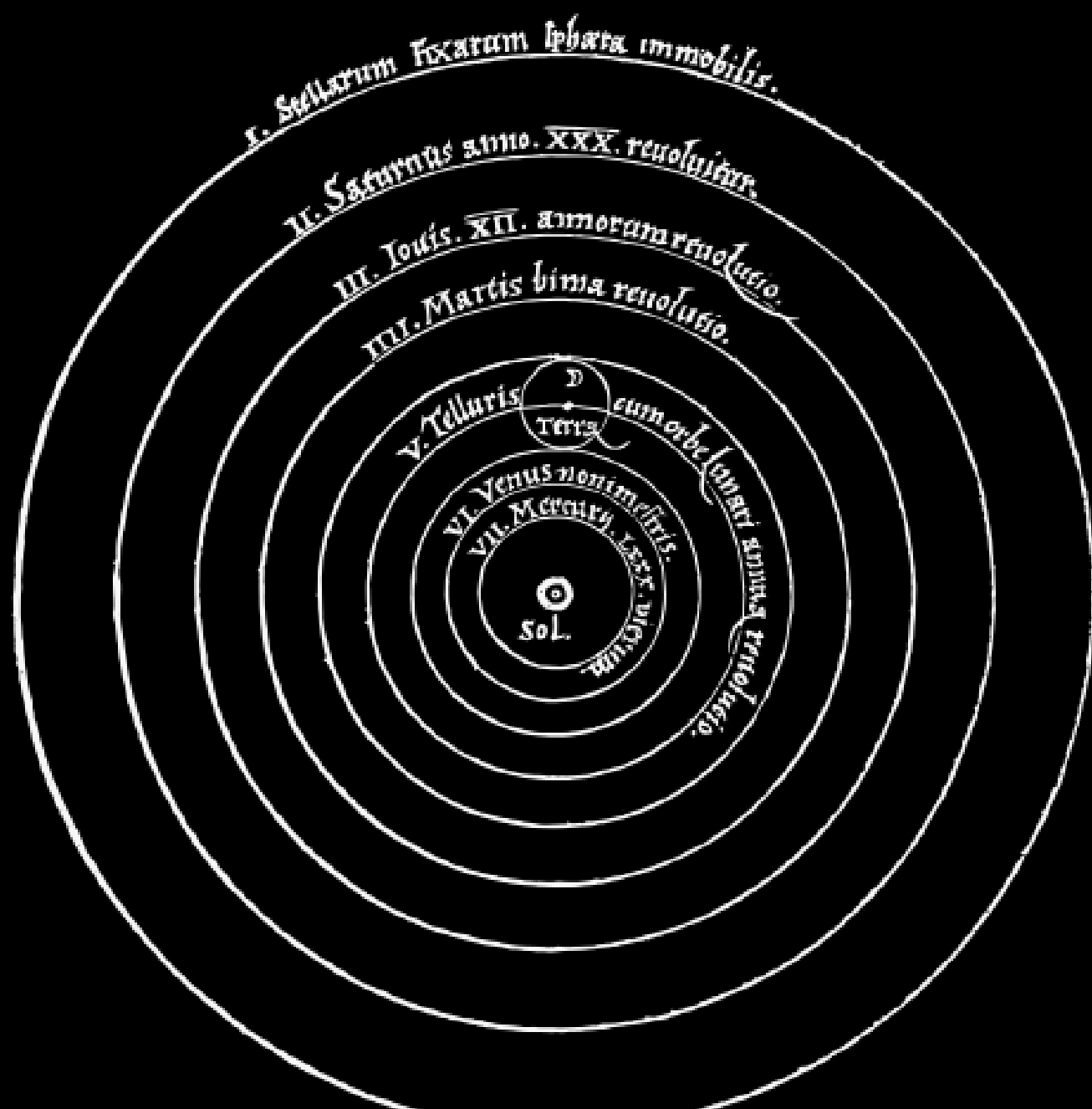
Simon Stevinus
Mathysen Com. Dreyerum.

Norimbergæ apud Ioh. Petreium,
Anno M. D. XLIII.

1420-2421



Nicolaus Copernicus 1473 - 1543



Tycho Brahe (1546 – 1601)

Danish nobleman and astronomer.

Passionate about planetary motion.

Made the most accurate measurements of planetary movements. (without a telescope!)

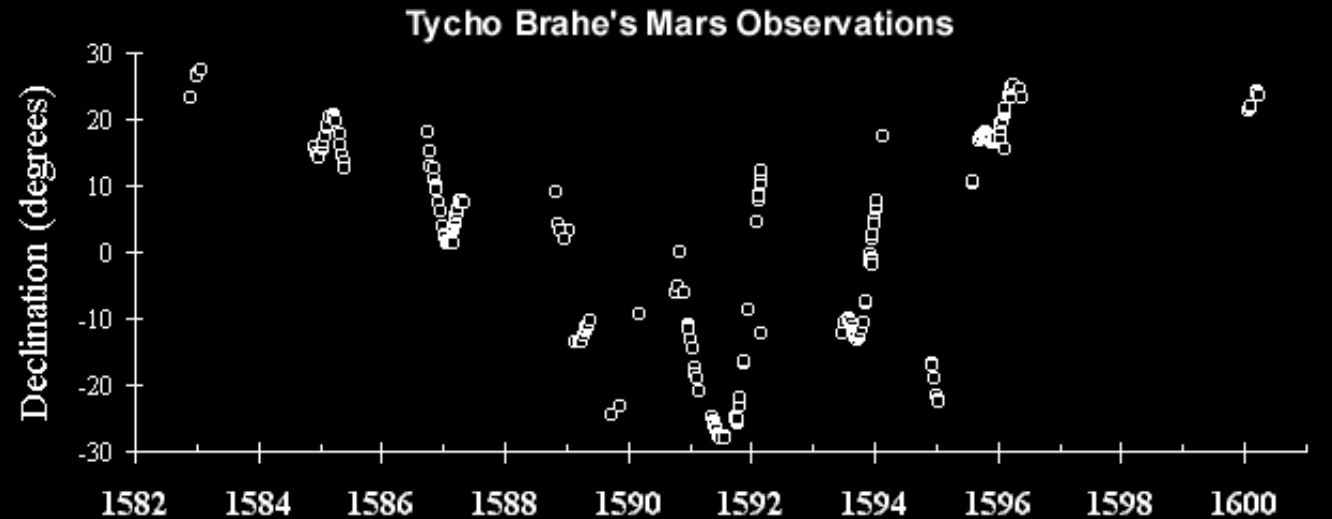


Image Copyright 2000, Wayne Paiko



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Kepler's laws

Based on observations of Tycho Brahe 1546 – 1601

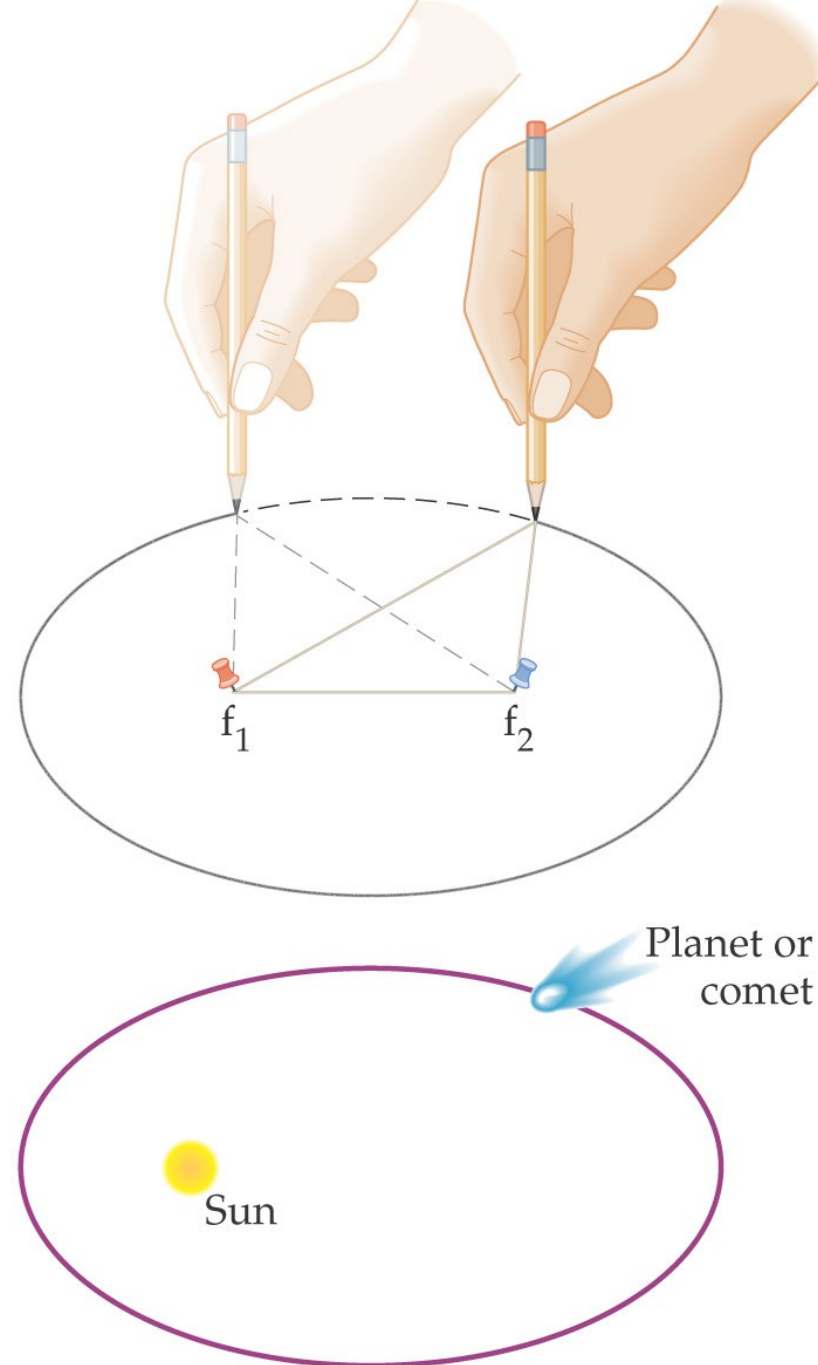
- I. Each of the planets moves on an elliptical path with the Sun at one focus of the ellipse (1609)
- II. For each of the planets, the straight line connecting the planet to the Sun sweeps out equal areas in equal times (1609)
- III. The squares of the periods of the planets are proportional to the cubes of the major axes of their orbits (1619)



**Johannes Kepler
(1571 - 1630)**

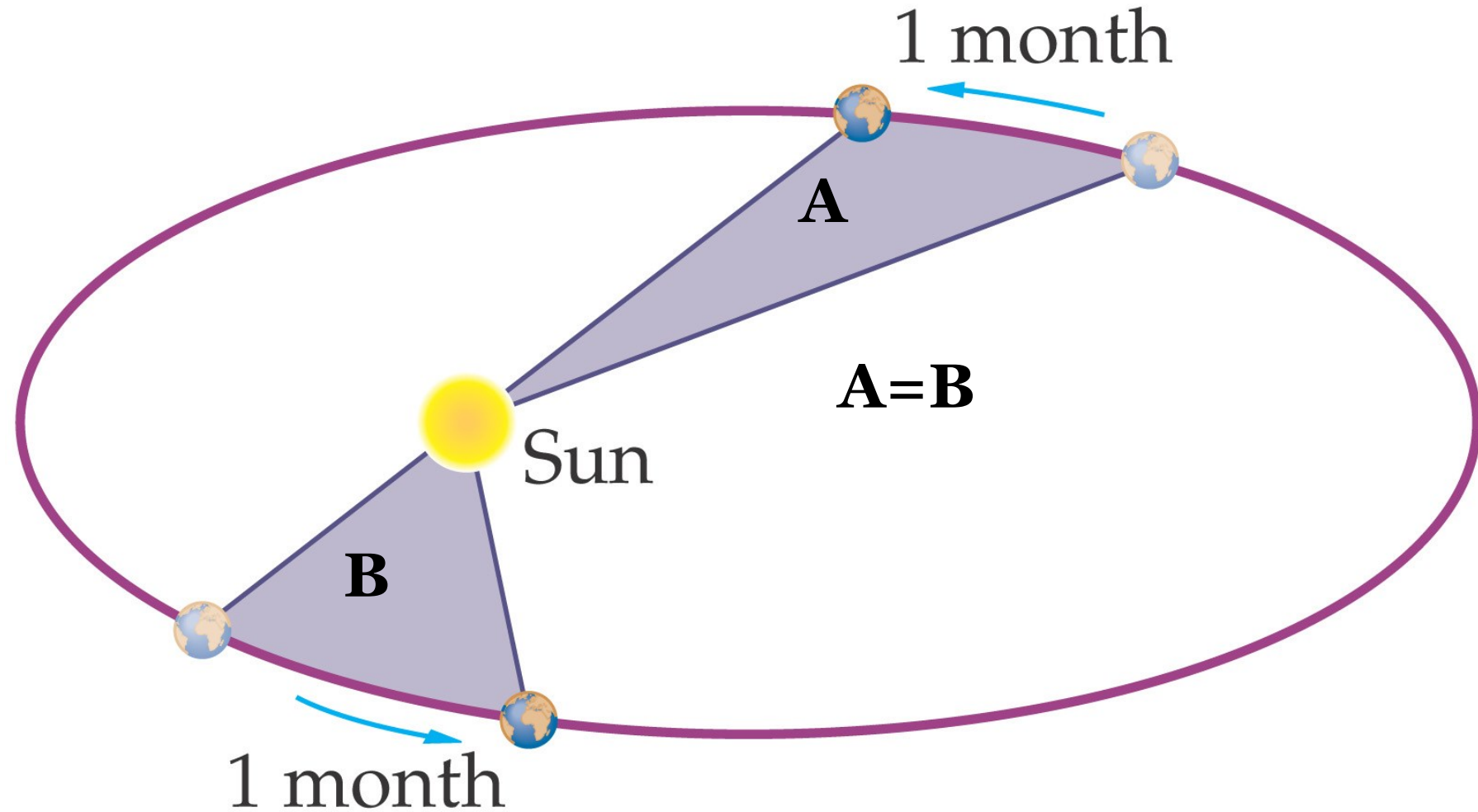
Kepler's

1



Kepler's

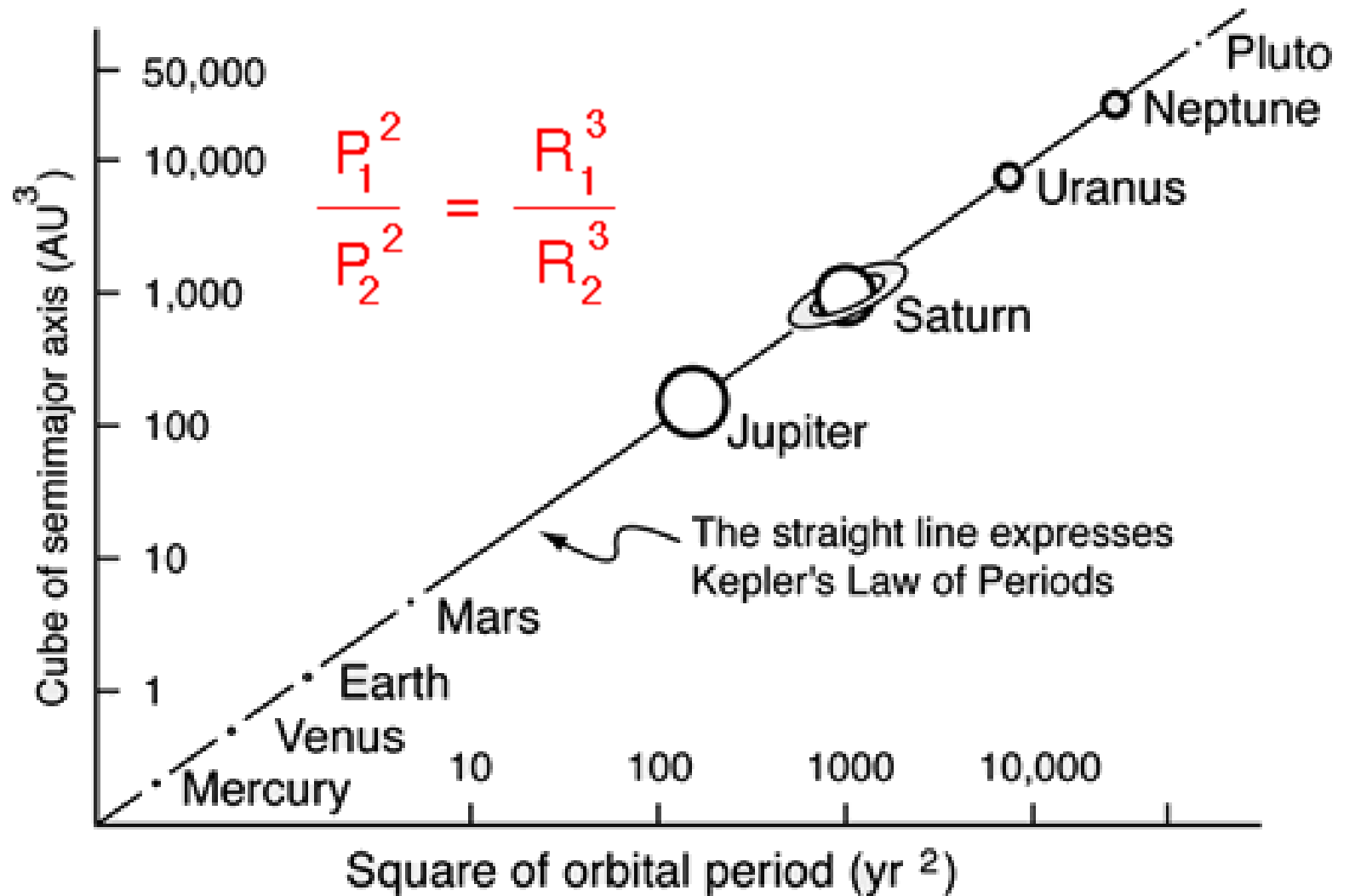
2

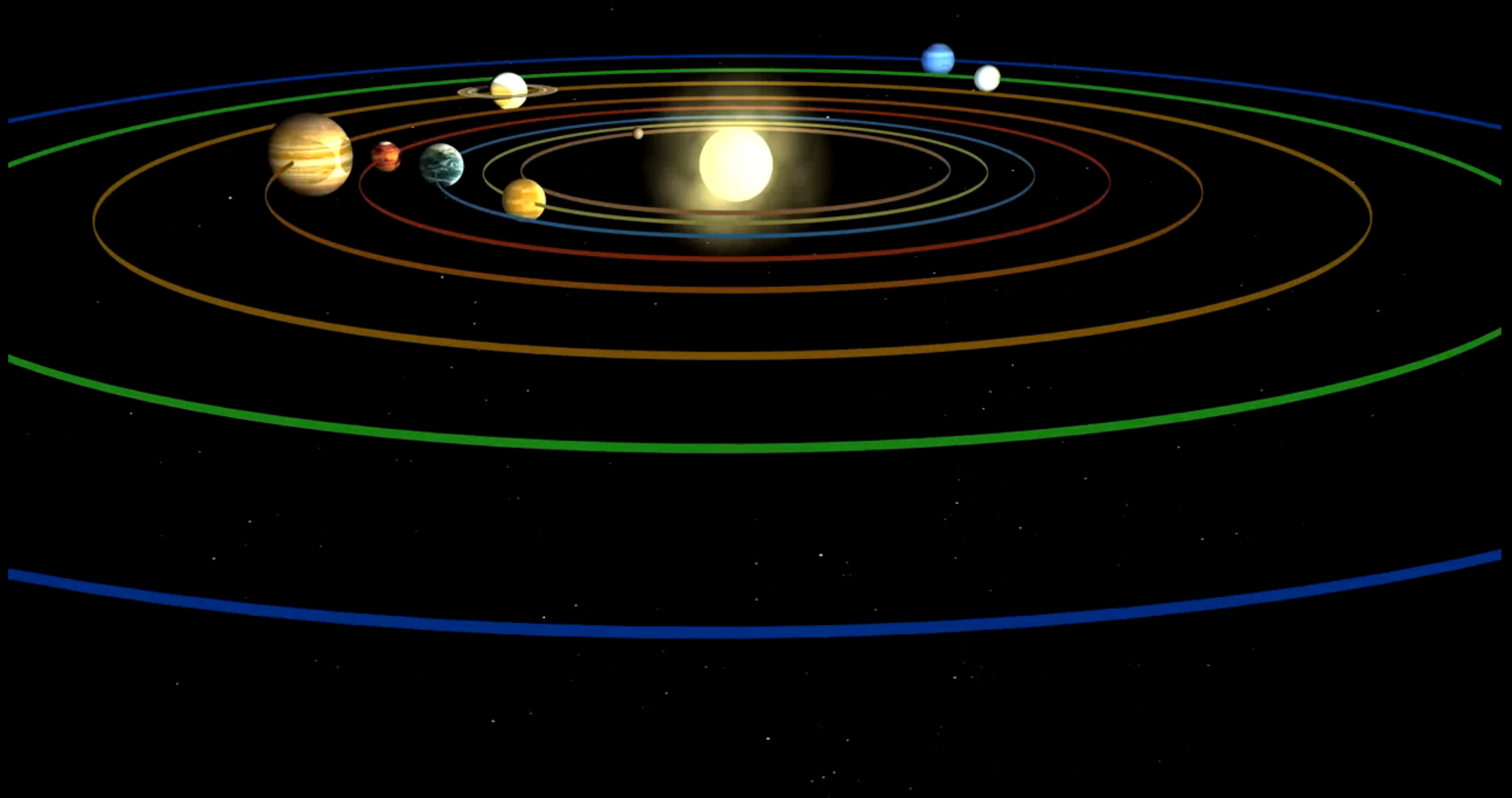


Kepler's

3

Probable examination question!





