

# Quantum Circuits

## Week 6

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### 1 Quantum algorithms

Our final topic in this course is a brief overview of the most promising algorithms that demonstrate a quantum advantage. Our treatment is far from complete, since the field of quantum algorithms has matured enough to be considered an independent research discipline.

### 2 Quantum Fourier transform (QFT)

The classical discrete Fourier transform is an operation that represents complex-valued vectors as a linear sum of the Fourier basis

$$y_m = \frac{1}{\sqrt{N}} \sum_{n=0}^{N-1} x_n e^{2\pi i m n / N}$$

where  $x = (x_0, \dots, x_{N-1})^\top, y = (y_0, \dots, y_{N-1})^\top \in \mathbb{C}$ , and  $e_m(n) = \frac{1}{\sqrt{N}} e^{2\pi i m n / N}$ ,  $n, m = 0, 1, 2, \dots, N-1$  are the Fourier basis.

Despite its similar mathematical representation, the QFT doesn't compute any complex coefficients as the DFT. Its primary role is to create an n-qubit superposition state via a basis transformation. Without this resource, none of the foretold quantum advantages can be manifested. QFT transforms basis states as

$$\mathcal{F} |n\rangle = \frac{1}{\sqrt{N}} \sum_{m=0}^{N-1} e^{2\pi i m n / N} |m\rangle$$

While for any arbitrary state  $|\Psi\rangle = \sum_{n=0}^{N-1} \alpha_n |n\rangle$

$$\begin{aligned}
\mathcal{F}|\Psi\rangle &= |\Phi\rangle \\
&= \sum_{m=0}^{N-1} \alpha_m |m\rangle \\
&= \frac{1}{\sqrt{N}} \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} \alpha_n e^{2\pi i mn/N} |m\rangle
\end{aligned}$$

where  $\alpha_m = \frac{1}{\sqrt{N}} \sum_{n=0}^{N-1} \alpha_n e^{2\pi i mn/N} = \sum_{n=0}^{N-1} \alpha_n (e^{2\pi i/N})^{mn}$ , such that  $W_N = \frac{1}{\sqrt{N}} e^{2\pi i/N}$ . Then by defining  $\Omega_N$  as the matrix

$$\Omega_N = \frac{1}{\sqrt{N}} W_N^{mn}, 0 \leq m, n \leq N-1,$$

we can write the matrix representation of the Fourier operator (ex).

By virtue of the Fourier *inverse theorem*, it is straight forward to show that the QFT operator is *unitary* (ex).

Letting  $N = 2^k$ , we can write the QFT of an n-qubit state as<sup>1</sup>

$$\frac{1}{\sqrt{2^k}} \bigotimes_{l=1}^k (|0\rangle + e^{i2\pi m/2^l} |1\rangle)$$

## 2.1 Example

Suppose that  $n = 3$ , in binary representation this can be written as a the 2-bit string “11”. Quantum mechanically we define this as a two-qubit tensor product state  $|1\rangle \otimes |1\rangle \equiv |11\rangle$ . The previous state can be transformed to the Fourier domain as

$$\begin{aligned}
\mathcal{F} &= \frac{1}{\sqrt{2^2}} \sum_{m=0}^3 e^{2\pi i(3m)/2^2} |m\rangle \\
&= \frac{1}{2} [ |00\rangle + e^{i3\pi/2} |01\rangle + e^{i3\pi} |10\rangle + e^{i9\pi/2} |11\rangle ]
\end{aligned}$$

Recalling our discussion in the past week, we can verify that the transformed vector indeed belongs to the 4-dimensional Fourier space.

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<sup>1</sup>Full derivation can be found in the Qiskit textbook.

## 2.2 Quantum phase estimation

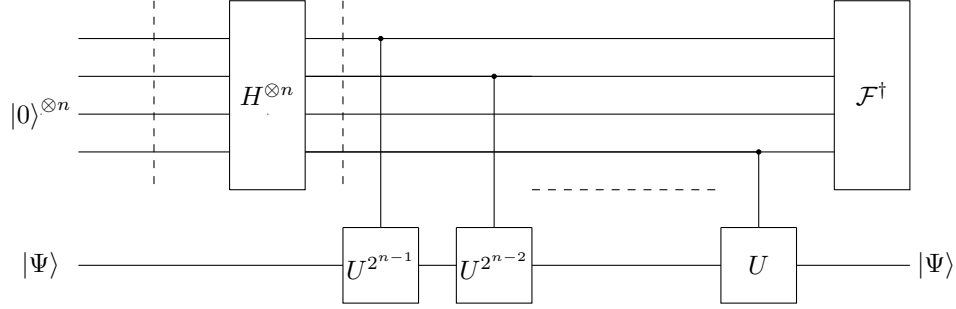


Figure 1: A sketch of the quantum phase estimation circuit.

