

ELEC-E4130

Lecture 15: Multilayer (Stratified Media) N layers, oblique incidence

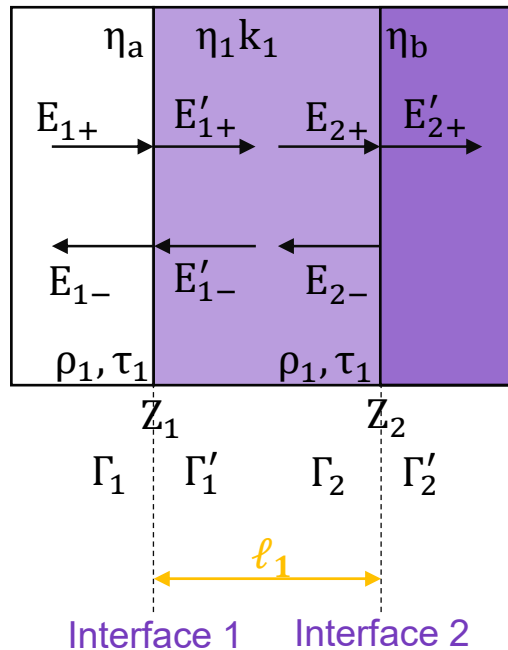


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Nov. 8, 2021

Reflection coefficient of slab



Reflection Coefficient Propagation

$$\Gamma_1 = \frac{\rho_1 + \Gamma_1'}{1 + \rho_1 \Gamma_1'}$$

$$\Gamma_1' = \Gamma_2 e^{-j2k_1 l_1}$$

$$\Gamma_1 = \frac{\rho_1 + \Gamma_2 e^{-j2k_1 l_1}}{1 + \rho_1 \Gamma_2 e^{-j2k_1 l_1}}$$

$$\Gamma_2 = \rho_2$$

$$\Gamma_1 = \frac{\rho_1 + \rho_2 e^{-j2k_1 l_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 l_1}}$$

Matching/propagation Matrices

$$\begin{bmatrix} E_{1+} \\ E_{1-} \end{bmatrix} = \frac{1}{\tau_1} \cdot \begin{bmatrix} 1 & \rho_1 \\ \rho_1 & 1 \end{bmatrix} \cdot \begin{bmatrix} e^{jk_1 l_1} & 0 \\ 0 & e^{-jk_1 l_1} \end{bmatrix} \cdot \frac{1}{\tau_2} \cdot \begin{bmatrix} 1 & \rho_2 \\ \rho_2 & 1 \end{bmatrix} \cdot \begin{bmatrix} E_{2+} \\ 0 \end{bmatrix}$$

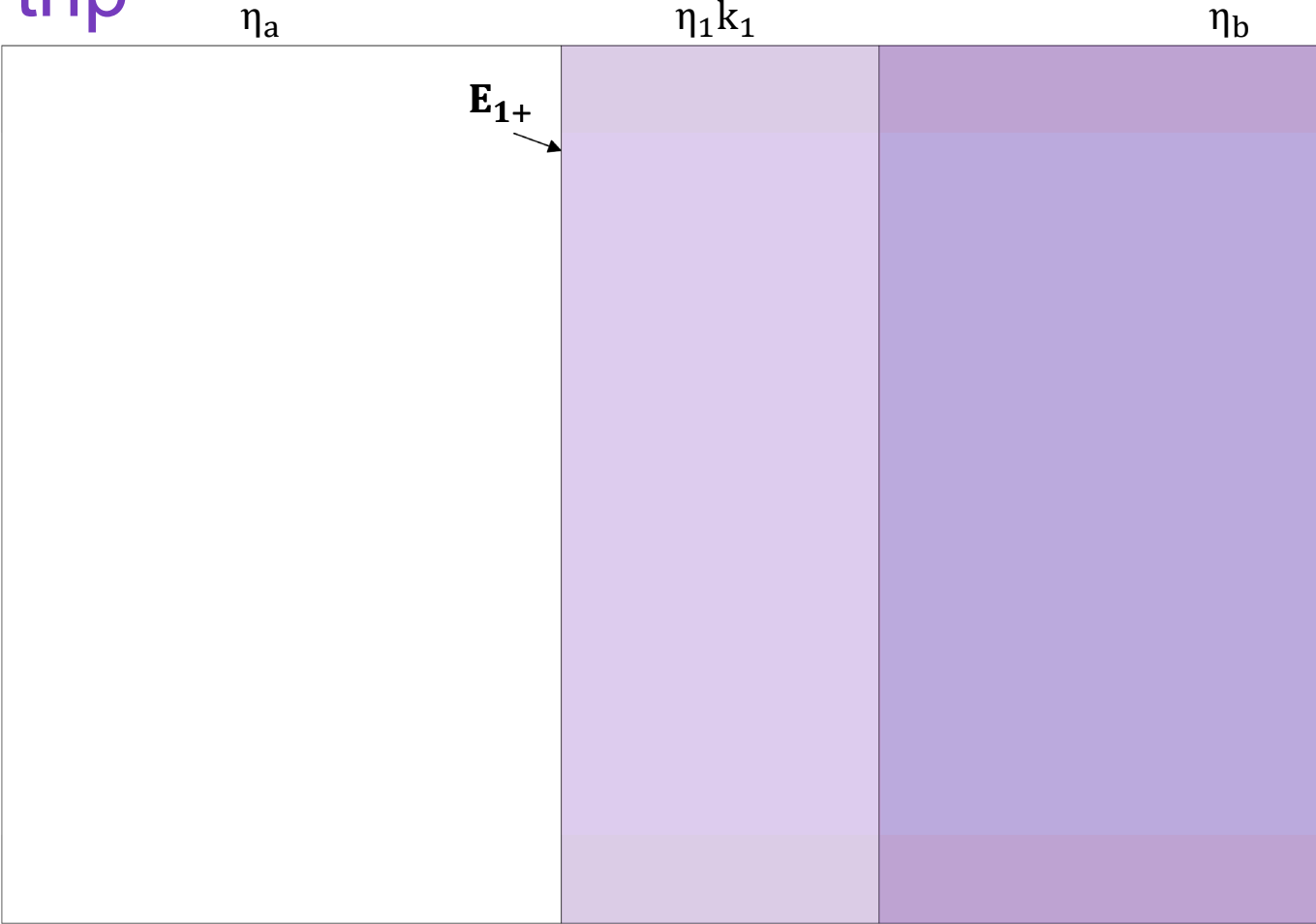
$$E_{1+} = \frac{e^{jk_1 l_1}}{\tau_1 \tau_2} (1 + \rho_1 \rho_2 e^{-j2k_1 l_1}) E_{2+}$$