The Solar Svstem

Newtonian Mechanics

Study of orbits of natural and artificial bodies in space

Based on Newton's laws (*transl. Andrew Motte, 1729*):

“**LAW I.** *Every body perseveres in its state of rest, or of uniform motion in a right line, unless it is compelled to change that state by forces impressed thereon.*”

“**LAW II.** *The alteration of motion is ever proportional to the motive force impressed; and is made in the direction of the right line in which that force is impressed.*”

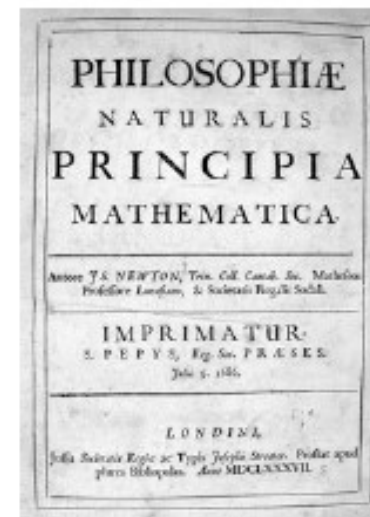
$$\frac{dp}{dt} = F \quad (p \equiv mv)$$

“**LAW III.** *To every action there is always opposed an equal reaction; or the mutual actions of two bodies upon each other are always equal, and directed to contrary parts.*”

$$F_{12} = -F_{21}$$



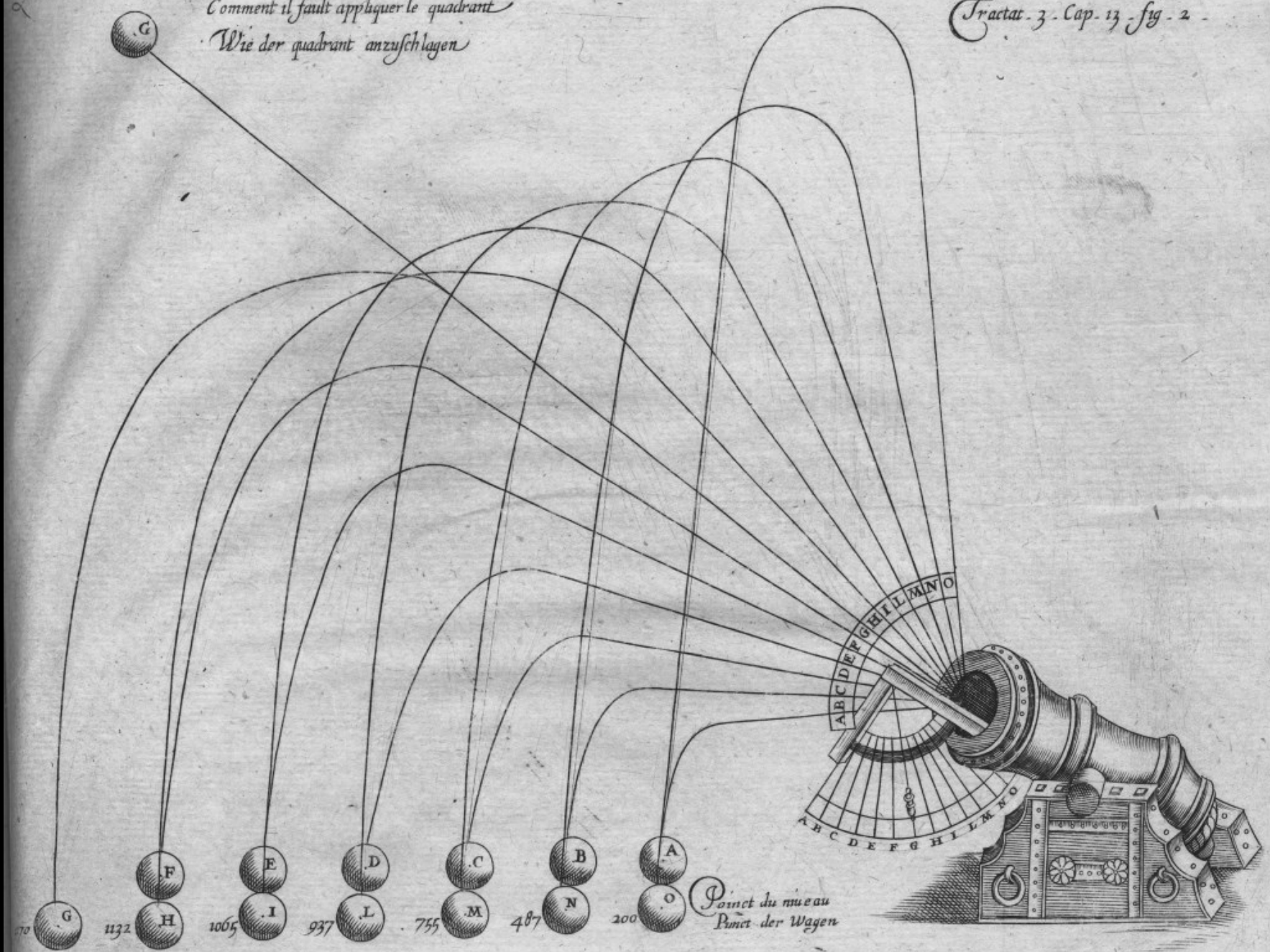
Isaac Newton
(1642 – 1727)



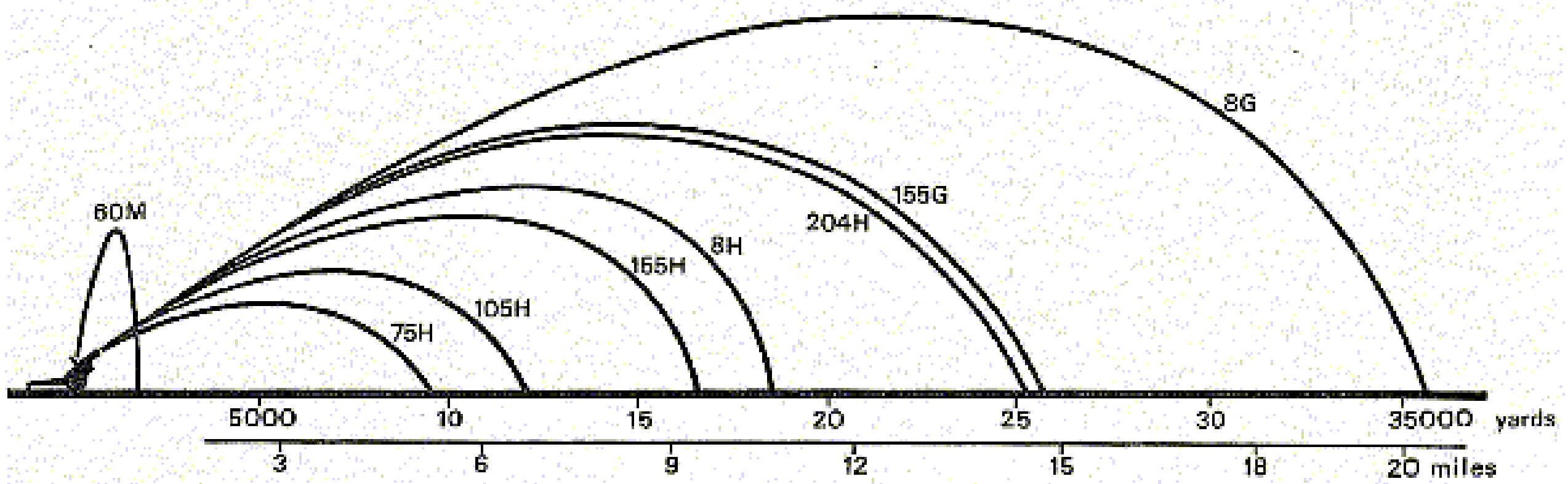
Principia, 1687

Comment il fault appliquer le quadrant
Wie der quadrant anzuschlagen

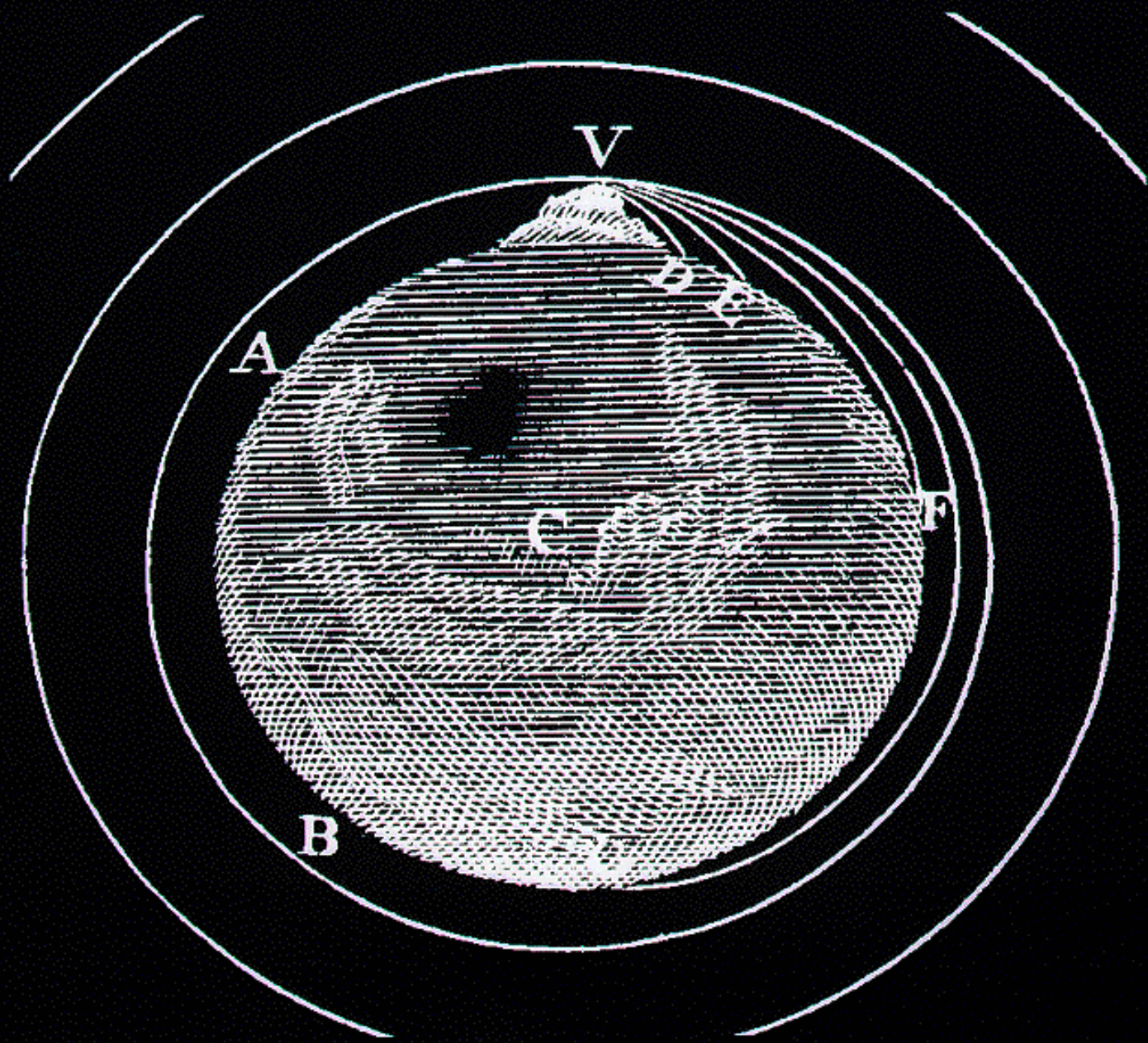
Tractat. 3. Cap. 13. fig. 2.



M = Mortar
H = Howitzer
G = Gun



after dinner, the weather being warm, we went
 into the garden, & drank thea with
 some apples; only he, & my
 other discourse, he told me, the
 same similitude, as when former
 gravitation came into his mind
 apples always descend perpendic
 ground, thought he to himself; or
 of an apple, as he sat in a chair
 why sh^d it not go sideways, or up
 slantly to the earths center? as
 Sen is, that the earth draws it.
 drawing power in matter. & the
 ing power in the matter of the o.
 the earths center, not in any sid
 therefore does this apple fall
 or toward the center. if matter
 ter; it must be in proportion of i
 therefore the apple draws the
 as the earth draws the apple.





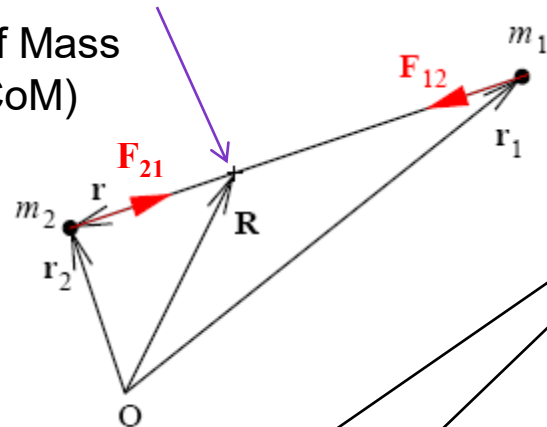
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Basics of classical orbital mechanics

Two body problem: Reduction to one-body problem; central forces

$$\mathbf{R} = \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2}$$

Center of Mass
CM (or CoM)



$$\mathbf{r}_1 = \mathbf{R} - \frac{m_2 \mathbf{r}}{m_1 + m_2}$$

$$\mathbf{r}_2 = \mathbf{R} + \frac{m_1 \mathbf{r}}{m_1 + m_2}$$

$$m_1 \ddot{\mathbf{r}}_1 = \mathbf{F}_{12};$$

$$m_2 \ddot{\mathbf{r}}_2 = \mathbf{F}_{21} = -\mathbf{F}_{12} \Rightarrow$$

$$m_2 \ddot{\mathbf{r}}_2 - m_1 \ddot{\mathbf{r}}_1 = 2\mathbf{F}_{21}$$

$$m \ddot{\mathbf{r}} = \mathbf{F}_{21} \quad m = \frac{m_1 m_2}{m_1 + m_2} \quad \text{reduced mass}$$

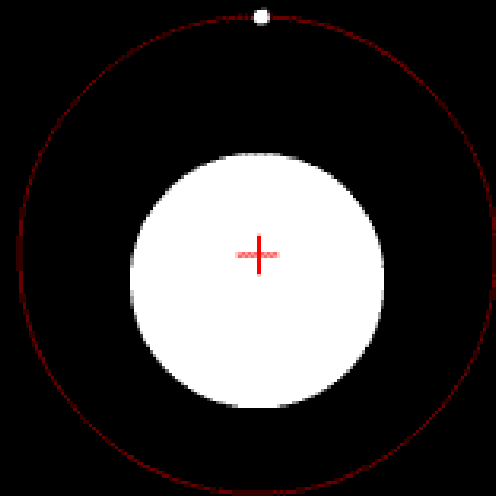
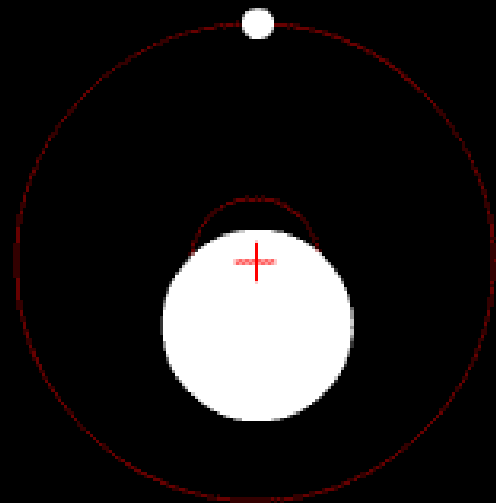
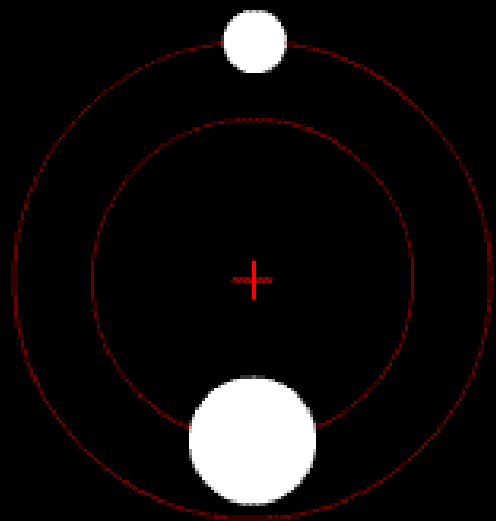
$$\mathbf{F}_{21} = f(r) \mathbf{e}_r \quad \text{central force}$$

Example

$$\mathbf{F} = -\frac{GMm}{r^2} \mathbf{e}_r \quad \text{Newton's law of gravity}$$

$$m_1 \gg m_2 \quad m \approx m_2 \quad r_1 \approx \mathbf{R} \quad r_2 \approx \mathbf{R} + \mathbf{r}$$





Motion under central force (e.g. gravity)

$$\mathbf{F}(\mathbf{r}) = f(r)\mathbf{e}_r \quad \text{points to the origin}$$

Calculate the moment of the force

$$\mathbf{r} \times \dot{\mathbf{p}} = \mathbf{r} \times \mathbf{F} = f(r)\mathbf{r} \times \mathbf{e}_r = 0$$

Calculate the time derivative of the angular momentum

$$\dot{\mathbf{L}} = \frac{d}{dt}(\mathbf{r} \times \mathbf{p}) = \mathbf{v} \times \mathbf{p} + \mathbf{r} \times \dot{\mathbf{p}} = 0$$

\mathbf{L} is a constant of motion

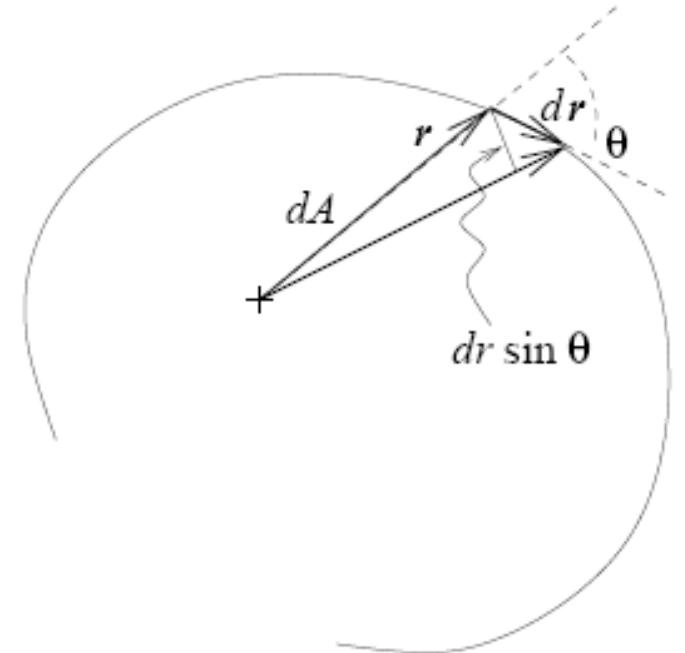
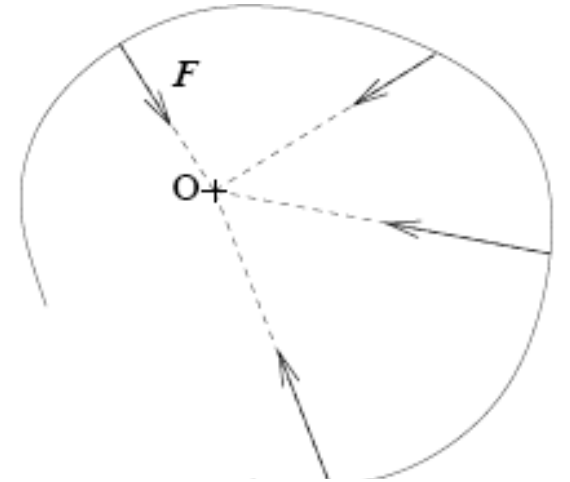
$\mathbf{r} \perp \mathbf{L}$ and $\mathbf{v} \perp \mathbf{L}$: the motion takes place on a plane

Calculate the area element

$$dA = \frac{1}{2} r \sin \theta dr = \frac{1}{2} |\mathbf{r} \times d\mathbf{r}|$$

$$\Rightarrow \frac{dA}{dt} = \frac{1}{2} |\mathbf{r} \times \mathbf{v}| = \frac{1}{2m} |\mathbf{L}|$$

The surface velocity dA/dt is constant (Kepler's second law)



Solutions of the equation of motion

Conic sections

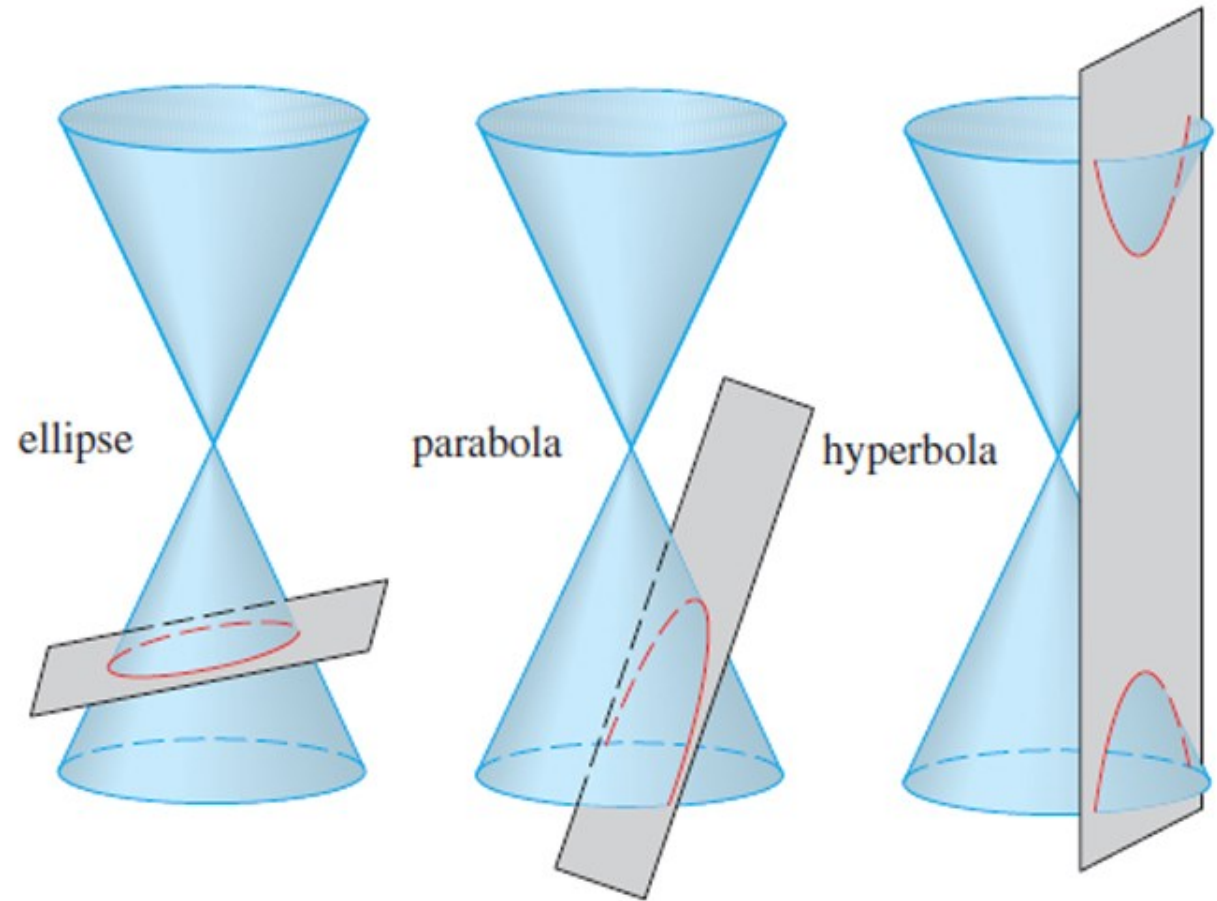
Solutions of the equation $m\ddot{\mathbf{r}} = -\frac{GMm}{r^2}\mathbf{e}_r$ are of the form

$$r = \frac{p}{1 + \varepsilon \cos \theta}; \quad p = \frac{l^2}{mk} \quad |\mathbf{L}| = l$$

$$\varepsilon = \sqrt{1 + \frac{2El^2}{mk^2}} \quad k = GMm$$

These are known as conic sections with ellipticity:

$$\begin{cases} \varepsilon = 0 & \text{circle} \\ 0 < \varepsilon < 1 & \text{ellipse} \\ \varepsilon = 1 & \text{parabola} \\ \varepsilon > 1 & \text{hyperbola} \end{cases}$$



Trajectories

If $e > 1$, the trajectory is a hyperbola

(an open trajectory)

If $e = 1$, the trajectory is a parabola

(an open trajectory)

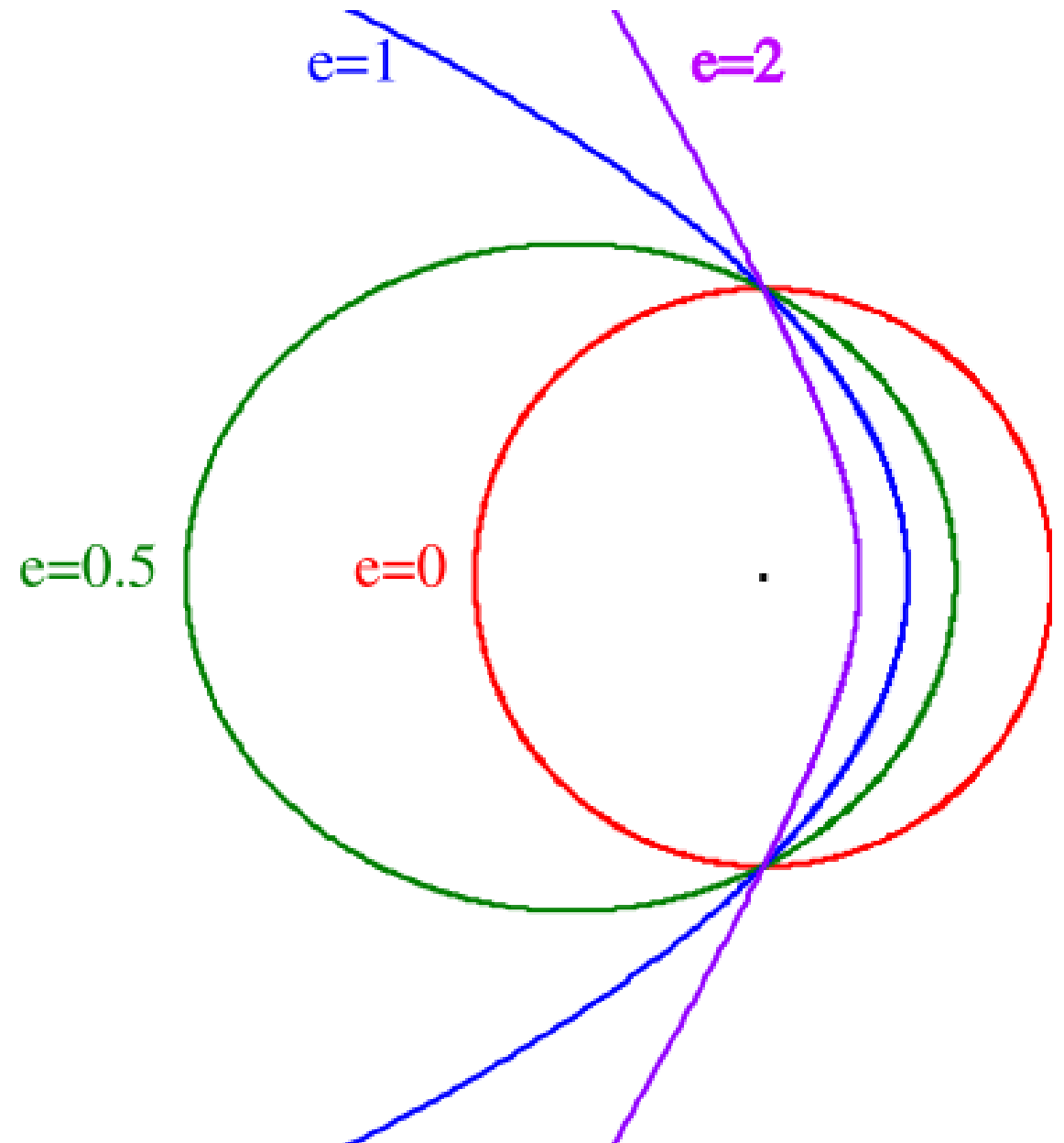
If $0 < e < 1$, the trajectory is an ellipse

(a closed trajectory; i.e., an orbit)

If $e = 0$, the orbit is circular.