A practical application that relies on the QFT is the problem of estimating the phase of an eigen-state of a unitary operator. A unitary operator  $U$  can be implemented as

$$U = e^{-i\gamma A} = e^{-i\gamma \sum_n a|a\rangle\langle a|}$$

where eigenstates of  $U$  have eigenvalues of the form  $e^{-i\theta}$ , such that  $\theta = a\gamma$ .

Projective measurements on the eigenstates wipe out the overall global phase of the state, that's why we introduce the phase estimation circuit to retrieve this information<sup>2</sup>

## 2.3 Procedure

- (1) We first prepare a huge superposition resource by applying the n-tensor product Hadamard operator  $H^{\otimes n}$  to a stack of auxiliary qubits all prepared in the ground state  $|0\rangle^{\otimes n}$ .
- (2) When  $U^{2^n}$  is applied to  $|\Psi\rangle$  it gives  $U^{2^n} |\Psi\rangle = e^{i2^n\theta} |\Psi\rangle$
- (2) At this stage the overall state of the system is written as

$$\left[ \frac{1}{(\sqrt{2})^n} (|0\rangle + |1\rangle)^{\otimes n} \right] \otimes |\Psi\rangle$$

- (3) Then applying  $n$  controlled-U operations gives

$$\left[ \frac{1}{(\sqrt{2})^n} (|0\rangle + e^{i\theta 2^{n-1}} |1\rangle) \otimes (|0\rangle + e^{i\theta 2^{n-2}} |1\rangle) \dots \otimes (|0\rangle + e^{i\theta} |1\rangle) \right] \otimes |\Psi\rangle$$

- (4) It is then straight forward to put the above expression in the QFT form, then applying  $\mathcal{F}^\dagger$  gives an expression that has a delta-peak at  $2^n\theta$ , thus the final state is projected to  $|2^n\theta\rangle |\Psi\rangle$ .

<sup>2</sup>For more details please consult “Quantum Computation and Quantum Information”-Nielsen & Chuang.

### 3 Grover’s algorithm

On the average, identifying an element in an unsorted database of  $N$  distinct elements requires  $N/2$  searches. Classically, the time taken by this operation is of the order of  $N$ . However, Grover showed in his protocol that by exploiting quantum parallelism, we can achieve a quadratic speed by performing the same task in  $\sqrt{N}$  steps. In our exercise sheet we go through the entire protocol in details and derive this result ourselves. If implemented efficiently, Grover’s algorithm would revolutionize the internet industry.

### 4 The variational quantum eigensolver (VQE)

VQE is a hybrid quantum-classical protocol that utilizes the variational quantum method to find an *optimal* upper bound for the ground state energy level of a quantum system, a molecule for example.

$$\lambda_{\min} \leq \lambda_{\theta} \equiv \langle \Psi(\theta) | H | \Psi(\theta) \rangle$$

#### 4.1 Procedure

- (1) Write down the **Hamiltonian** of the system in terms of its Pauli decomposition. For example, the Hamiltonian of a Hydrogen atom can be represented by two qubits, and hence its Pauli components<sup>3</sup> are

$$H = \gamma_0 I + \gamma_1 Z_0 + \gamma_2 Z_1 + \gamma_3 Z_0 Z_1 + \gamma_4 Y_0 Y_1 + \gamma_5 X_0 X_1$$

- (2) Choose a wave-function ansatz  $|\Psi\rangle$ .
- (3) Encode a *parameter*  $\theta$  into the ansatz via a parameterized quantum circuit described by a unitary transformation  $U(\theta)$

$$U(\theta) |\Psi\rangle = |\Psi(\theta)\rangle$$

- (4) Calculate the expected value of the **Hamiltonian** in the state  $|\Psi(\theta)\rangle$
- (5) Optimize  $\theta$  via classical algorithms till you find the minimum eigenvalue.

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<sup>3</sup>“Scalable Quantum Simulation of Molecular Energies”