(a closed trajectory; i.e., an orbit)

This is just a special case of the ellipse



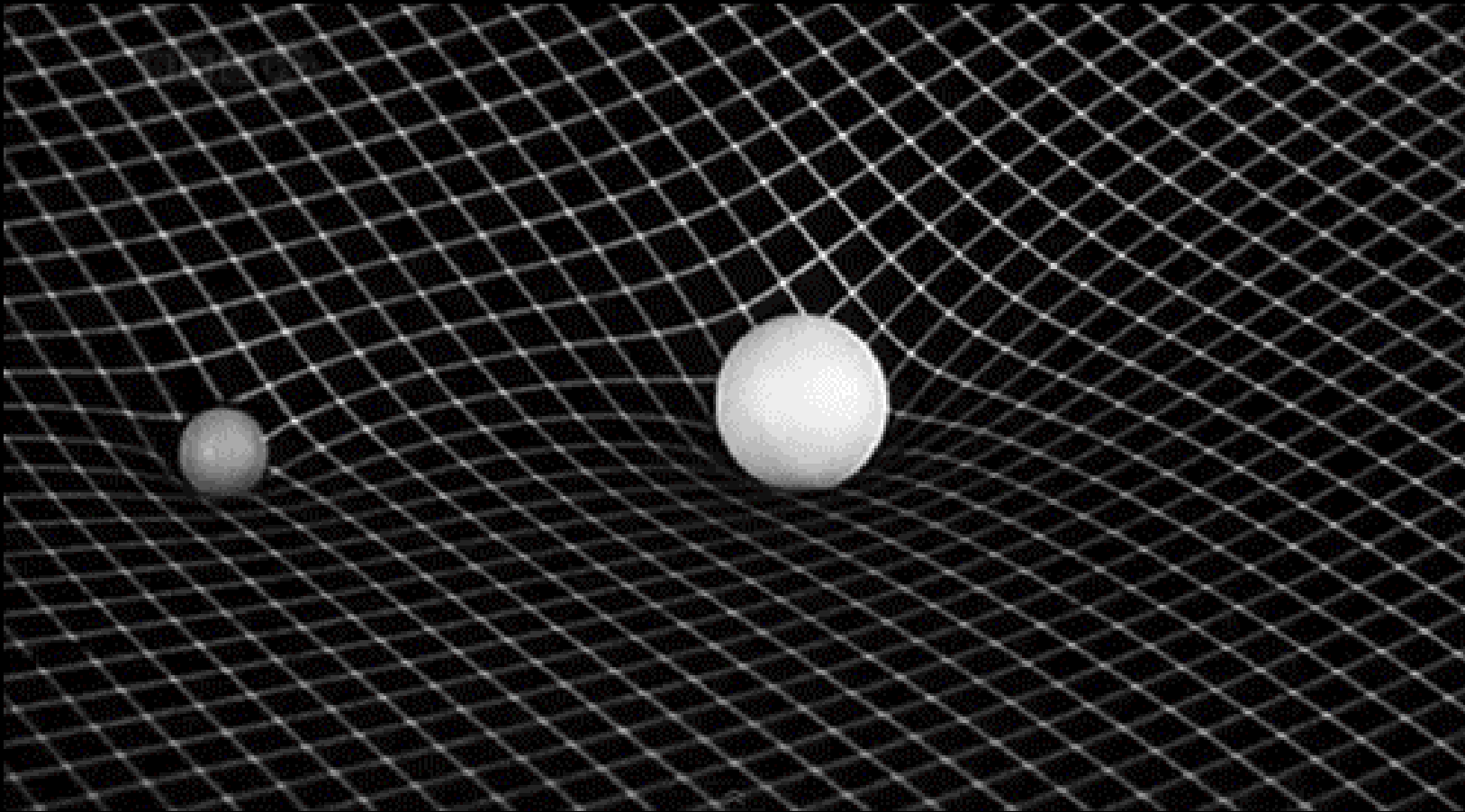
A?

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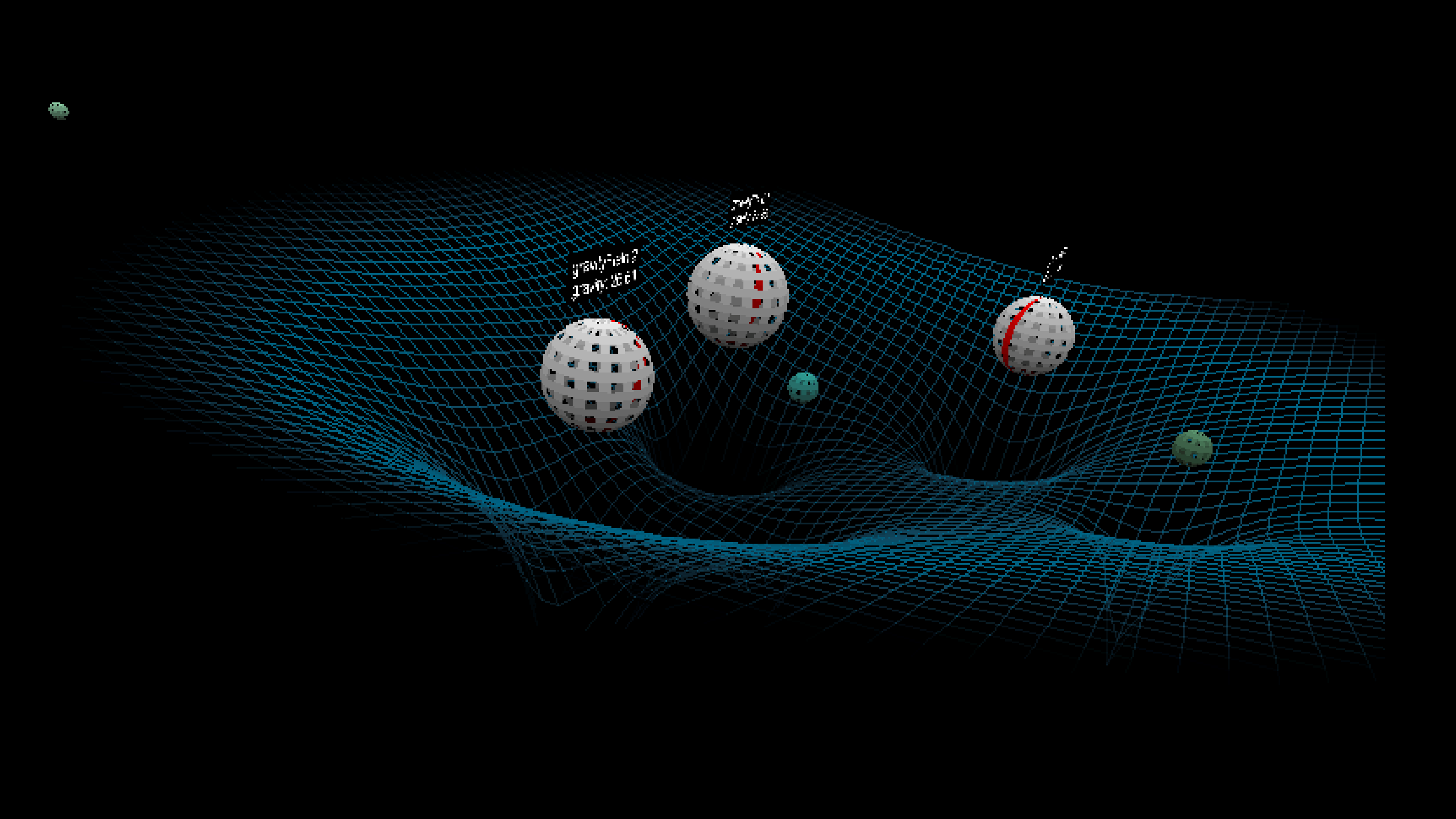
Modern interpretation of gravity

General Relativity (1915)

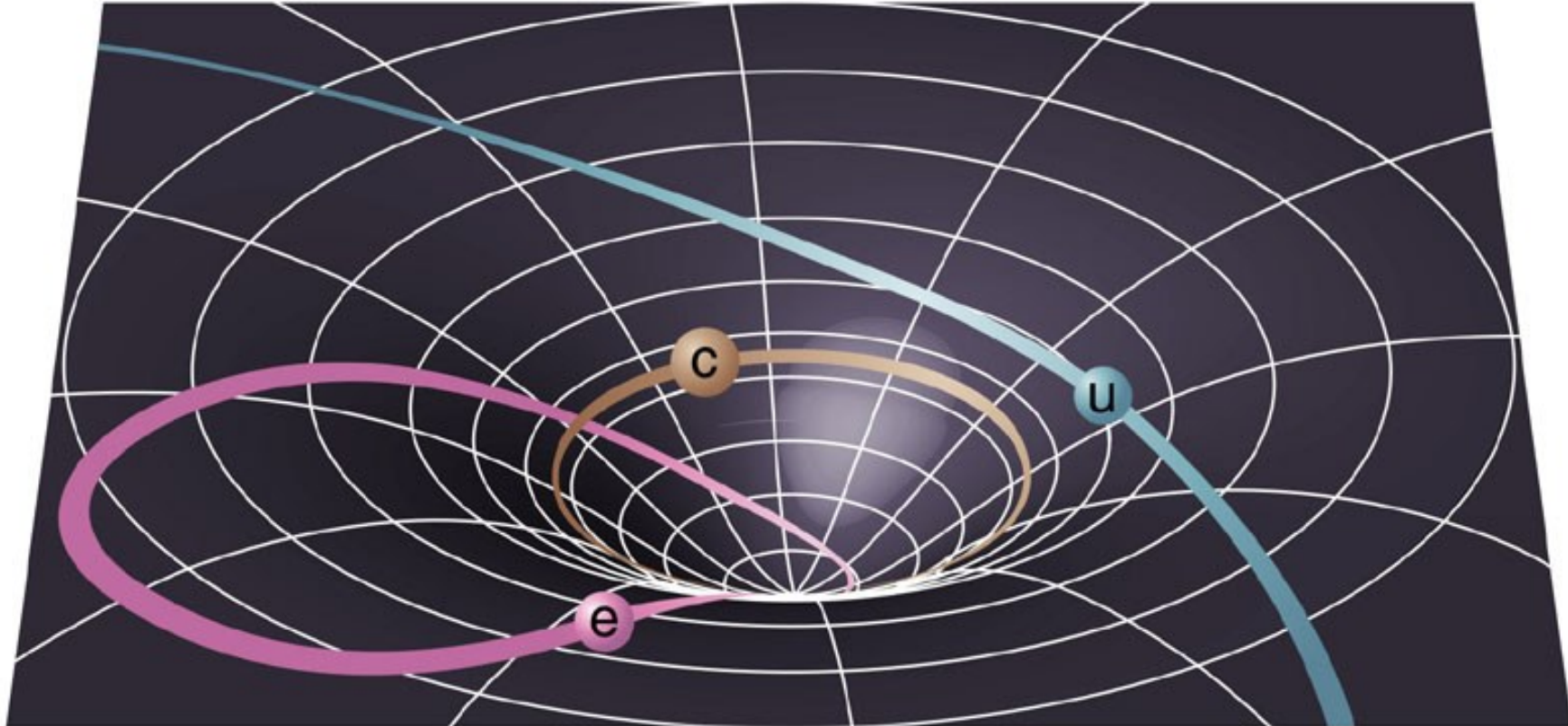
The observed gravitational effect between masses results from their warping of spacetime



Albert Einstein



- c circular orbit
- e elliptical orbit
- u unbound orbit



Mike Gruntman
file: mikegruntman-02.wmv
run time 5 min 30 sec

Educational Use Only

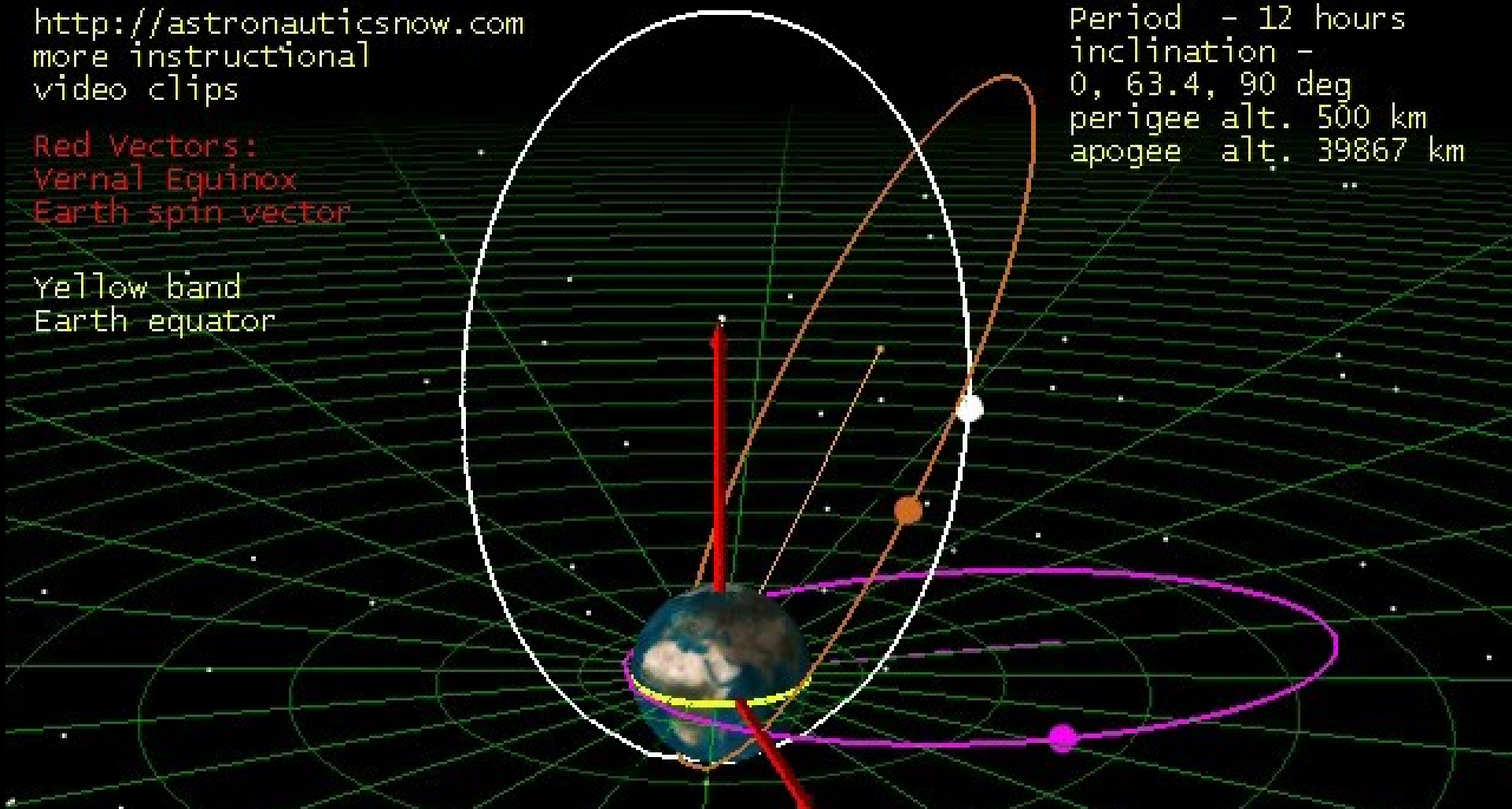
Rotation of Apsides
effect of J2

<http://astronauticsnow.com>
more instructional
video clips

Red Vectors:
Vernal Equinox
Earth spin vector

Yellow band
Earth equator

Orbits:
Period - 12 hours
inclination -
0, 63.4, 90 deg
perigee alt. 500 km
apogee alt. 39867 km



Grid 100000.0 km (10000.0 km)

Earth Inertial Axes

3 Jan 2008 13:20:00.000

Time Step: 600.00 sec

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Orbits and their properties

Elliptical orbits and Kepler's laws

K-I $r = \frac{p}{1 + \varepsilon \cos \theta}; \quad \varepsilon = \sqrt{1 + \frac{2El^2}{mk^2}} < 1;$

$$2a = \frac{p}{1 - \varepsilon} + \frac{p}{1 + \varepsilon} = \frac{2p}{1 - \varepsilon^2}$$

$$c = a - p/(1 + \varepsilon) = \varepsilon a$$

$$b^2 = a^2 - c^2 = a^2(1 - \varepsilon^2) = ap$$

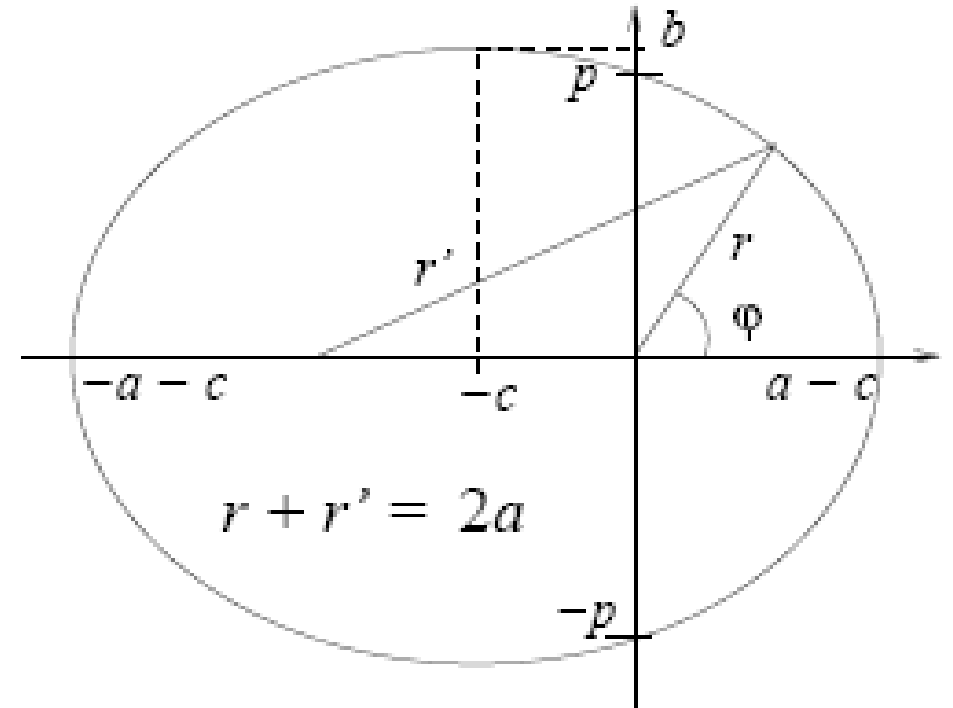
K-II During an infinitesimal time period dt

$$dA = \frac{1}{2}r r d\varphi \Rightarrow \dot{A} = \frac{1}{2}r^2 \dot{\varphi} = \frac{l}{2m} = \text{constant}$$

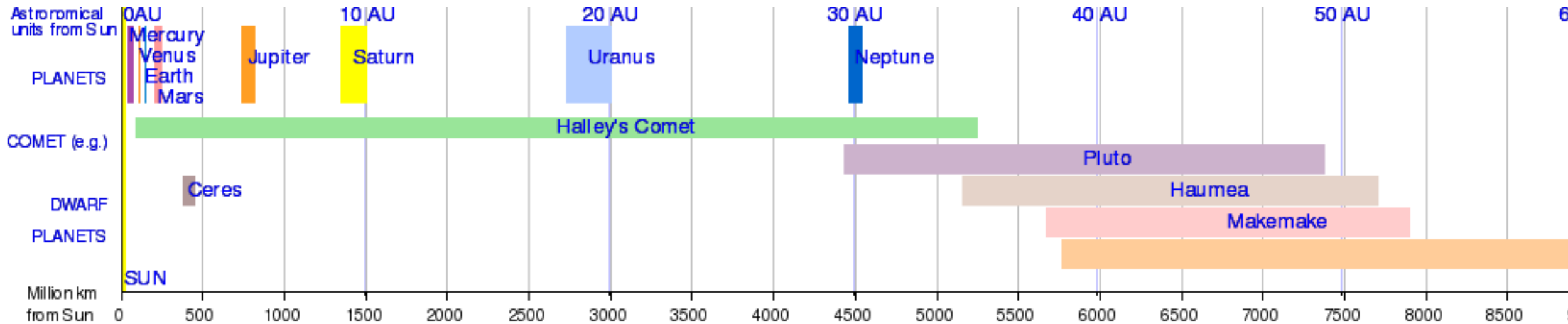
K-III From above $\frac{b^2}{a} = p = \frac{l^2}{mk} \Rightarrow b = \frac{\sqrt{al}}{\sqrt{mk}}$ and the area of the ellipse is

$$A = \pi ab$$

$$A = \int_0^T \dot{A} dt = \frac{lT}{2m} = \pi a \frac{\sqrt{al}}{\sqrt{mk}} \Rightarrow T = 2\pi \sqrt{\frac{m}{k}} a^{3/2}$$



Ellipticity of planetary orbits





Kinetic + potential energy

Orbital energy is constant

Kinetic plus potential energy

$$E = \frac{-GMm}{r} + \frac{M(v_M)^2}{2} + \frac{m(v_m)^2}{2}$$

kinetic

$$E = \frac{-GMm}{r} + \mu \frac{v^2}{2}$$

Potential

Parabolic and hyperbolic orbits: Escape velocity

For an open conic section

thus $E \geq 0$.

$$\varepsilon = \sqrt{1 + \frac{2El^2}{\mu k^2}} \geq 1 \quad \mu = \frac{mM}{m+M}$$

and the minimum velocity to escape is given by $\frac{1}{2}\mu v_e^2 = \frac{GMm}{r}$

$$v_e = \sqrt{\frac{2GMm}{\mu r}} = \sqrt{\frac{2G(M+m)}{r}} \approx \sqrt{\frac{2GM}{r}}$$

From the surface of the Earth $v_e = 11.2$ km/s

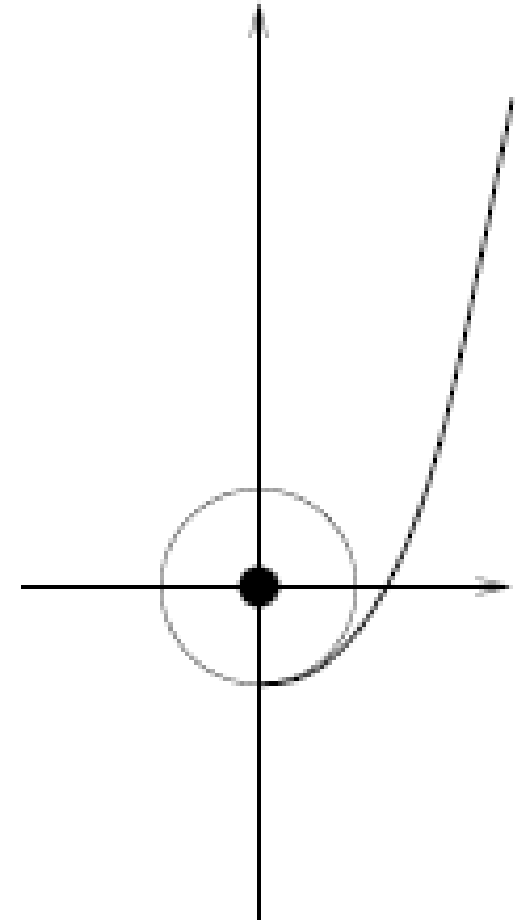
From the surface of the Sun $v_e = 618$ km/s

Escape from geostationary orbit $r = 6.6 R_E$

$$v_{\varphi 0} = 2\pi r / (24 \text{ h}) = 3.06 \text{ km/s}$$

$$v_e = 11.2 \text{ km/s} / \sqrt{6.6} = 4.36 \text{ km/s}$$

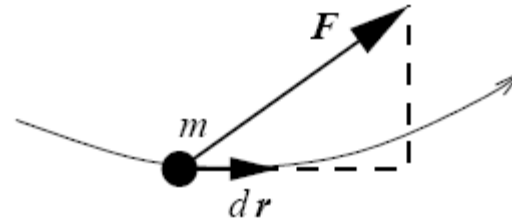
$\Delta v = v_e - v_{\varphi 0} = 1.3$ km/s in the direction of the motion



How to change the orbit

Recall the work done by a force

$$W = \int_{r(t)}^{r(t+\Delta t)} \mathbf{F} \cdot d\mathbf{r}$$
$$= \int_t^{t+\Delta t} \mathbf{F} \cdot \frac{d\mathbf{r}}{dt} dt$$

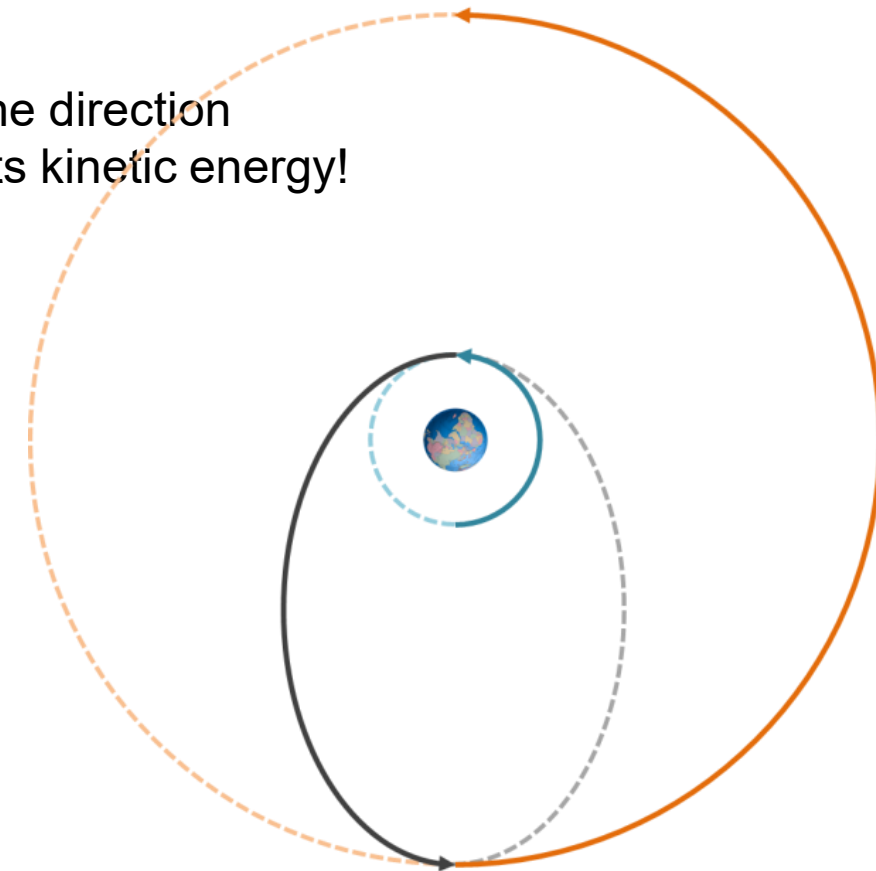


Thus only the component of force in the direction of the spacecraft motion can change its kinetic energy!

To lift the apogee, fire the rocket in perigee
To lift the perigee, fire the rocket in apogee

To reach from one circular orbit to another with least energy:

Hohmann transfer orbit



Vis-viva equation

Orbital energy is constant

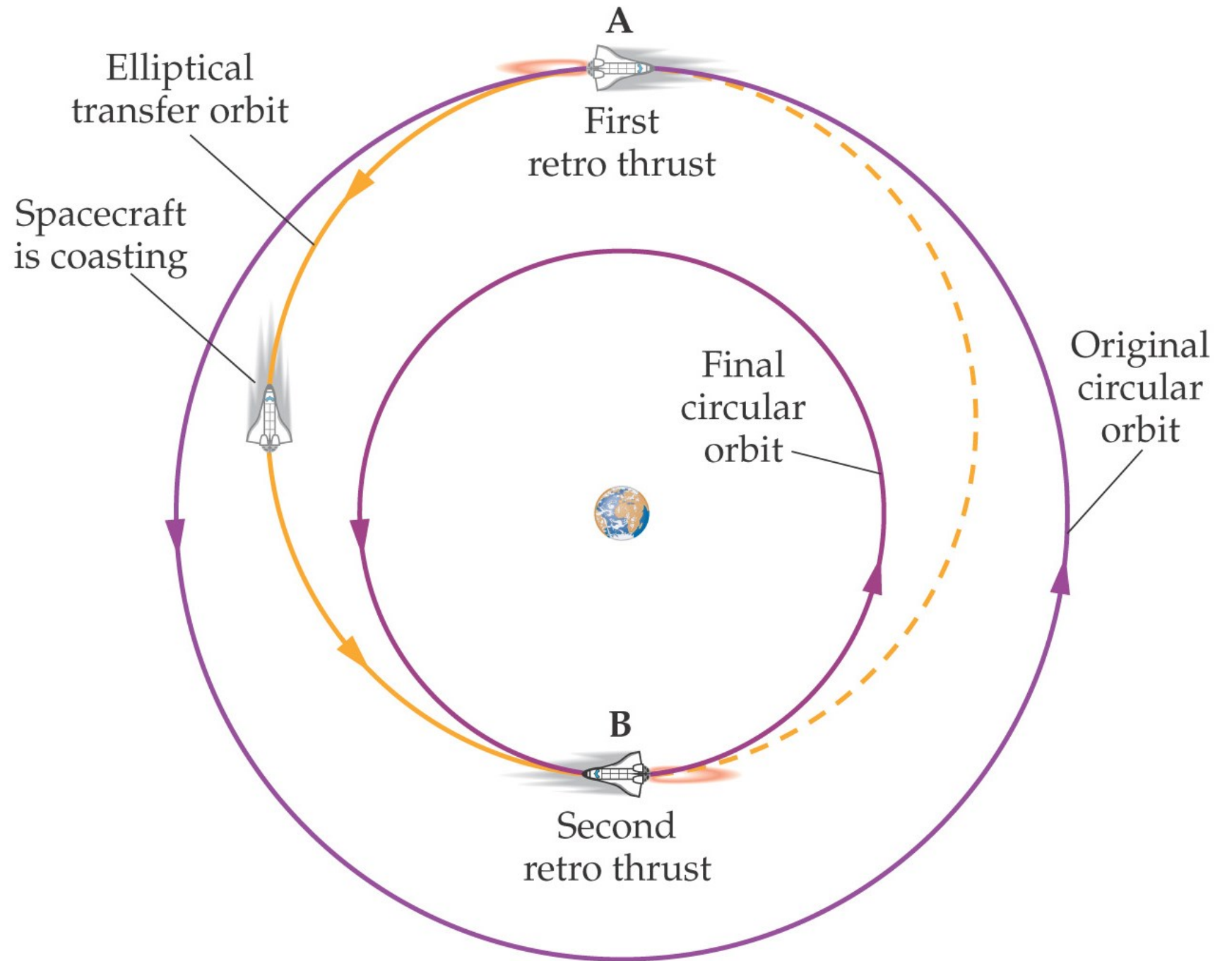
Kinetic plus potential energy

$$E = \frac{-GMm}{r} + \frac{M(v_M)^2}{2} + \frac{m(v_m)^2}{2}$$

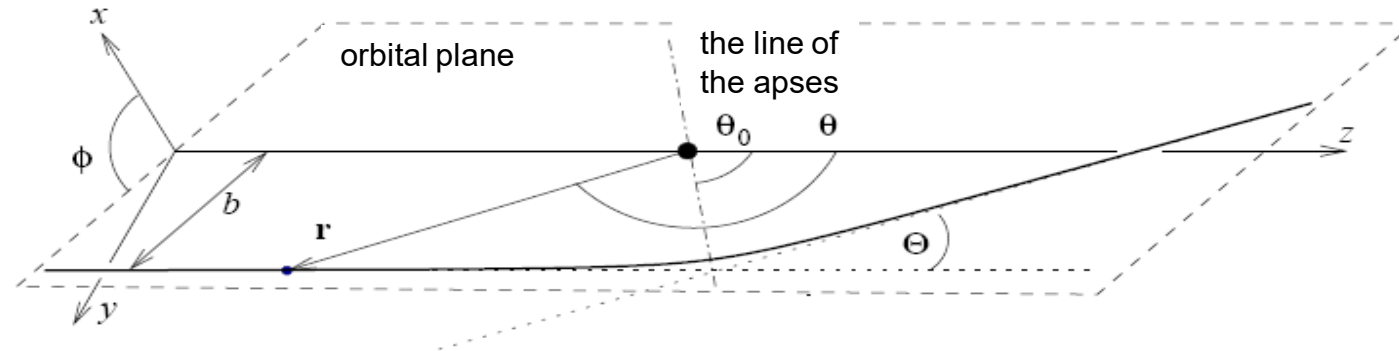
$$E = \frac{-GMm}{r} + \mu \frac{v^2}{2}$$

$$v^2 = G(M+m) \left(\frac{2}{r} - \frac{1}{a} \right)$$

Hohmann transfer orbit



Hyperbolic orbits: Scattering in the gravitational field of a planet



$$r = \frac{p}{1 + \varepsilon \cos(\theta - \theta_0)}$$

$$r \rightarrow \infty \text{ when } \theta \rightarrow \pi \quad 1 + \varepsilon \cos(\pi - \theta_0) = 0 \Rightarrow \cos \theta_0 = \frac{1}{\varepsilon}$$

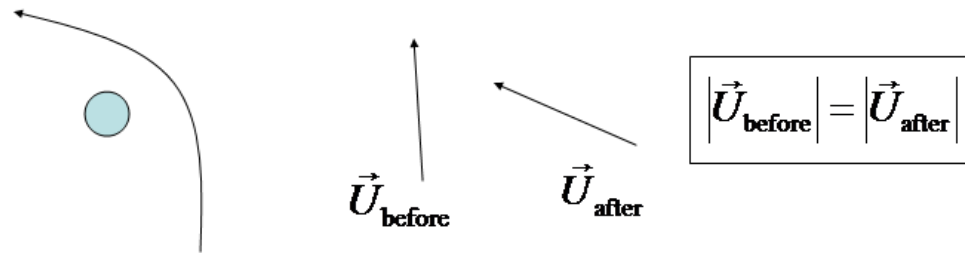
$$\text{Symmetry: } \Rightarrow \Theta = \pi - 2\theta_0$$

$$\cot \frac{\Theta}{2} = \sqrt{\varepsilon^2 - 1} = \sqrt{\frac{2El^2}{mk^2}} = \frac{2Eb}{|k|} \quad b = \frac{|k|}{2E} \cot \frac{\Theta}{2} \quad \text{impact parameter}$$

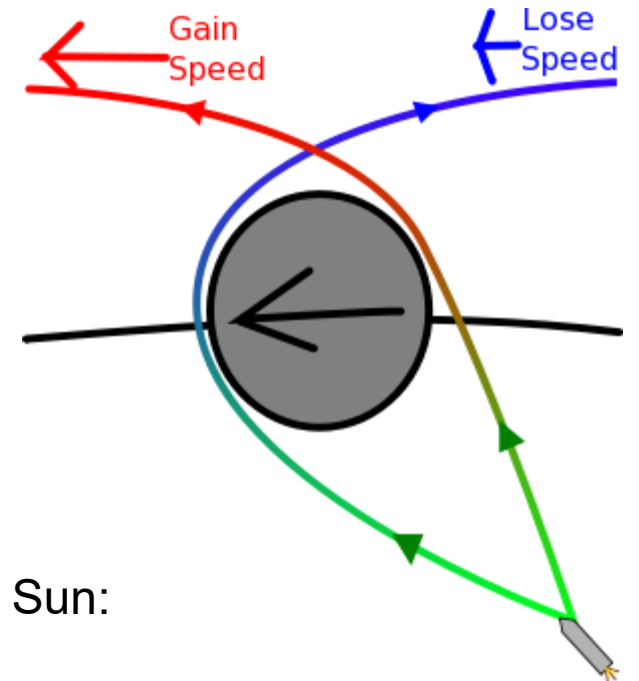
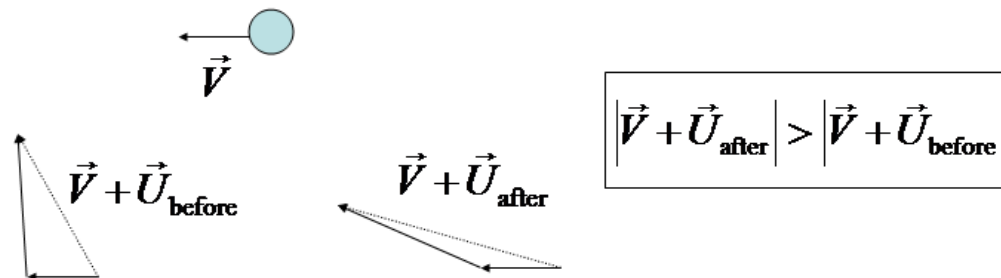


Using a planet to accelerate / decelerate a spacecraft

Frame of Reference: Moving with Planet



Frame of Reference: Planet Moving Left



In the frame of the Sun:

To accelerate
take over the planet from behind

To decelerate
let the planet take over you

Cassini gravity assisted trajectory

VENUS 1 FLYBY
26 APR 1998

VENUS 2 FLYBY
24 JUN 1999

VENUS
TARGETING
MANEUVER
3 DEC 1998

LAUNCH
15 OCT 1997

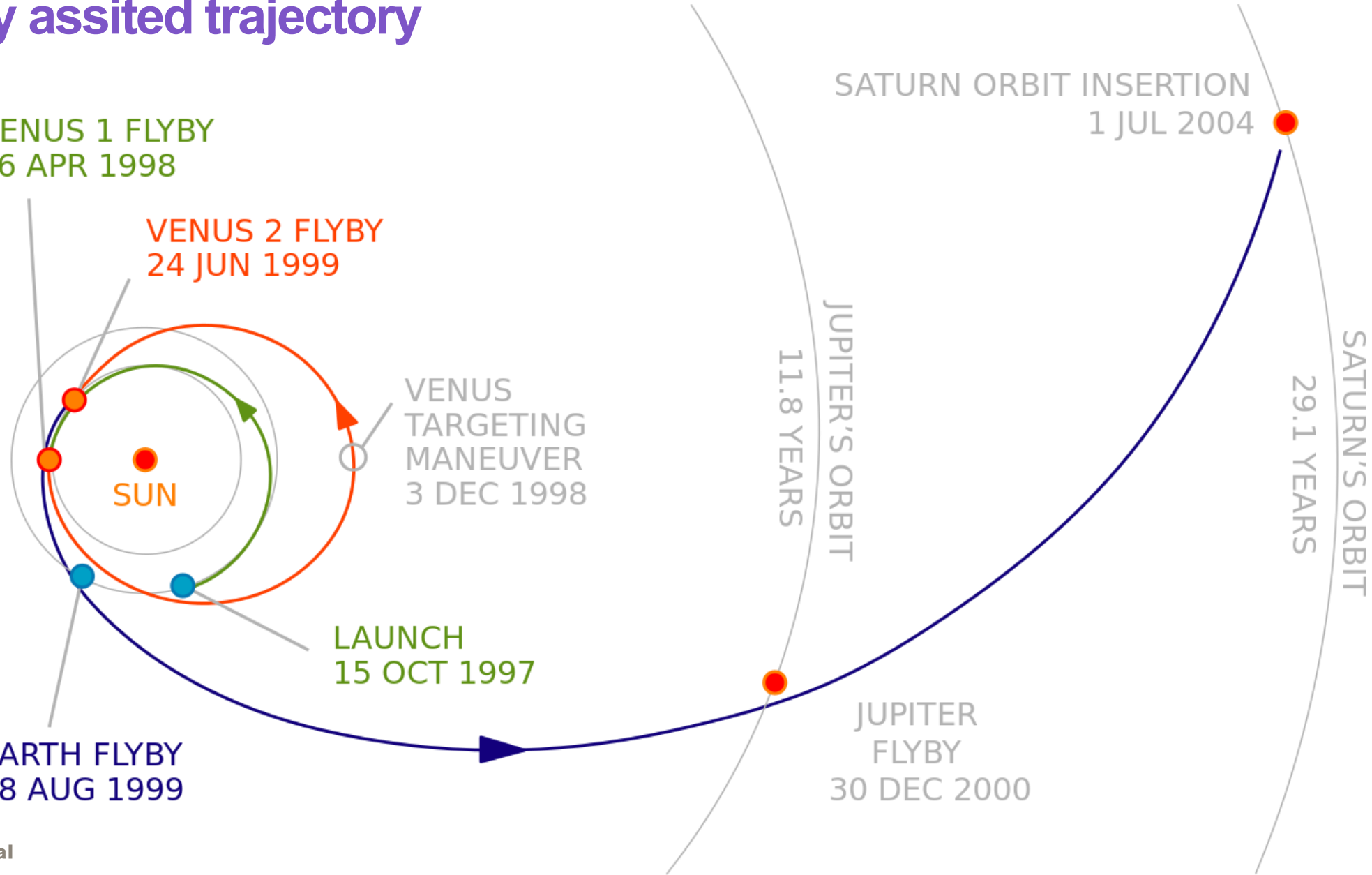
EARTH FLYBY
18 AUG 1999

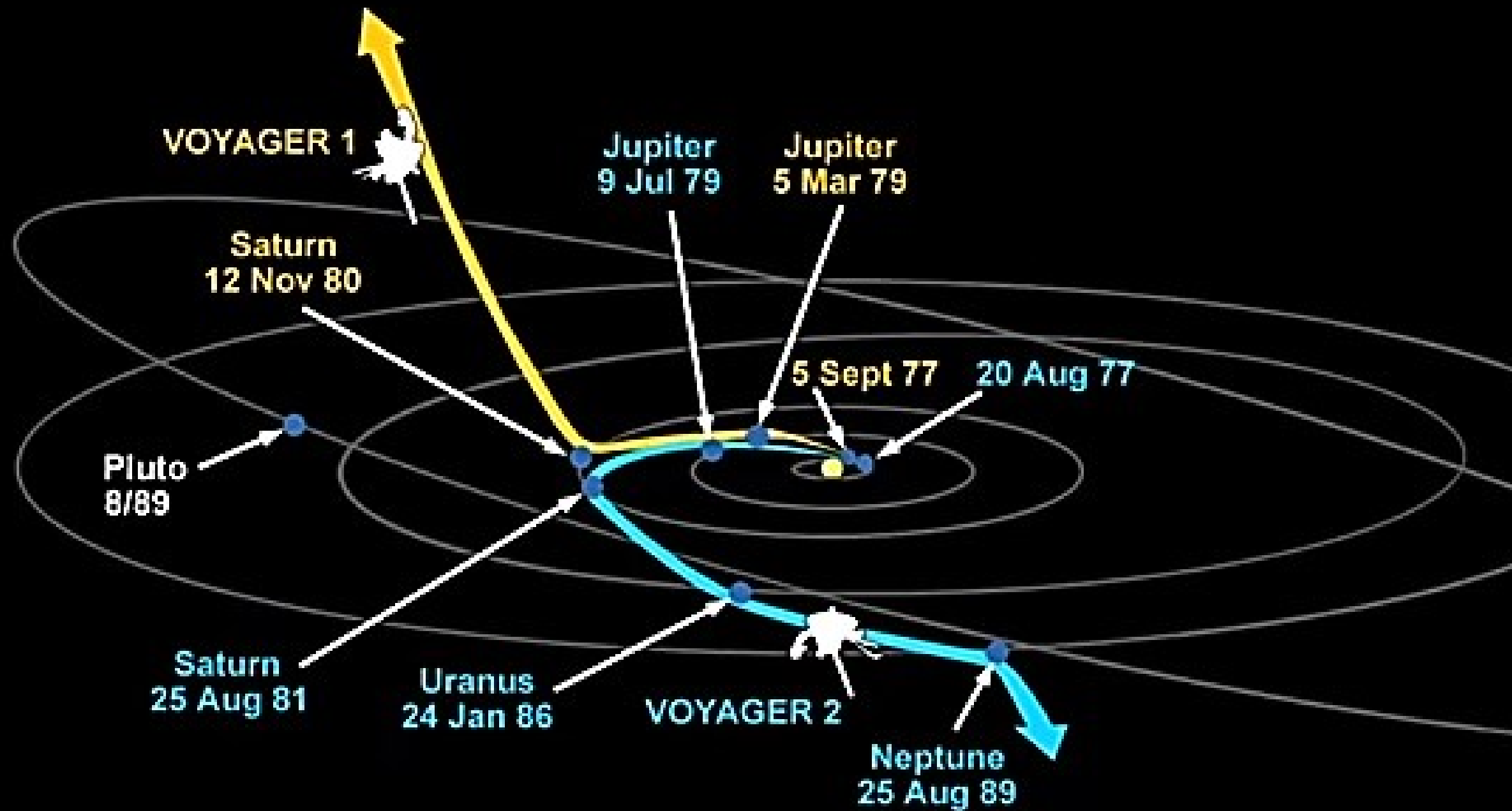
SATURN ORBIT INSERTION
1 JUL 2004

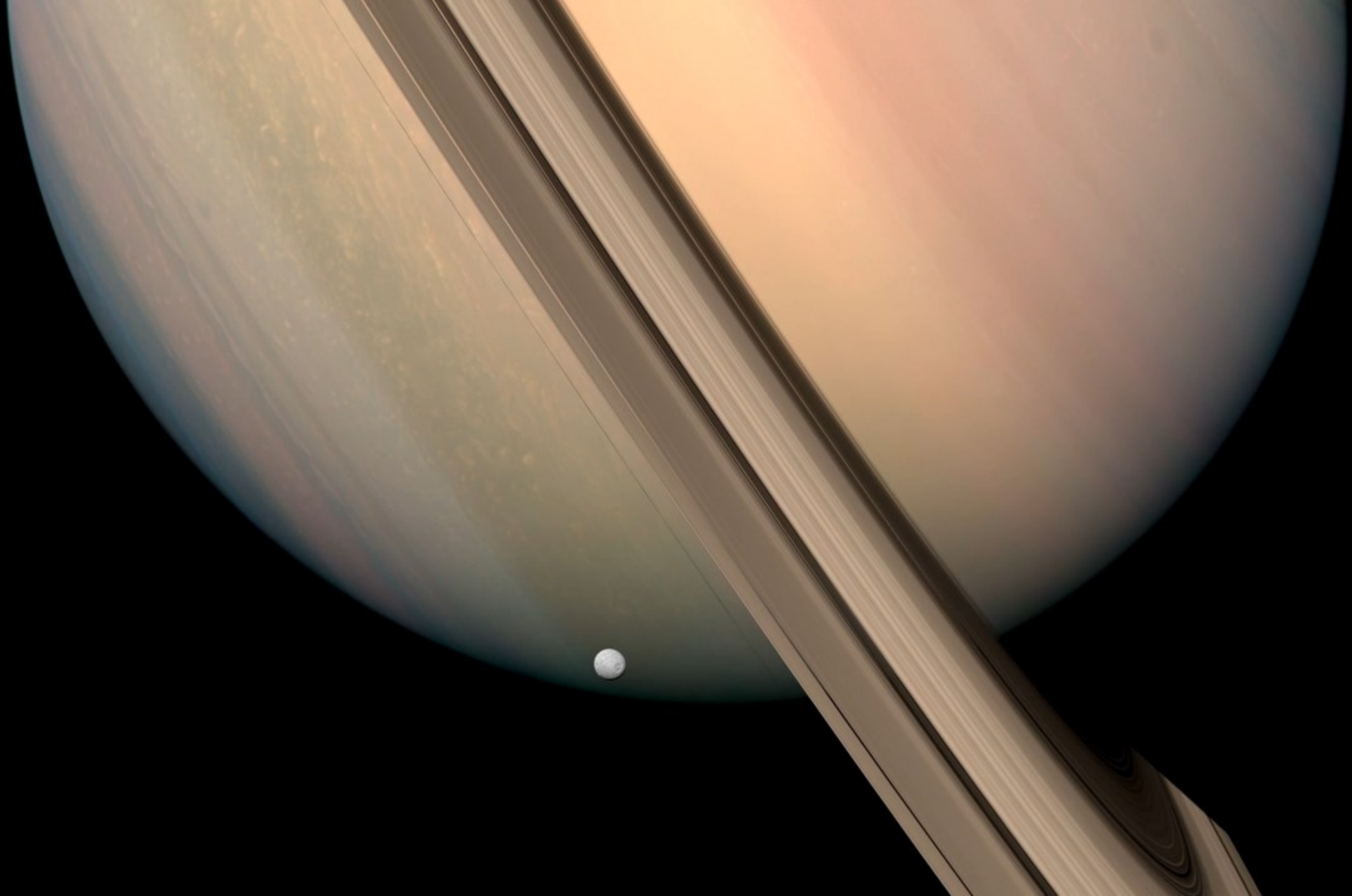
JUPITER'S ORBIT
11.8 YEARS

JUPITER
FLYBY
30 DEC 2000

SATURN'S ORBIT
29.1 YEARS









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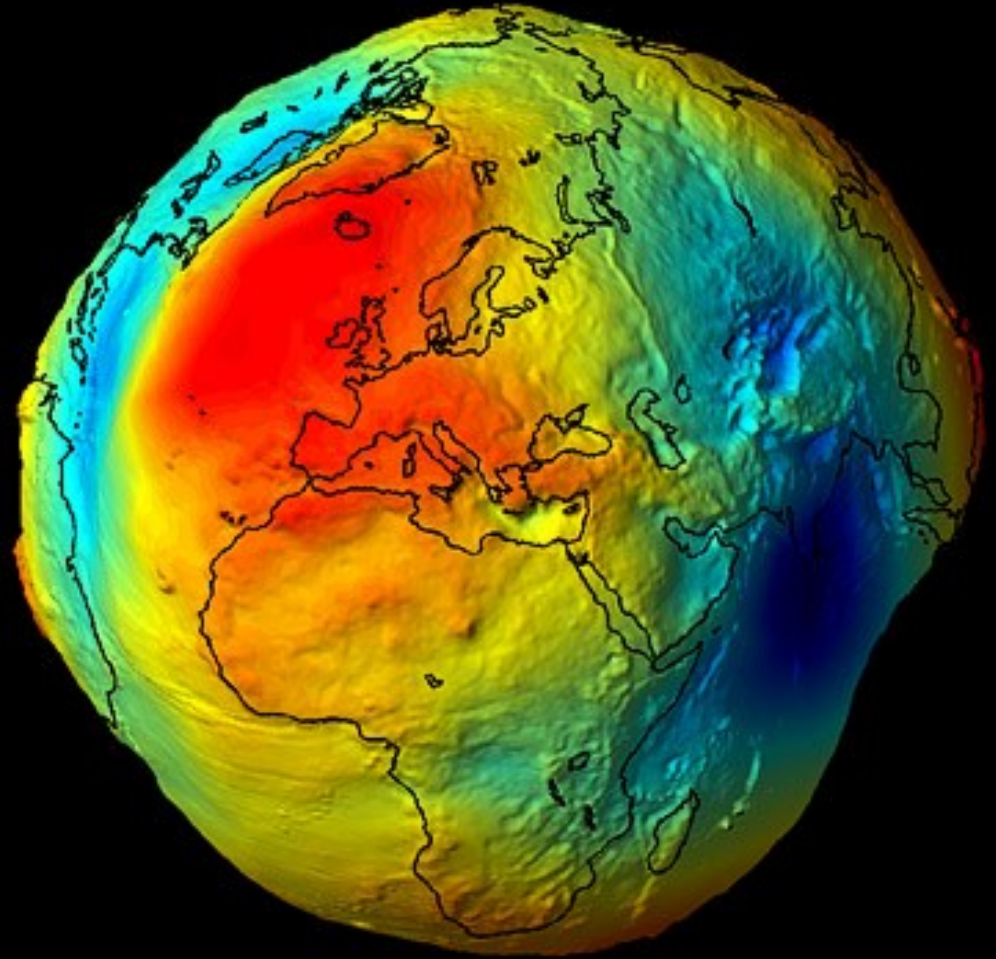
Orbits around real celestial body

Potato #345

Kevin Abosch (2010)



Earth gravity field measured by GOCE



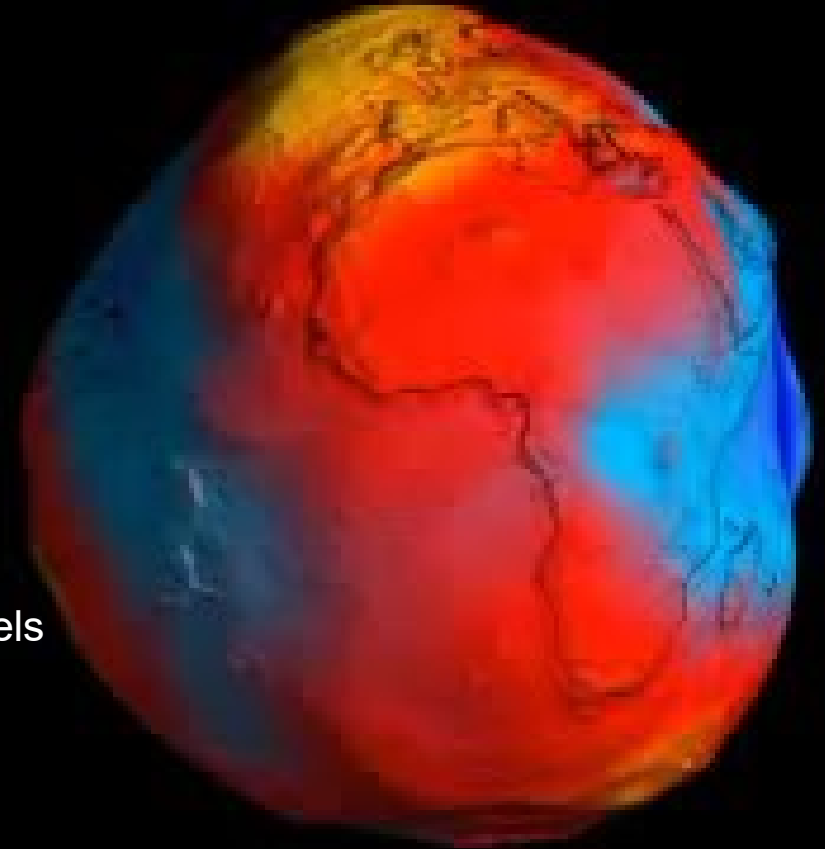


Perturbations to orbits

The two-body approach is a theorist's dream

Other effective forces arise from:

- atmospheric drag
 - largest at perigee and makes an elliptical orbit more circular
 - steady deceleration on the circular orbit lowers the altitude and finally the satellites return to the atmosphere
- radiation pressure of the Sun
 - increased need for orbital corrections
 - telecommunication satellites on GEO have typically large solar panels
- inhomogeneous gravitational field
 - oblateness
 - uneven distribution of matter
- other celestial bodies
 - Moon, Jupiter



Gravitational perturbations

The gravitational potential can be represented as a spherical harmonic expansion

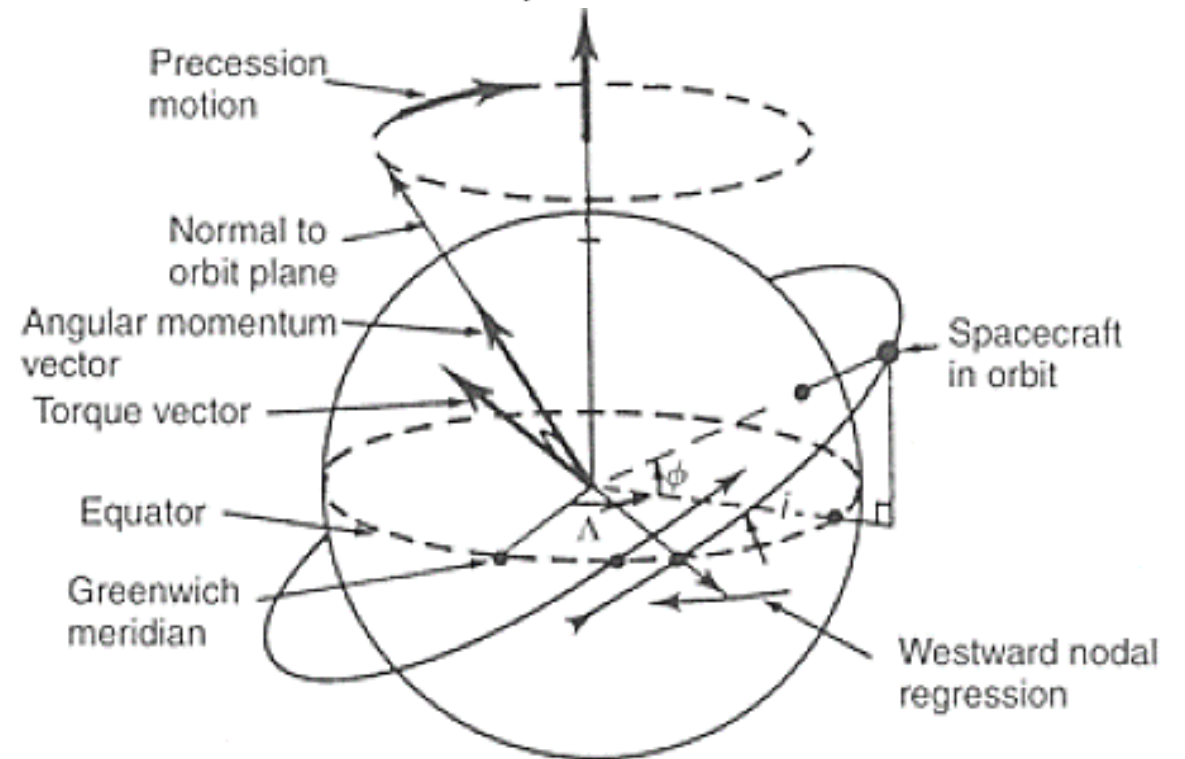
$$U(r, \Phi, \Lambda) = \frac{\mu}{r} \left\{ -1 + \sum_{n=2}^{\infty} \left[\left(\frac{R_E}{r} \right)^n J_n P_{n0}(\cos \Phi) + \sum_{m=1}^n \left(\frac{R_E}{r} \right)^n (C_{nm} \cos m\Lambda + S_{nm} \sin m\Lambda) P_{nm}(\cos \Phi) \right] \right\}$$

The most important contribution is due to the oblateness

- it turns the angular momentum vector
- the nodal line rotates

$$\Delta\Omega = -\frac{3\pi J_2 R_E^2}{p^2} \cos i \frac{\text{rad}}{\text{rev}}$$

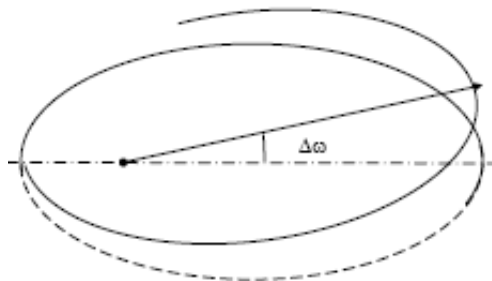
A shift of 360 deg / year results in a Sun-synchronous orbit



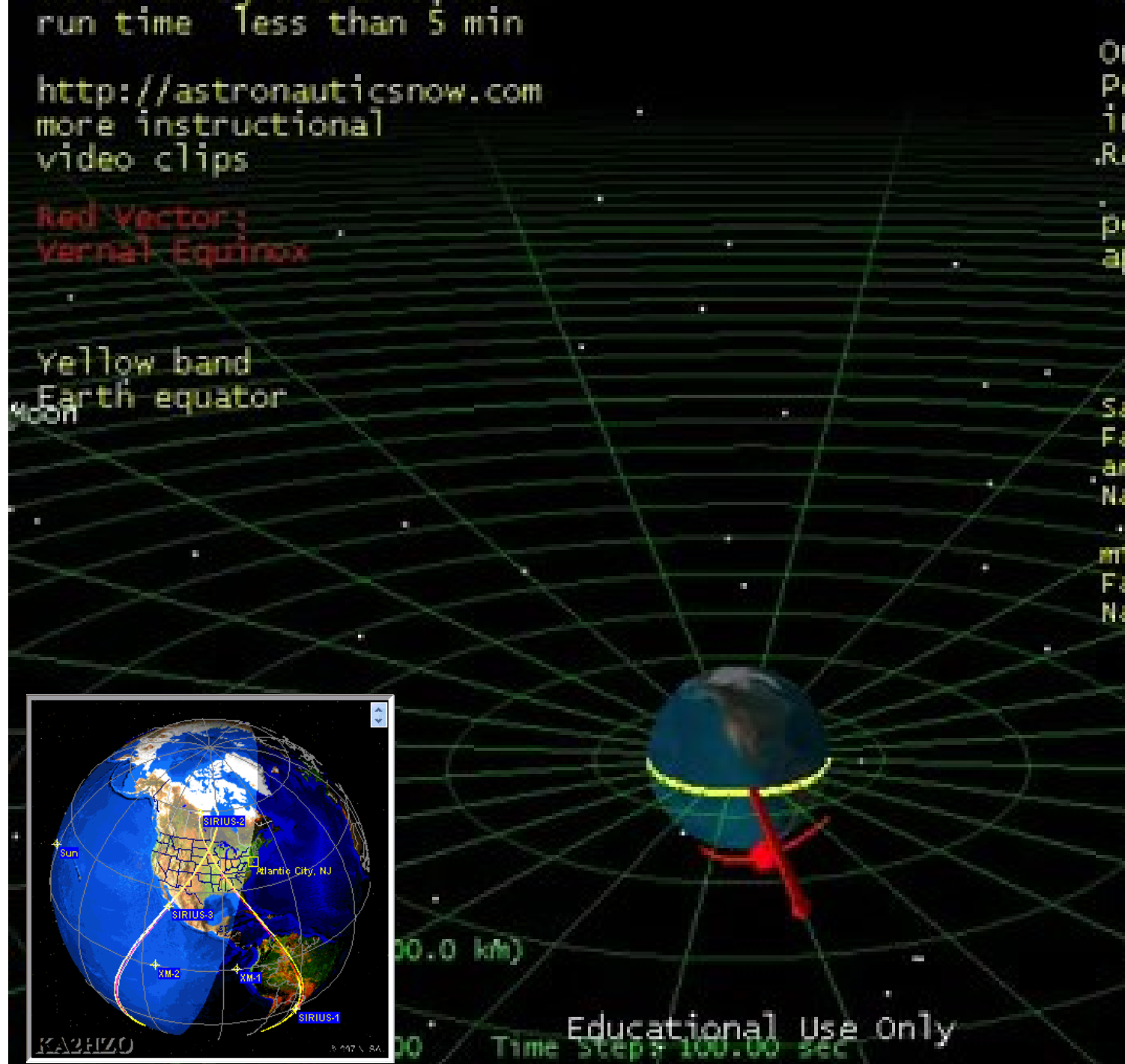
Gravitational perturbations

The oblateness also causes precession of the line of the apsides

$$\Delta\omega = 3\pi \frac{J_2 R_E^2}{p^2} \left(2 - \frac{5}{2} \sin^2 i \right) \frac{\text{rad}}{\text{rev}}$$



Inclination of 63.4 deg is a special case: so-called **Molniya orbits**



Mike Gruntman
file: mikegruntman-06.wmv
run time 5 min 20 sec

Educational Use Only

Prograde and
Retrograde orbits

Regression of Nodes
effect of J2

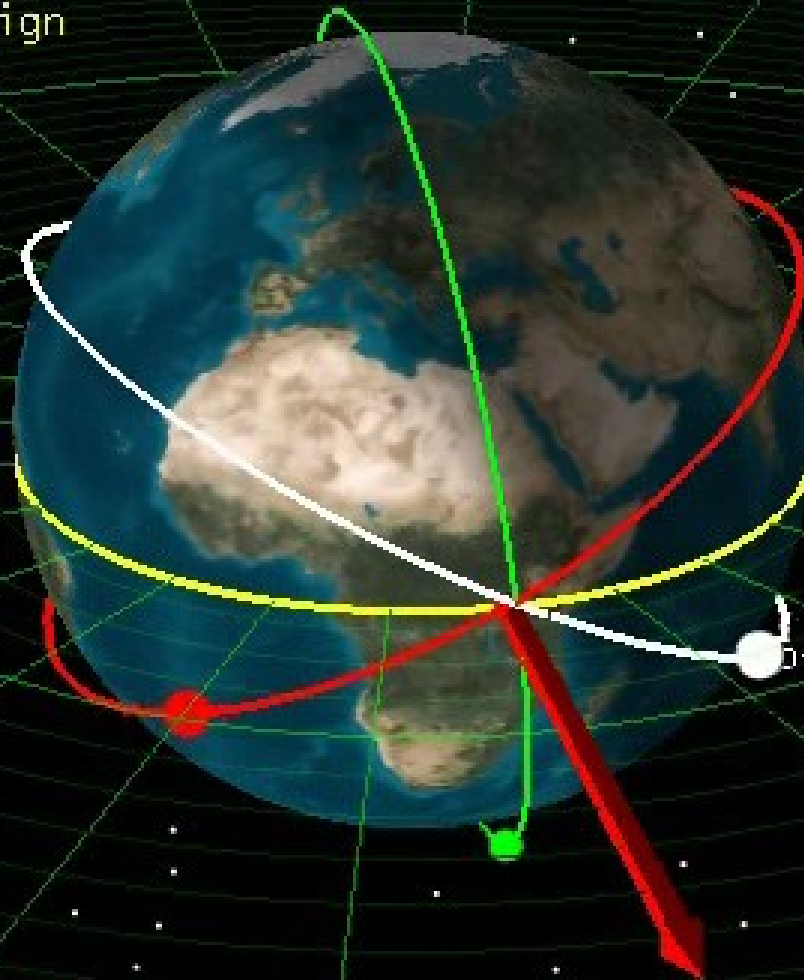
<http://astronauticsnow.com>
video clips of interest
for space mission design
and spacecraft design

Orbit inclination:
red - 28 deg
white - 152 deg
green - 97 deg

Red vector
Vernal Equinox

400-km altitude
circular orbits

Yellow band
Earth equator



Grid 10000.0 km (1000.0 km)

Earth Inertial Axes

1 Jan, 2008 15:13:00.000

Time step: 60.00 sec

Educational Use Only



Effects of other celestial bodies

Three-body problem

For more than 2 bodies the equation of motion

$$m_i \ddot{\mathbf{r}}_i = - \sum_{j \neq i} \frac{G m_i m_j (\mathbf{r}_i - \mathbf{r}_j)}{|\mathbf{r}_i - \mathbf{r}_j|^3} \quad i = 1, 2, \dots, N$$

soon becomes intractable

Poincaré (ca. 1890): the system is very sensitive to initial conditions

➤ chaos (which was reinvented in the 1960s)

K. F. Sundman (1912): there is a unique solution for the 3-body problem

➤ a very slowly converging expansion in powers of $t^{1/3}$

➤ impractical for calculations in celestial mechanics



Reduced three-body problem

Lagrange points

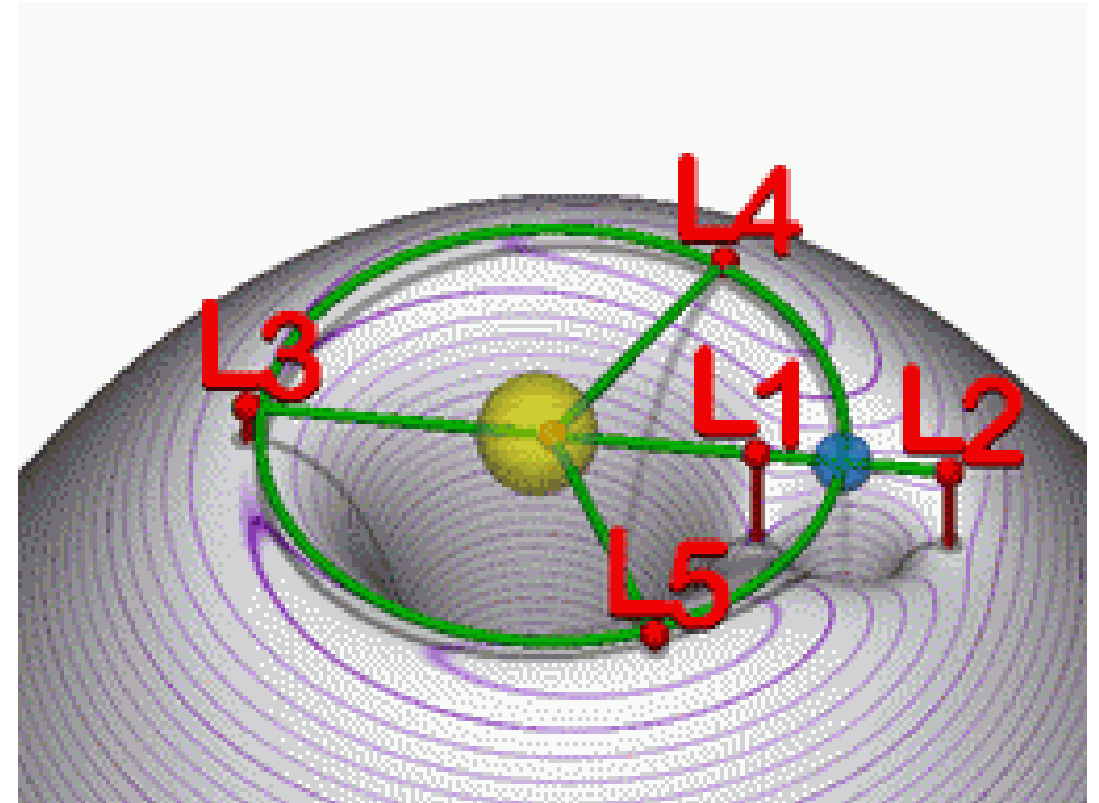
Two large and one small body

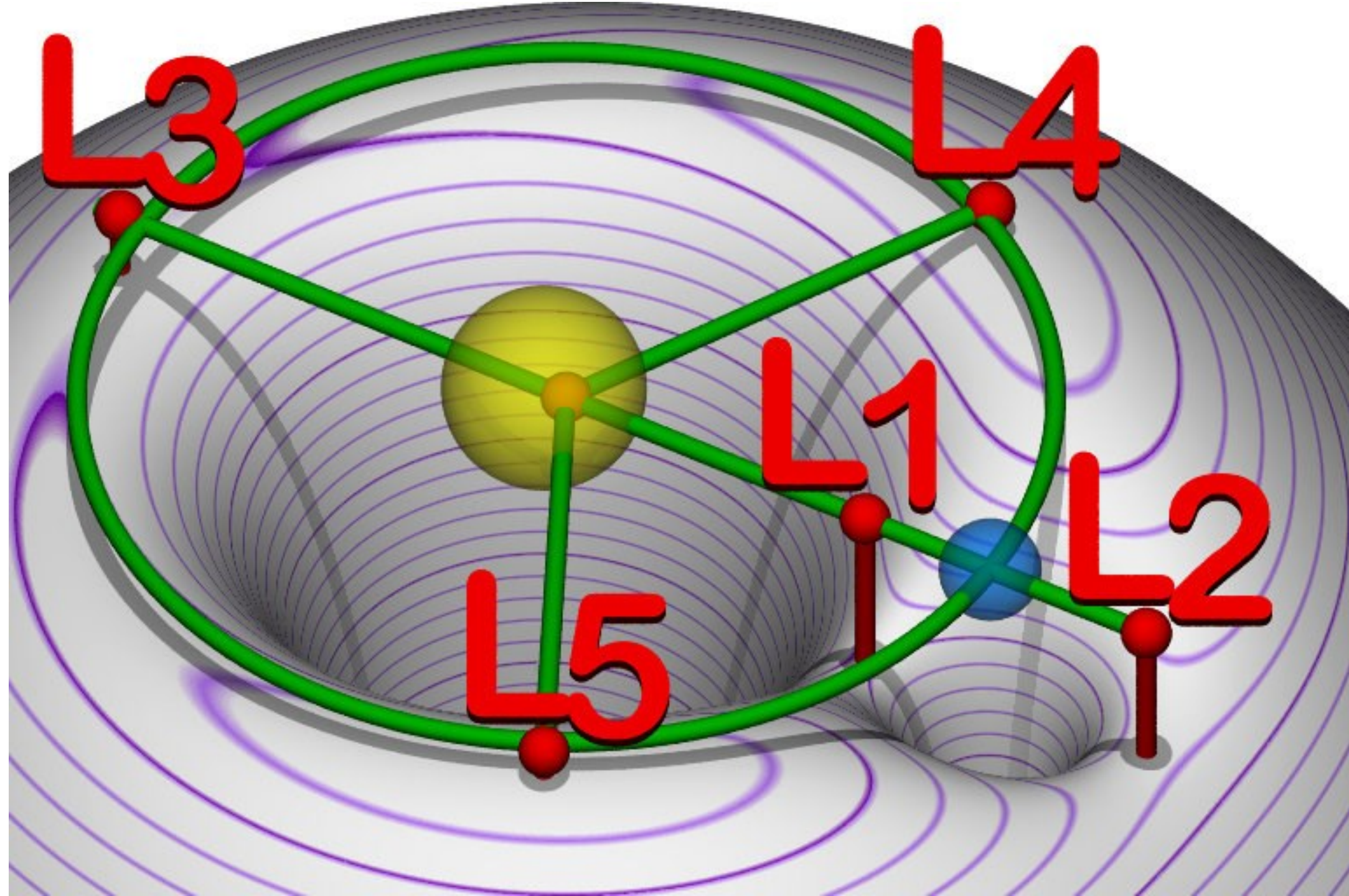
- solve the 2-body problem for the large masses
- consider the motion of the small body in the gravitational potential of the large ones.

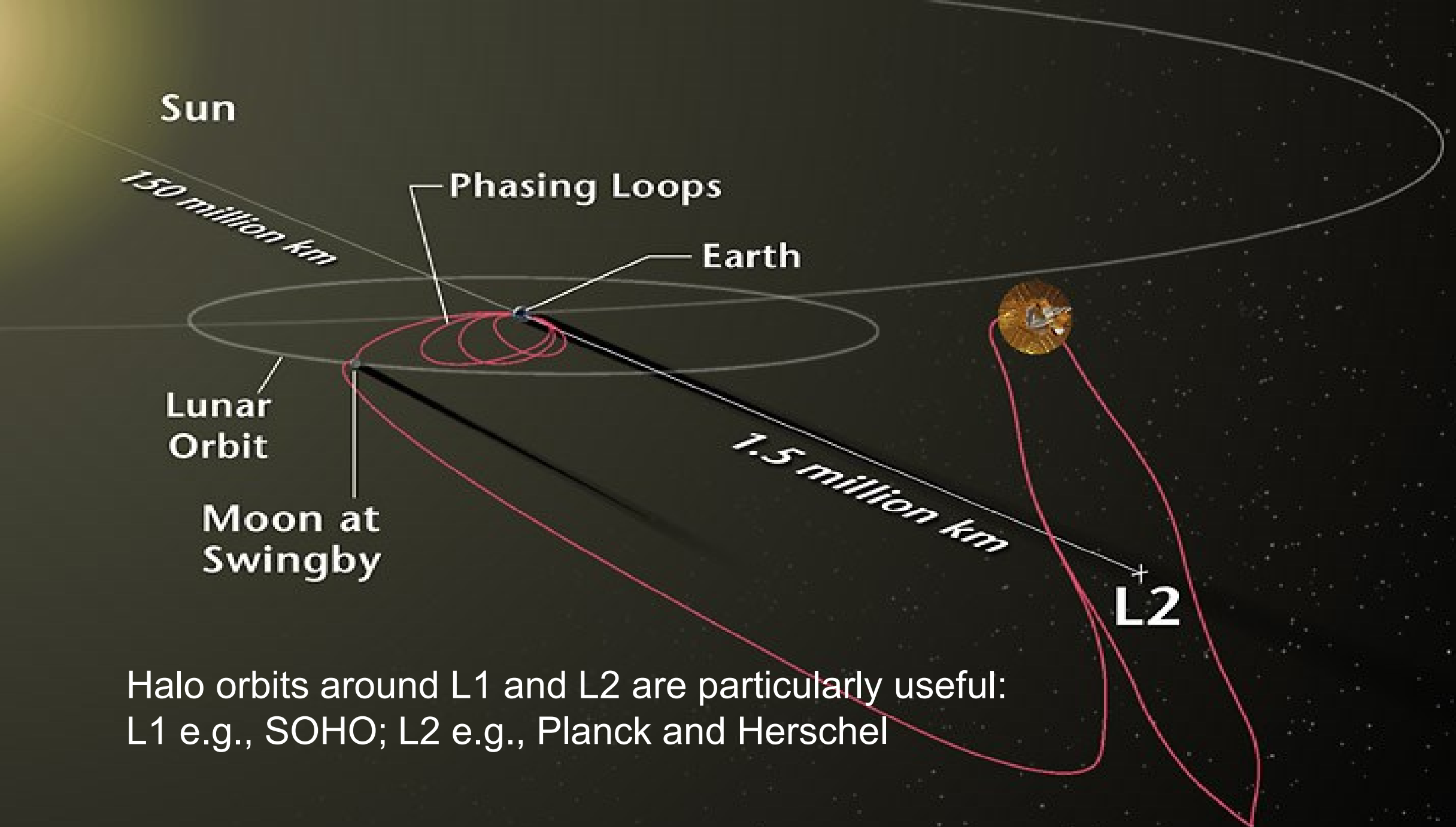
Does not work in the
Sun-Earth-Moon system

but OK for
Sun-Earth-spacecraft

Five Lagrange points

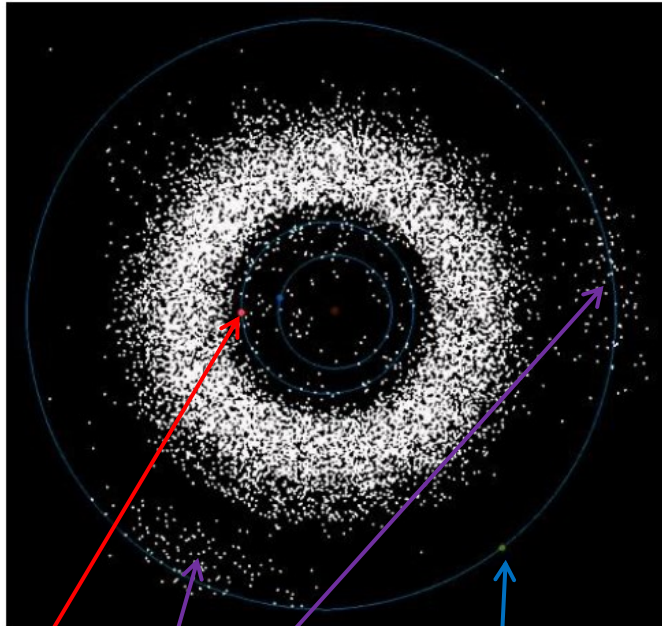






Halo orbits around L1 and L2 are particularly useful:
L1 e.g., SOHO; L2 e.g., Planck and Herschel

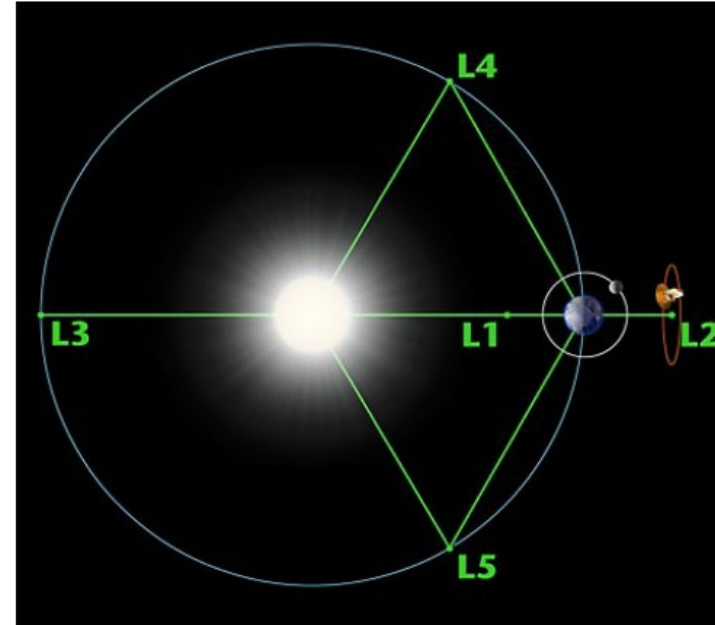
Good to know about Lagrange points



Mars

Jupiter

Trojan asteroids around the Sun–Jupiter system's L4 and L5



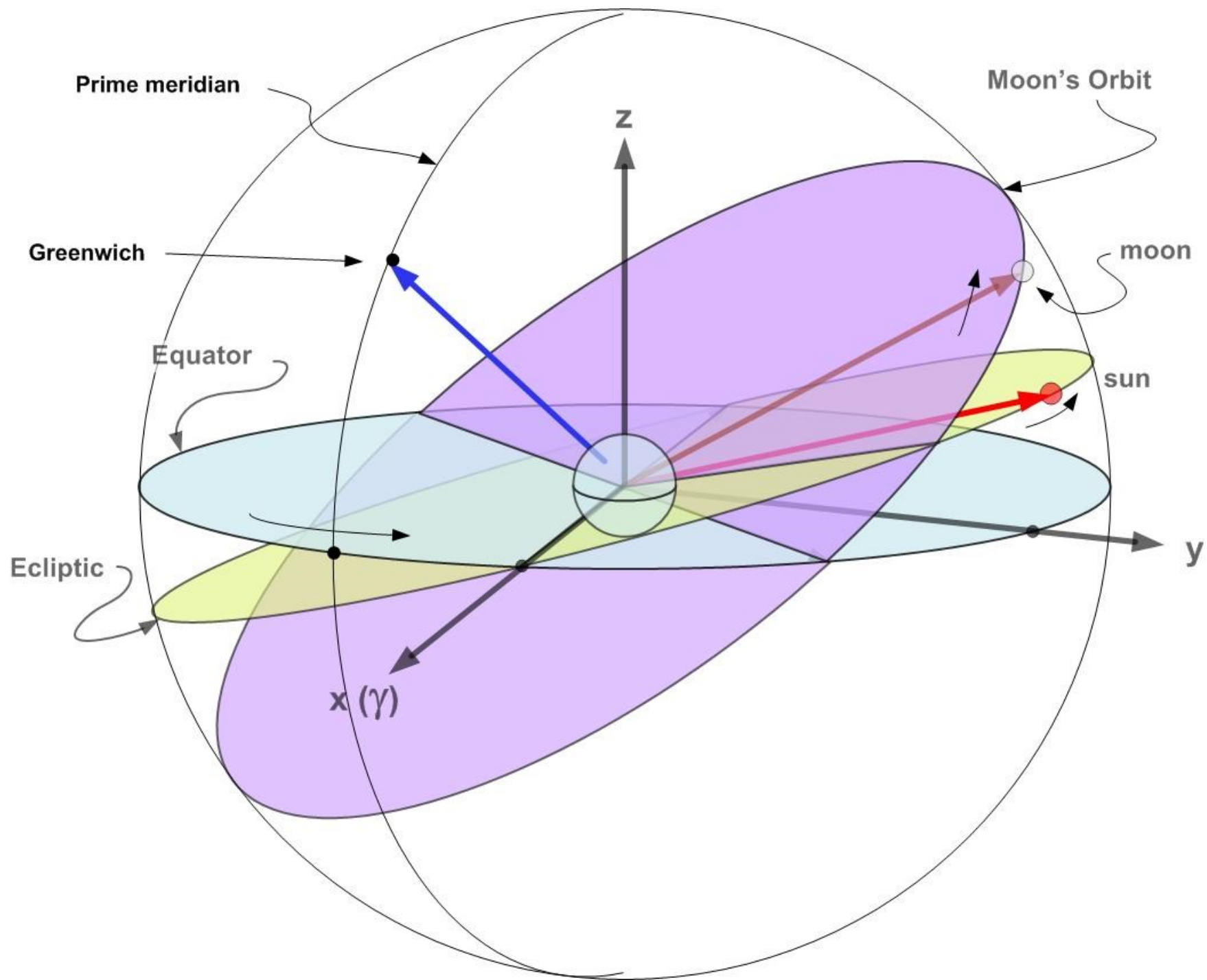
- People often have two misconceptions:
- L1 is NOT the point where gravitational forces of the Sun and the Earth balance each other
 - How can a S/C stay around L2?

These are three-body problems!!



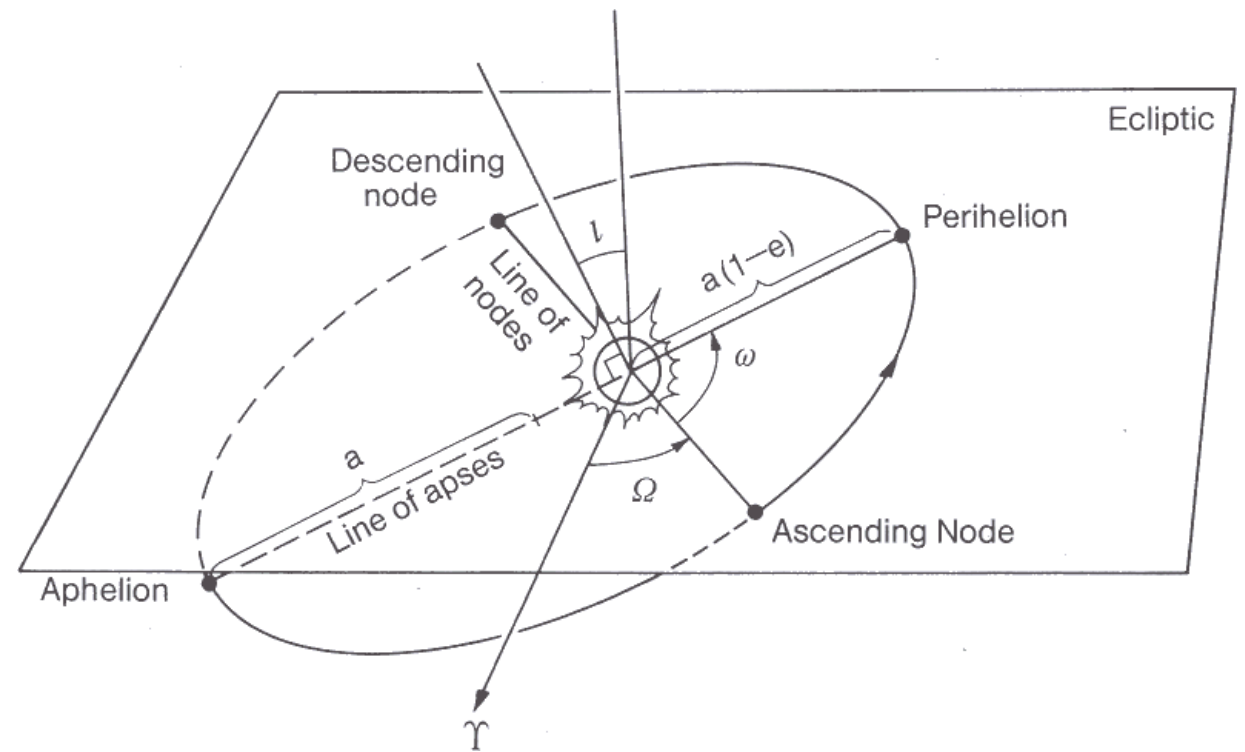
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Describing orbital motion



Orbital elements: Astronomer's view

1. Semi major axis a
2. Eccentricity ε (or e)
3. Inclination i (or ι)
4. Right ascension of the ascending node Ω
(from vernal equinox γ)
5. Argument of perihelion ω ; or length of the perihelion $\varpi = \Omega + \omega$
6. Perihelion time τ



Orbital elements: Satellite operator's view

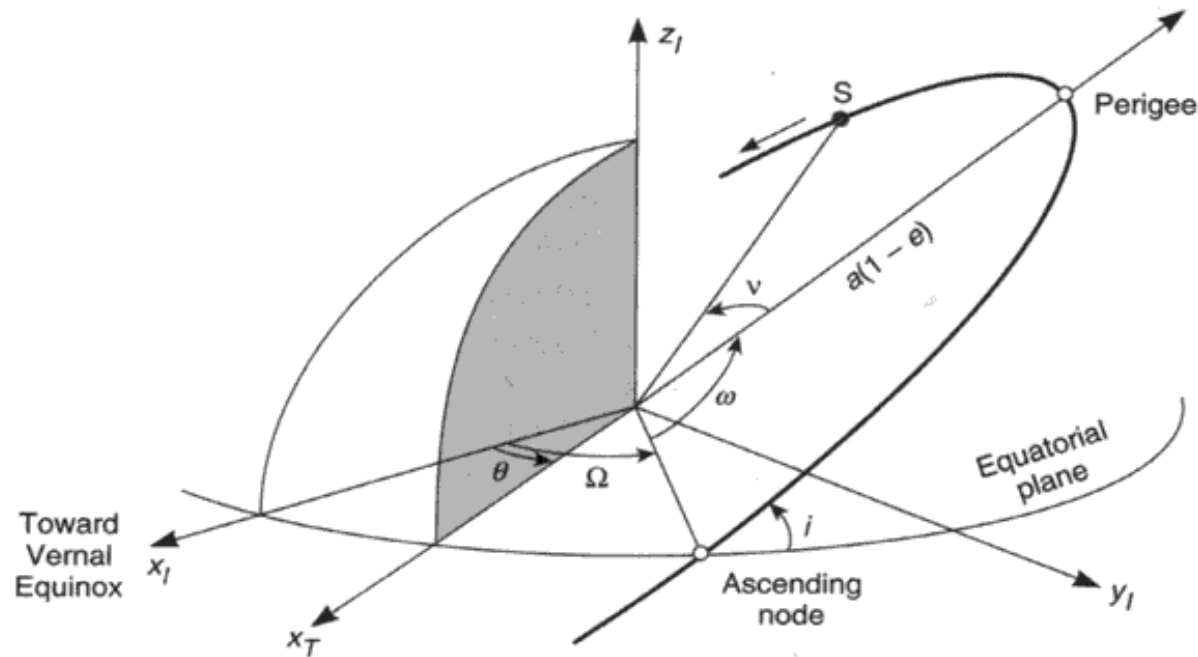


Figure 3.11 Characterization of an ideal orbit and the satellite position by Keplerian elements: $\{a, e, i, \Omega, \omega, \text{ and } \nu\}$.

Also other sets of elements are used; the number of independent elements is always 6

Orbital Elements

Semi-major axis a (size of the ellipse)

- the longest axis of the ellipse going through the two foci.

Eccentricity e (oblateness of the ellipse)

- the elongation of the ellipse

Inclination of the orbit i (orbit position relative to equator)

- the orbit plane is tipped relative to the reference equatorial plane.

Argument of Perigee ω (place of the perigee)

- an angle in the orbital plane between the ascending node and peri-apsis, measured in the direction of the satellite's motion

Right ascension of the ascending node Ω (ellipse rotation around poles)

- the angle measured to the point where the orbit crosses the equatorial plane relative to a reference direction known as the vernal equinox.

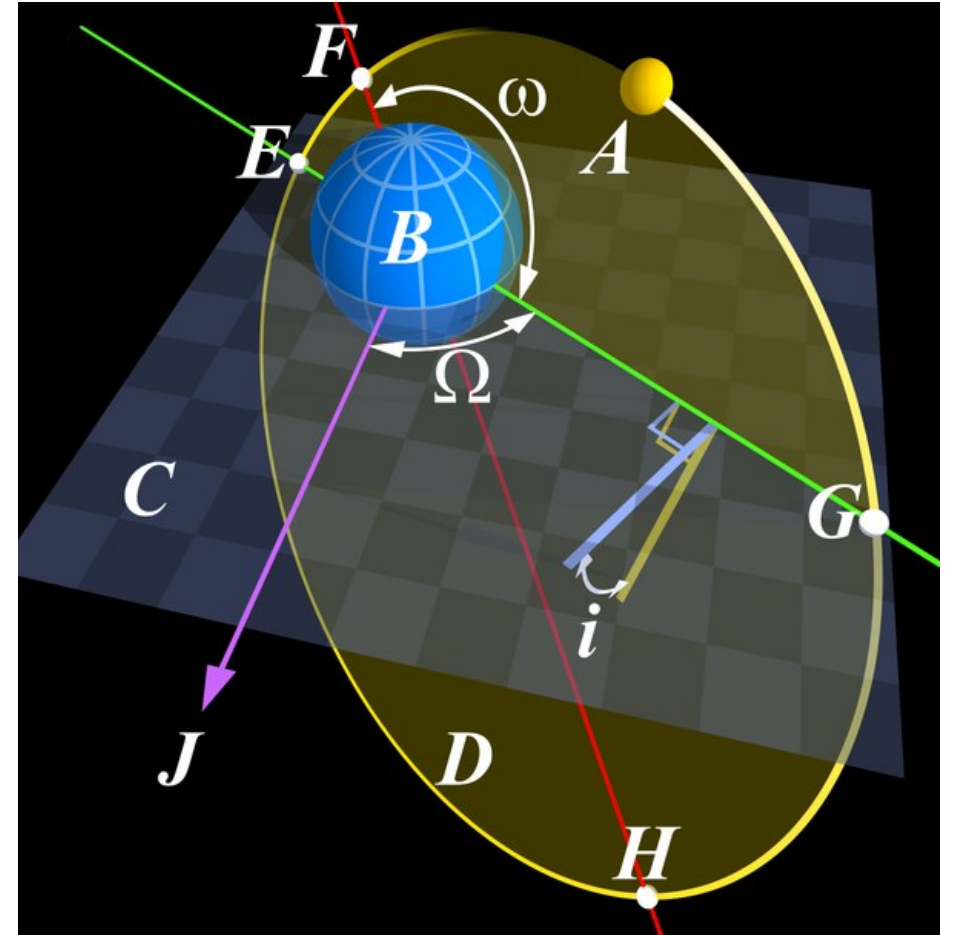
True anomaly ν (place of the satellite on orbit)

- the angular distance of a point in an orbit past the point of peri-apsis, measured in degrees

Parameters of Elliptical Orbit

Legend:

- A – Minor, orbiting body
- B – Major body being orbited by A
- C – Reference plane, e.g. the
- D – Orbital plane of A
- E – Descending node
- F – Periapsis
- G – Ascending node
- H – Apoapsis
- i – Inclination
- J – Reference direction; for orbits in or near the ecliptic, usually the vernal point
- Ω – Longitude of the ascending node
- ω – Argument of the periapsis
- The red line is the line of apsides; going through the periapsis (F) and apoapsis (H); this line coincides with the major axis in the elliptical shape of the orbit
- The green line is the node line; going through the ascending (G) and descending node (E); this is where the reference plane (C) intersects the orbital plane (D).



[Source: Wikipedia]

Source: www.wikipedia.org



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Artificial satellites around the Earth

Satellite orbit classification

LEO: low Earth orbit

- lowest stable orbit at about 180 km
- often polar orbits
- Earth observation and military satellites

MEO: medium altitudes

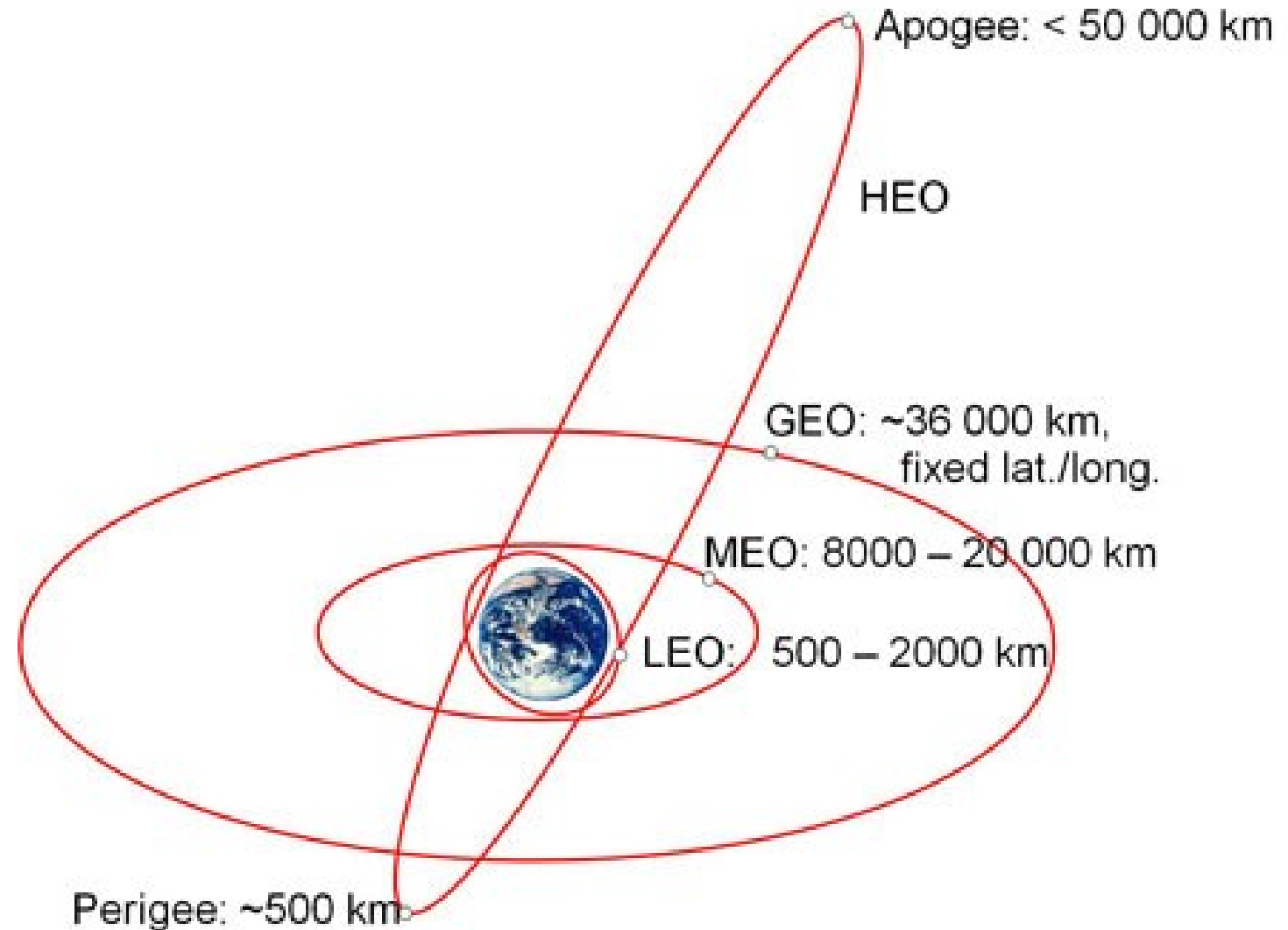
- e.g. GPS and Galileo

GEO: geostationary orbit

- Altitude: 35786 km, geocentric distance 6.6 RE

HEO: highly elliptical orbit

- Apogee above GEO
- some scientific S/C
- Molnya



Types of Orbits

Altitude classifications

- **Low Earth orbit (LEO):** Geocentric orbits with altitudes up to 2,000 km
- **Medium Earth orbit (MEO):** Geocentric orbits ranging in altitude from 2,000 km to just below geosynchronous orbit at 35,786 km.
- **High Earth orbit:** Geocentric orbits above the altitude of geosynchronous orbit 35,786 km

Centric classifications

- **Galactocentric orbit:** An orbit about the center of a galaxy. The Sun follows this type of orbit
- **Heliocentric orbit:** An orbit around the Sun. In our Solar System, all planets, comets, and asteroids are in such orbits, as are many artificial satellites and pieces of space debris. Moons by contrast are not in a heliocentric orbit but rather orbit their parent planet.
- **Geocentric orbit:** An orbit around the planet Earth, such as that of the Moon or of artificial satellites.

Types of Orbits

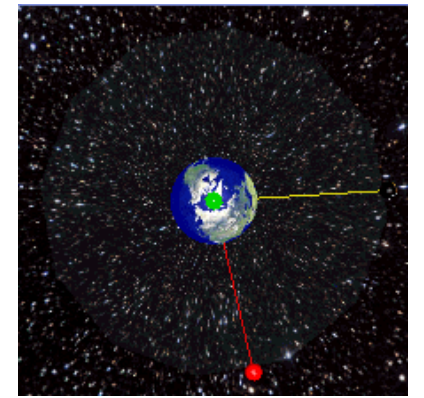
Eccentricity classifications

- There are two types of orbits: ***closed (periodic) orbits***, and ***open (escape) orbits***. Circular and elliptical orbits are closed. Parabolic and hyperbolic orbits are open.
- ***Circular orbit***: An orbit that has an eccentricity of 0 and whose path traces a circle.
- ***Elliptic orbit***: An orbit with an eccentricity greater than 0 and less than 1 whose orbit traces the path of an ellipse.
- ***Parabolic orbit***: An orbit with the eccentricity equal to 1. Such an orbit also has a velocity equal to the escape velocity and therefore will escape the gravitational pull of the planet. If the speed of a parabolic orbit is increased it will become a ***hyperbolic orbit***.

Inclination classifications

- ***Inclined orbit***: An orbit whose inclination in reference to the equatorial plane is not 0.
- ***Non-inclined orbit***: An orbit whose inclination is equal to zero with respect to some plane of reference.

Types of Orbits



Synchronicity classifications

Synchronous orbit:

An orbit whose period is a rational multiple of the average rotational period of the body being orbited and in the same direction of rotation as that body. This means the track of the satellite, as seen from the central body, will repeat exactly after a fixed number of orbits.

- **Geosynchronous orbit (GSO):** An orbit around the Earth with a period equal to one sidereal day, which is Earth's average rotational period of 23 hours, 56 minutes, 4.091 seconds. For a nearly circular orbit, this implies an altitude of approximately 35,786 km. The orbit's inclination and eccentricity may not necessarily be zero.
- **Geostationary orbit (GEO):** A circular geosynchronous orbit with an inclination of zero. To an observer on the ground this satellite appears as a fixed point in the sky

Two Line Elements (TLE)

North American Aerospace Defense Command



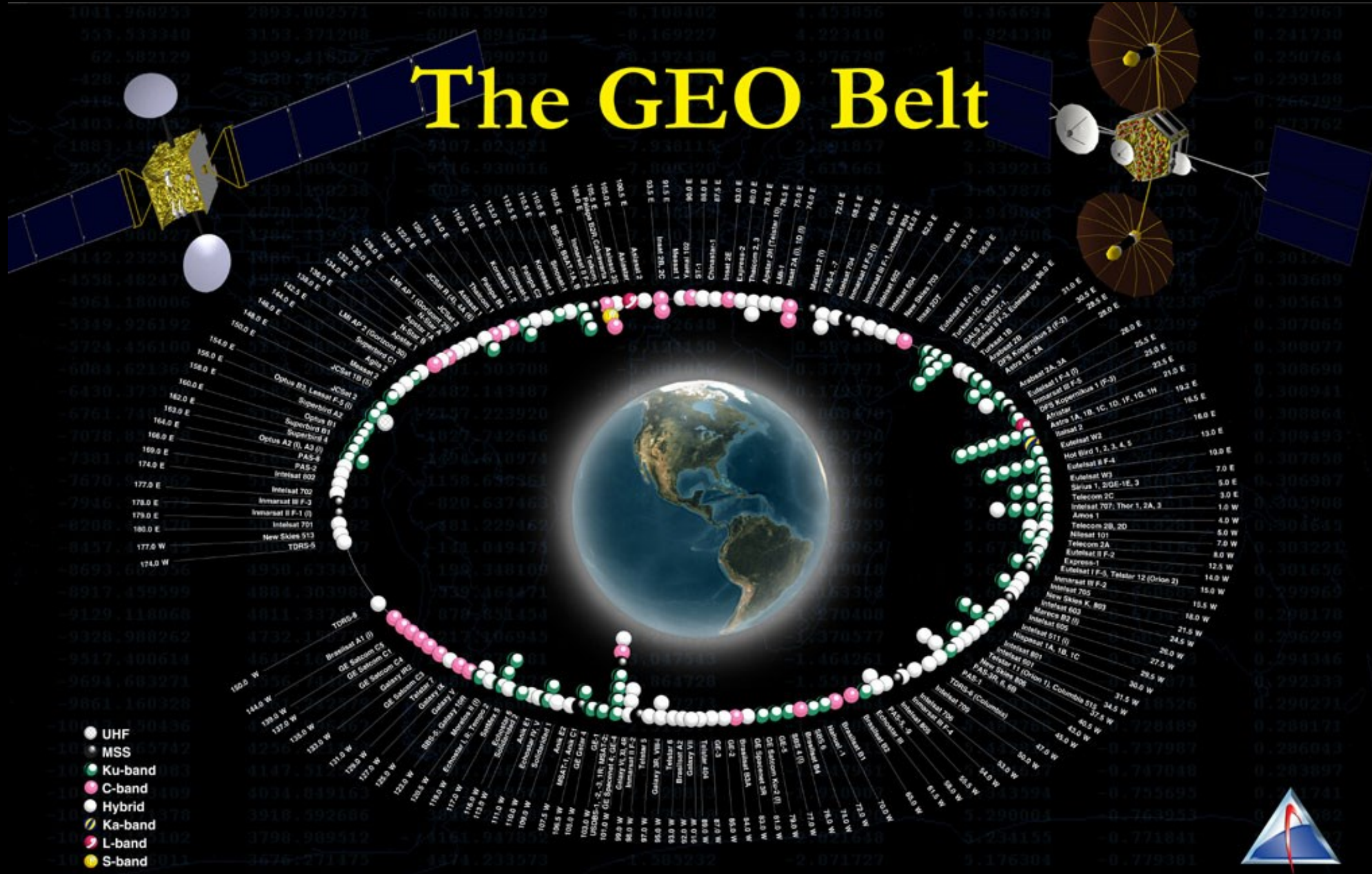
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ISS (ZARYA)
1 25544U 98067A 14260.21469767 .00018739 00000-0 33124-3 0 989
2 25544 51.6477 3.4590 0002230 112.1112 330.1591 15.50366882905605
TIANGONG 1
1 37820U 11053A 14260.09847079 .00037692 00000-0 34424-3 0 3932
2 37820 42.7690 118.2614 0007048 254.3650 242.2970 15.67126632170471
FLOCK 1-12
1 39528U 98067DT 14260.08687349 .00182287 00000-0 90427-3 0 6018
2 39528 51.6355 341.6772 0005790 353.2487 6.8436 15.80898745 33599
FLOCK 1-18
1 39556U 98067DZ 14260.09390708 .00171334 00000-0 89242-3 0 5596
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FLOCK 1-20
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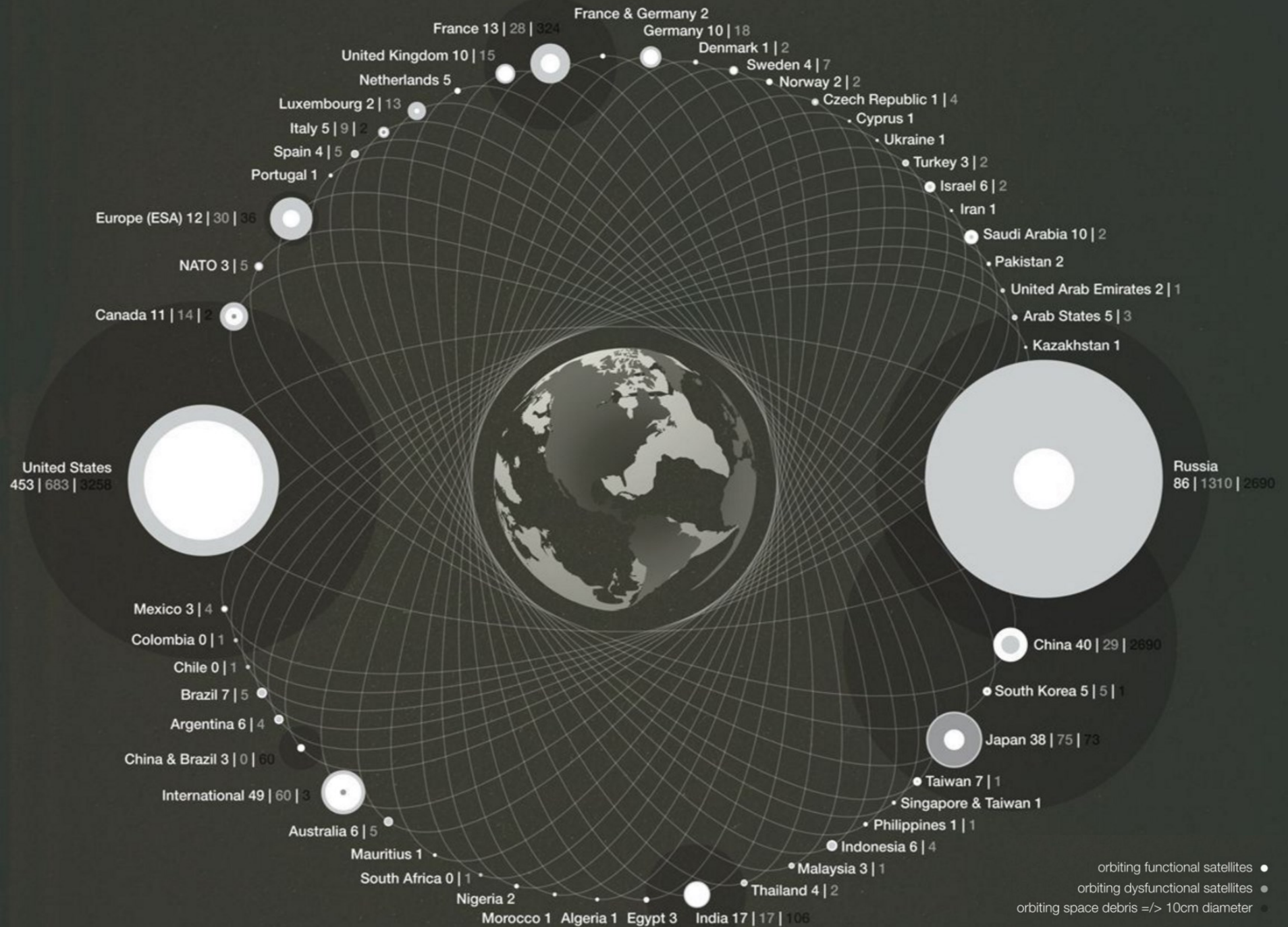




'Urgent need'
to remove space debris

The GEO Belt





● orbiting functional satellites
 ● orbiting dysfunctional satellites
 ● orbiting space debris => 10cm diameter

WASTE IN SPACE

Currently, a thick band of levitating space junk—composed primarily of broken satellite pieces and discarded rocket boosters—skirts the Earth. Two or three times a day, a satellite circling our planet narrowly misses a torrent of the orbital debris. This phenomenon has jeopardized not only current space travelers, but future missions as well.

WHAT IS SPACE DEBRIS?

Nonfunctional, human-made materials in orbit caused by everything from spent booster stages to satellite collisions and explosions.

73%

of tracked debris reside in low-Earth orbit (LEO), 1,200 miles above our planet's surface.

HOW MUCH SPACE JUNK IS UP THERE?

The amount of space debris larger than four inches in diameter in Earth's orbit being tracked by the U.S. Space Surveillance Network:

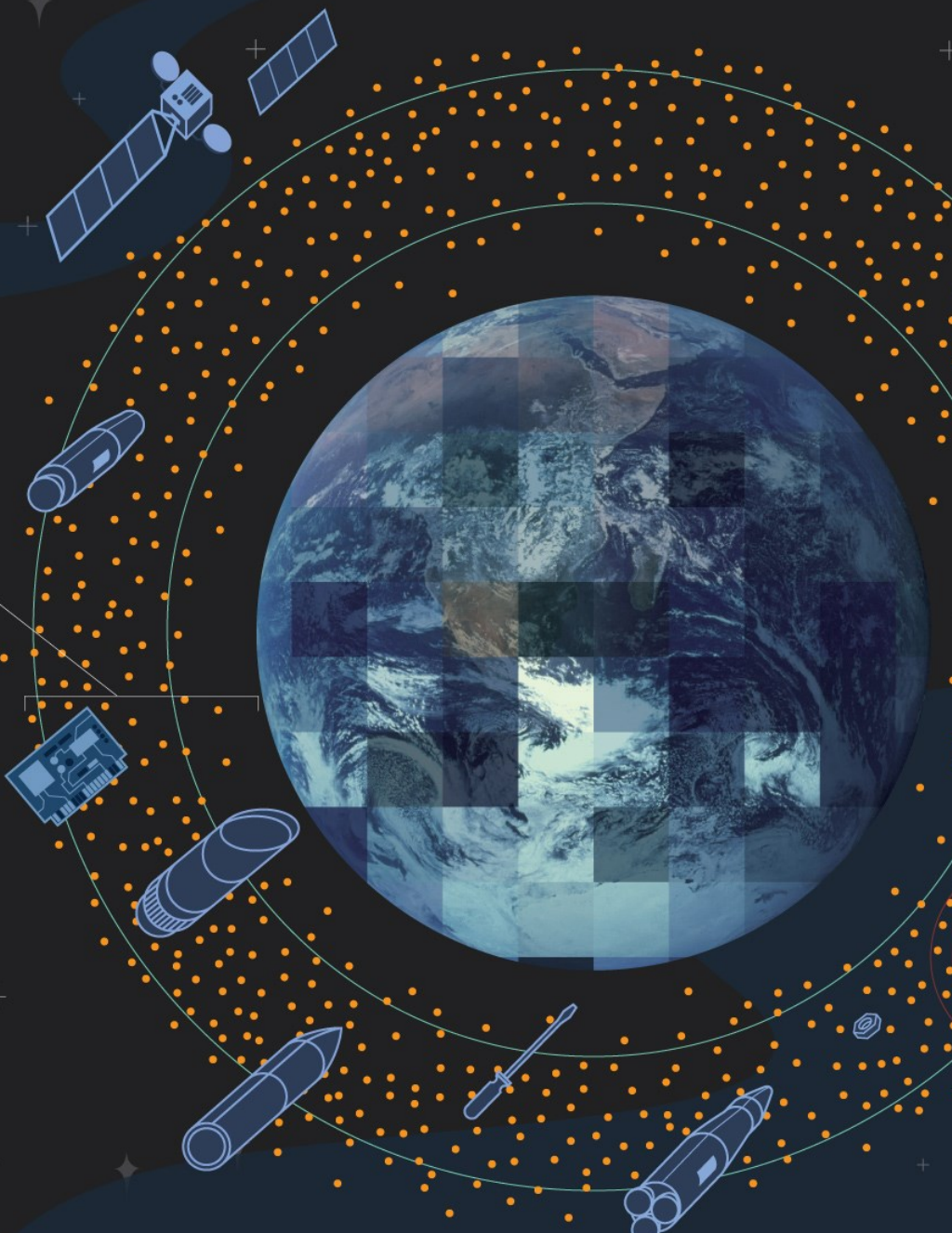
More than **21,000** objects



500,000 objects

Estimated amount larger than one centimeter in diameter—or the size of a marble.

There are another tens of millions of paint chip-like pieces that measure smaller than a centimeter.



WHY IT'S A SERIOUS PROBLEM

Traveling at such hyper velocities, any particle of space junk presents a considerable threat to spaceflight for any nation. And with more hardware flying around Earth's orbit, the potential of collisions between spacecraft and large orbital trash only continues to grow.

FASTER THAN THE SPEED OF SOUND

The speed of sound travels at approximately **768 mph** on a normal day.

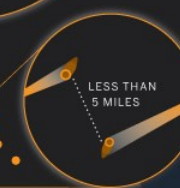
In order to remain in orbit, the fragments in space have to move along at least **20 times that speed**, and can go up to almost

18,000 mph.



TOO CLOSE FOR COMFORT

About 1,000 times a day, satellites and debris pass less than 5 miles from each other. Considering how expansive space is, this distance is striking.



COLLISIONS & EXPLOSIONS INCREASE DEBRIS

CHINA'S ANTI-SATELLITE MISSION

In 2007, China intentionally destroyed one of their weather satellites in space, and the event led to a



900-piece cloud of debris.

THE FIRST MAJOR IMPACT

February 10, 2009:

The 15,000 mph collision of the private Iridium 33 satellite and Cosmos 2251, a Russian military spacecraft, left a trail of approximately 2,000 pieces of low-Earth orbit debris.



Together, these two events combined increased the number of debris in low-Earth orbit by

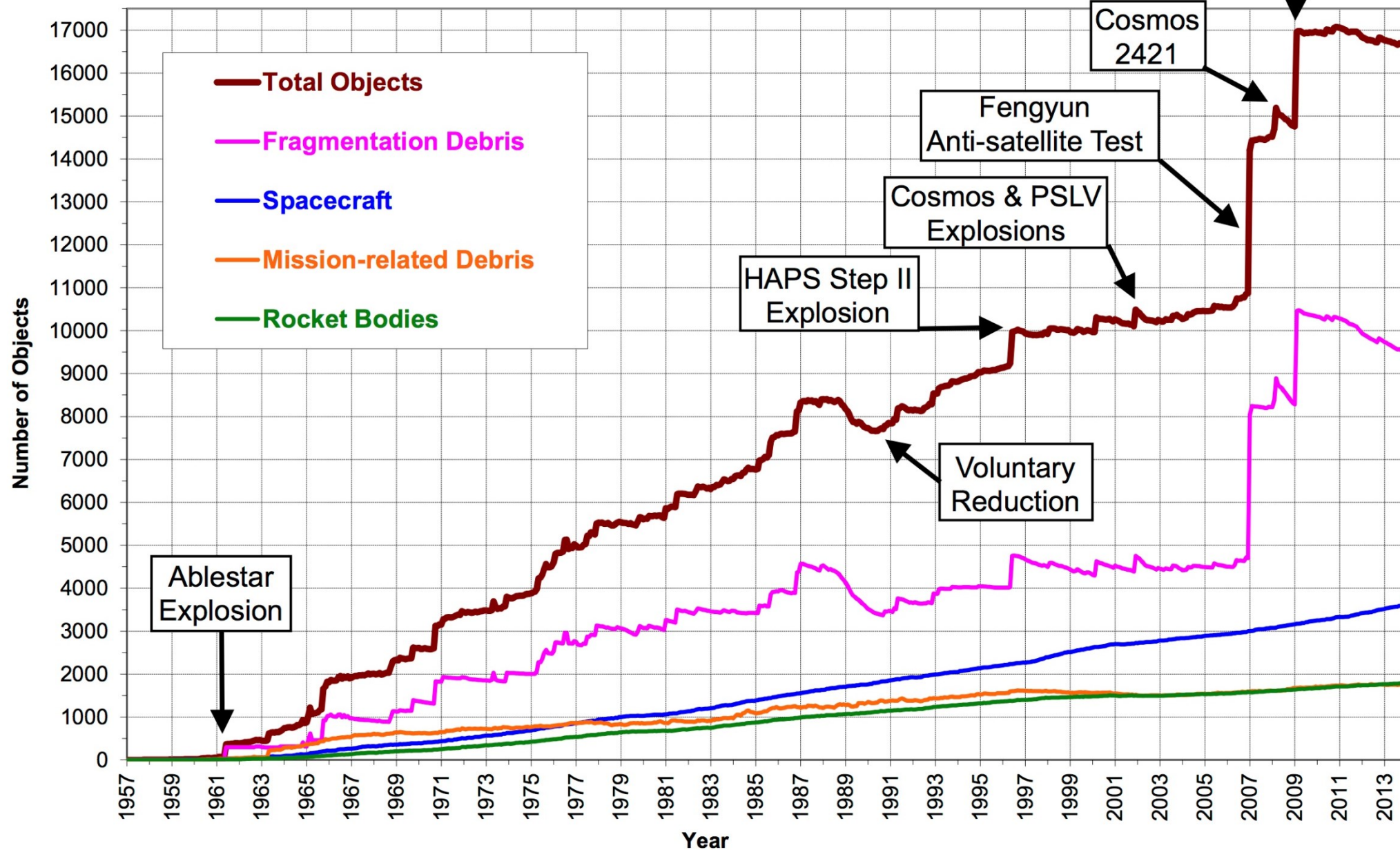
more than 60%



That's taking into account everything that has accumulated over the past 50 years.

Space Debris

Monthly Number of Objects in Earth Orbit by Object Type



Assignment

- **Draw Earth, and a satellite orbiting the Earth.**
 - **Draw the orbit of the satellite.**
 - **Mark rotation direction of the Earth.**
 - **Change your picture with your neighbour. Discuss.**
-
- **Draw the Sun and the Earth in summer and winter.**
 - **Draw the same satellite orbiting the Earth.**
 - **Change your picture with your neighbour. Discuss.**



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Additional notes

Geostationary orbit (GEO)

24-hour orbit

- satellite always on the same longitude
- good for, e.g., telecommunications and broadcasting
- good global coverage

Orbit in the equatorial plane

- poor coverage at high latitudes

GEO is becoming crowded

Periapsis & Apoapsis

The point on the orbit nearest the occupied focus is called the periapsis

The point farthest from the occupied focus is called the apoapsis.

If the central body is the Earth, these points are also referred to as the perigee and apogee, respectively. Similarly, if the central body is the sun, the names change to perihelion and aphelion.

Some links, used in lecture

Gravity

<https://www.youtube.com/watch?v=MTY1Kje0yLg>

<https://www.youtube.com/watch?v=a3OQ7ek7t68>

Orbit elements made easy

<http://www.amsat.org/amsat/keps/kepmodel.html>

Satellite orbits

<https://www.youtube.com/watch?v=4K5FyNbV0nA>

<https://www.youtube.com/watch?v=Hcm7oQwpZfg>

<https://www.youtube.com/watch?v=uZc0YJjyWGM>

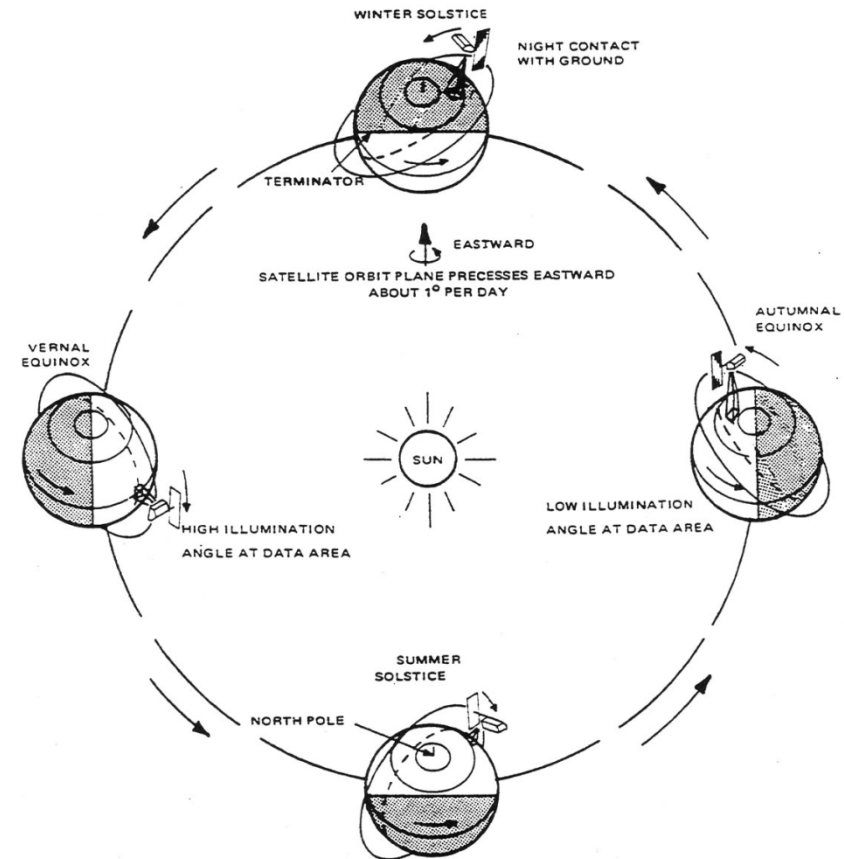


Sun synchronous orbits

Combined effect of the precession of the orbit and the motion around the Sun

Same geographic location (e.g. equator) passed always at the same local time

- useful in many Earth-observation applications



Useful relations

- Angular Momentum

$$\vec{h} = \vec{r} \times \vec{v}$$

- Node Line Vector

$$\vec{N} = \hat{K} \times \vec{h}$$

\vec{N} is line of nodes;

\hat{K} is reference axis

- Argument of the periapsis

$$\omega = \cos^{-1} \frac{\vec{N} \cdot \vec{e}}{|\vec{N}| |e|}$$

- Inclination

$$i = \cos^{-1} \frac{h_z}{h}$$

- True anomaly

$$\theta = \cos^{-1} \frac{\vec{e} \cdot \vec{r}}{|e| |r|}$$

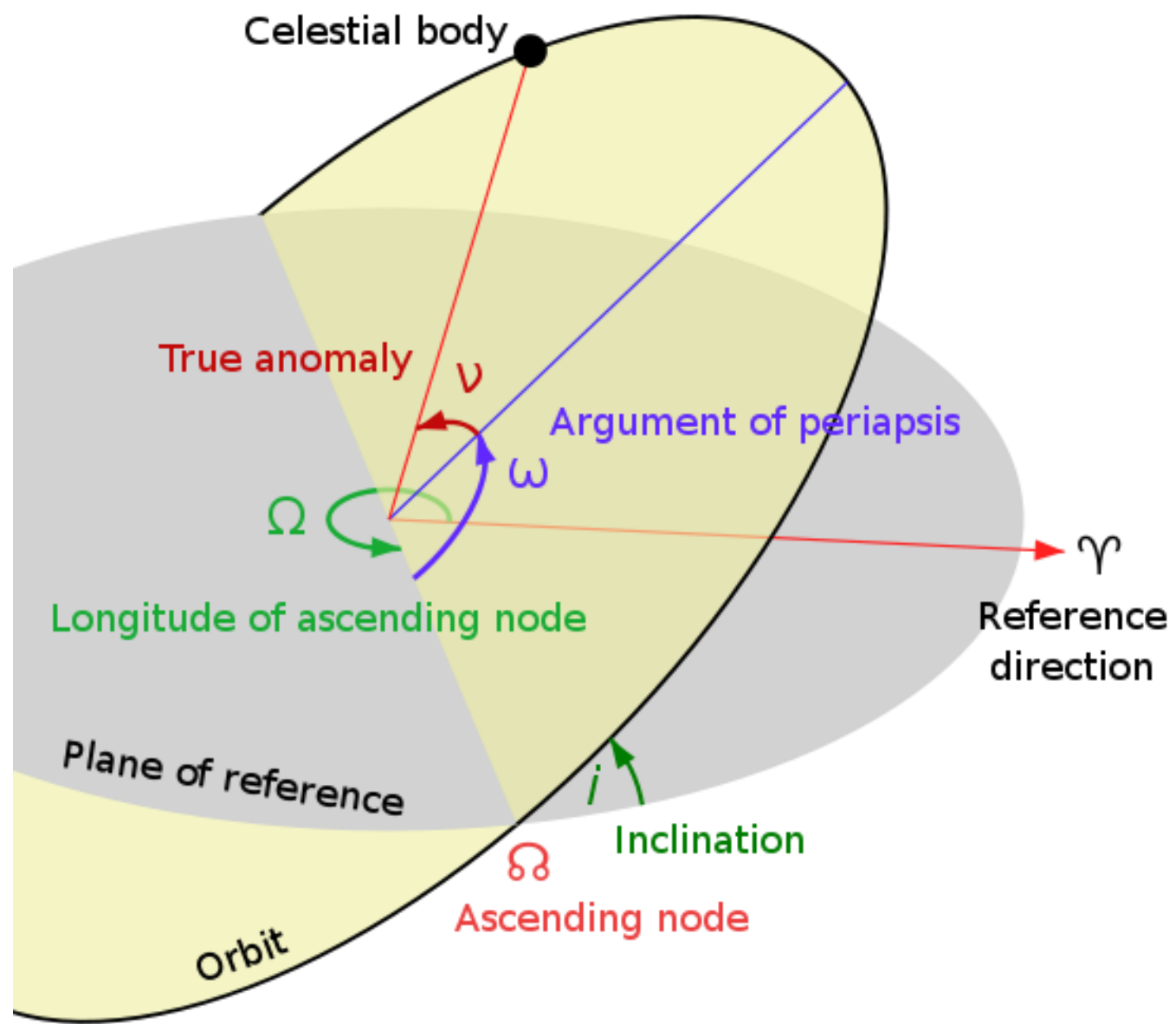
- Right Ascension of the ascending node

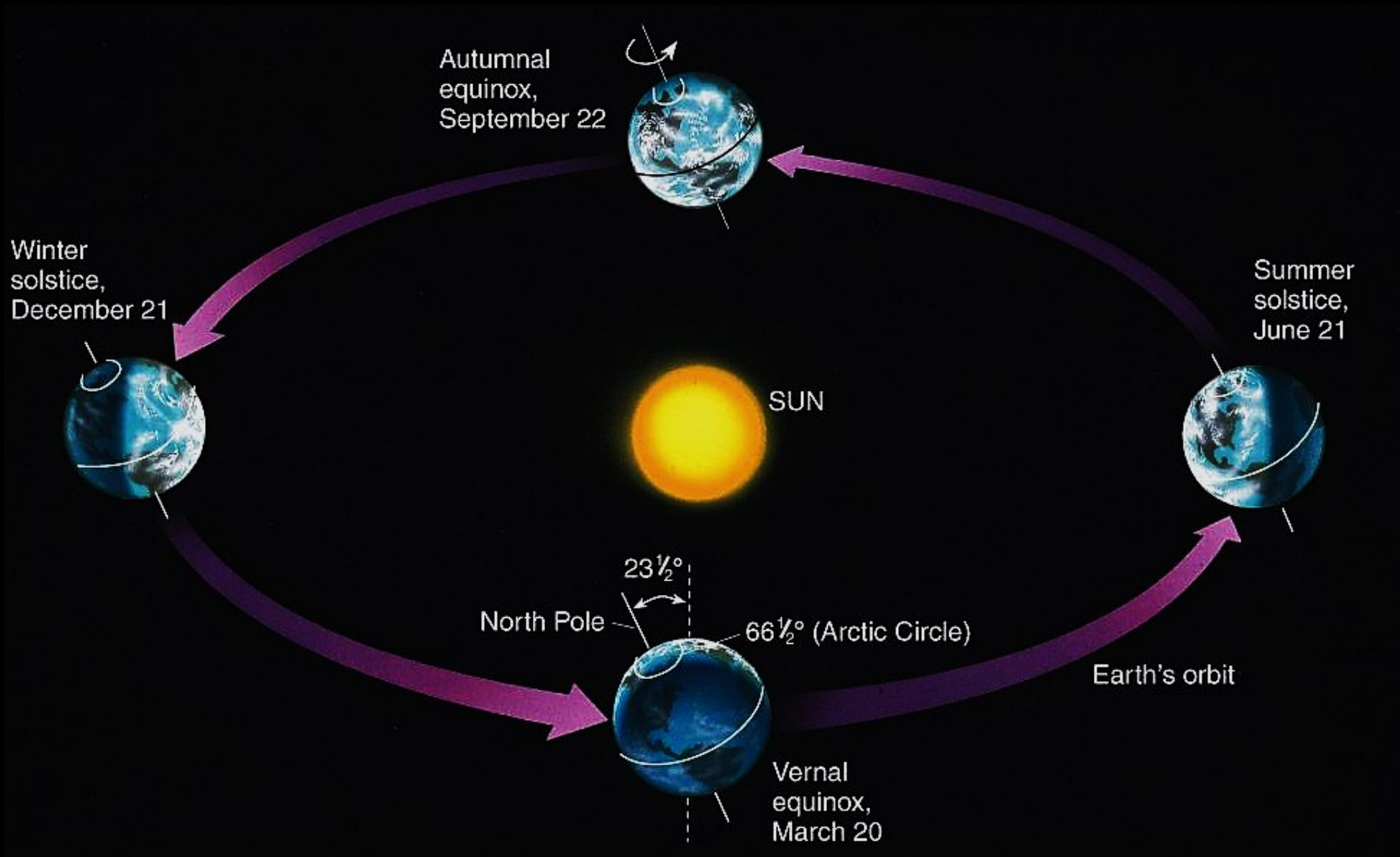
$$\Omega = \cos^{-1} \frac{\vec{I} \cdot \vec{N}}{|\vec{N}|}$$

\vec{I} is reference axis

- Eccentricity

$$\vec{e} = \frac{1}{\mu} \left[\left(v^2 - \frac{\mu}{r} \right) \vec{r} - r v_r \vec{v} \right];$$
$$v_r = \frac{\vec{r} \cdot \vec{v}}{r}$$





Line of Nodes

Ascending Node

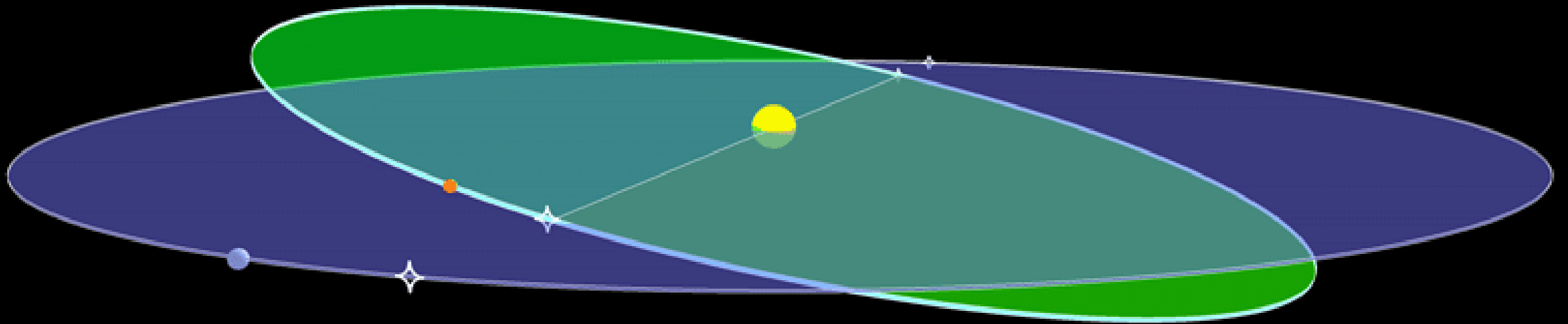
- For geocentric and heliocentric orbits, the **ascending node** (or **north node**) is where the orbiting object moves north through the plane of reference.

Descending node

- For geocentric and heliocentric orbits, the **descending node** (or **south node**) is where it moves south through the plane of reference.

In the case of objects outside the Solar System, the ascending node is the node where the orbiting secondary passes away from the observer, and the descending node is the node where it moves towards the observer.

Line of Nodes



[Source: Wikipedia]

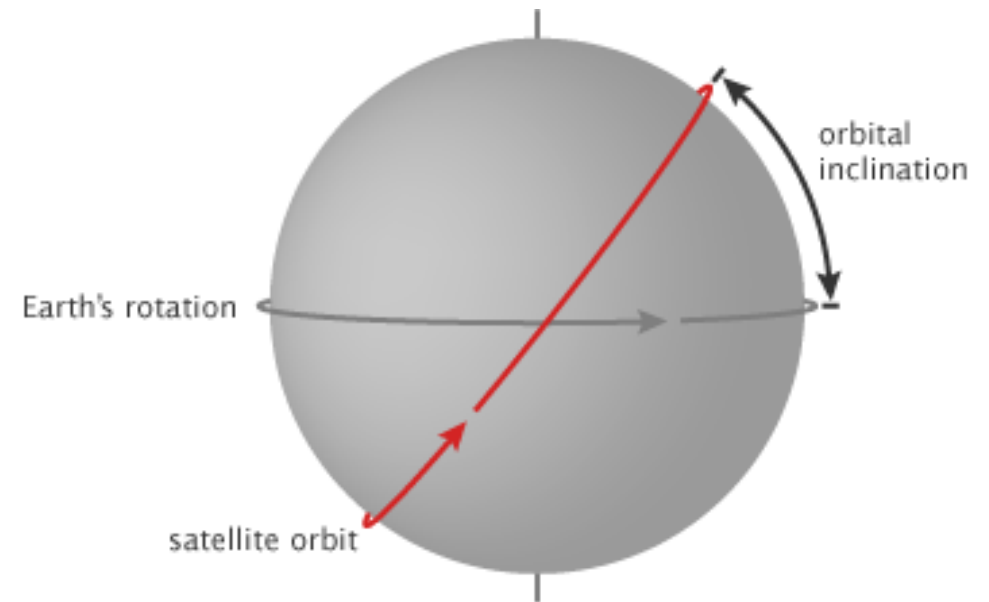
Inclination Angle

Prograde Orbit

- A spacecraft or other body moves in the same direction as the planet rotates.
- An orbit inclination between 0 and 90 degrees will generate a prograde orbit.

Retrograde Orbit

- A spacecraft or other body moves in the opposite direction to the planet's rotation
- An orbit inclination between 90 and 180 degrees is generate a retrograde orbit.



[Source: NASA]

- An object with an inclination of 90 degrees is said to be neither prograde nor retrograde, but it is instead called a polar orbit. (since it does not have an east-west directional component)