$$E_{1-} = \frac{e^{jk_1 l_1}}{\tau_1 \tau_2} (\rho_1 + \rho_2 e^{-j2k_1 l_1}) E_{2+}$$

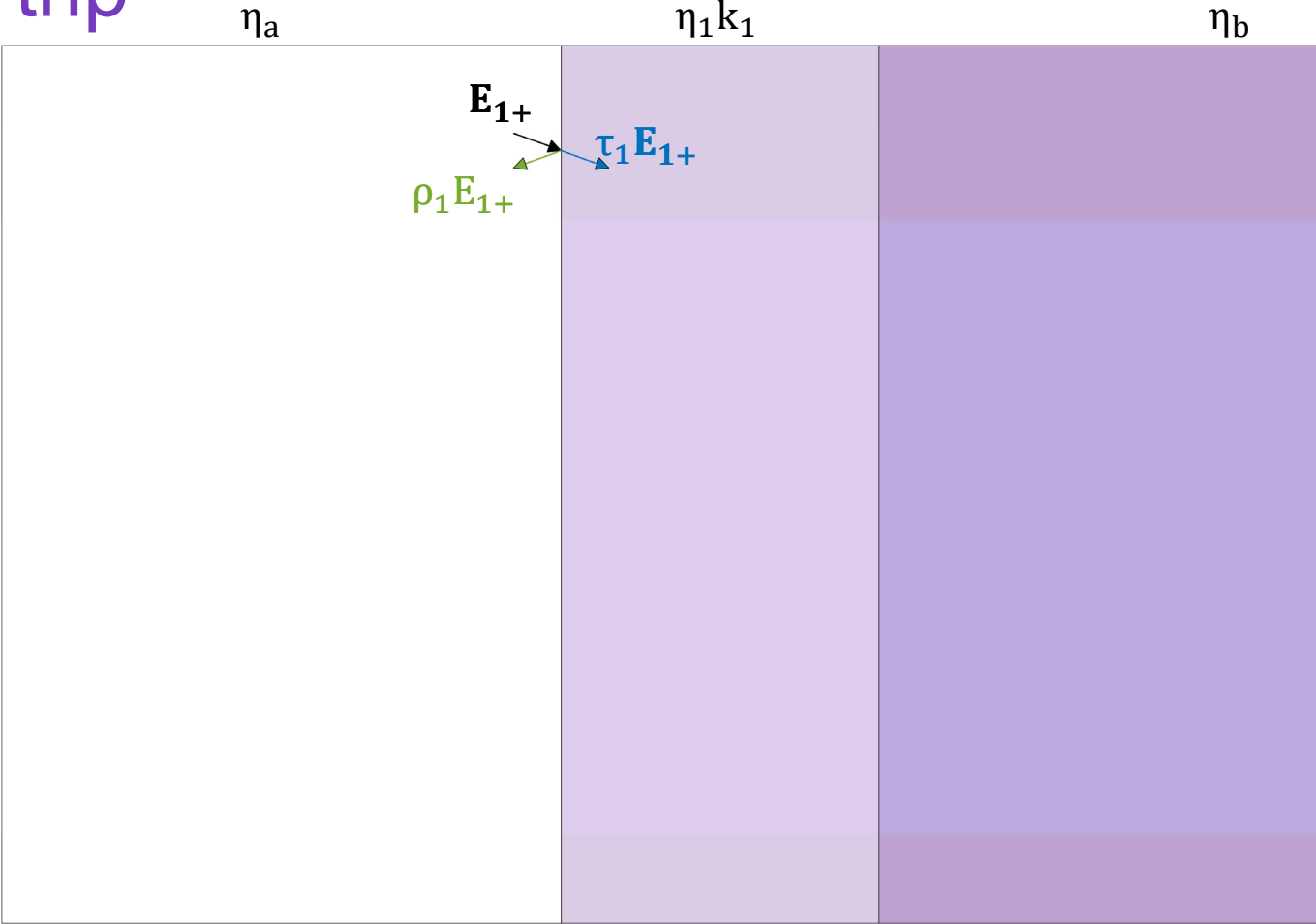
$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \frac{\rho_1 + \rho_2 e^{-j2k_1 l_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 l_1}}$$

Single dielectric slab via geometric series

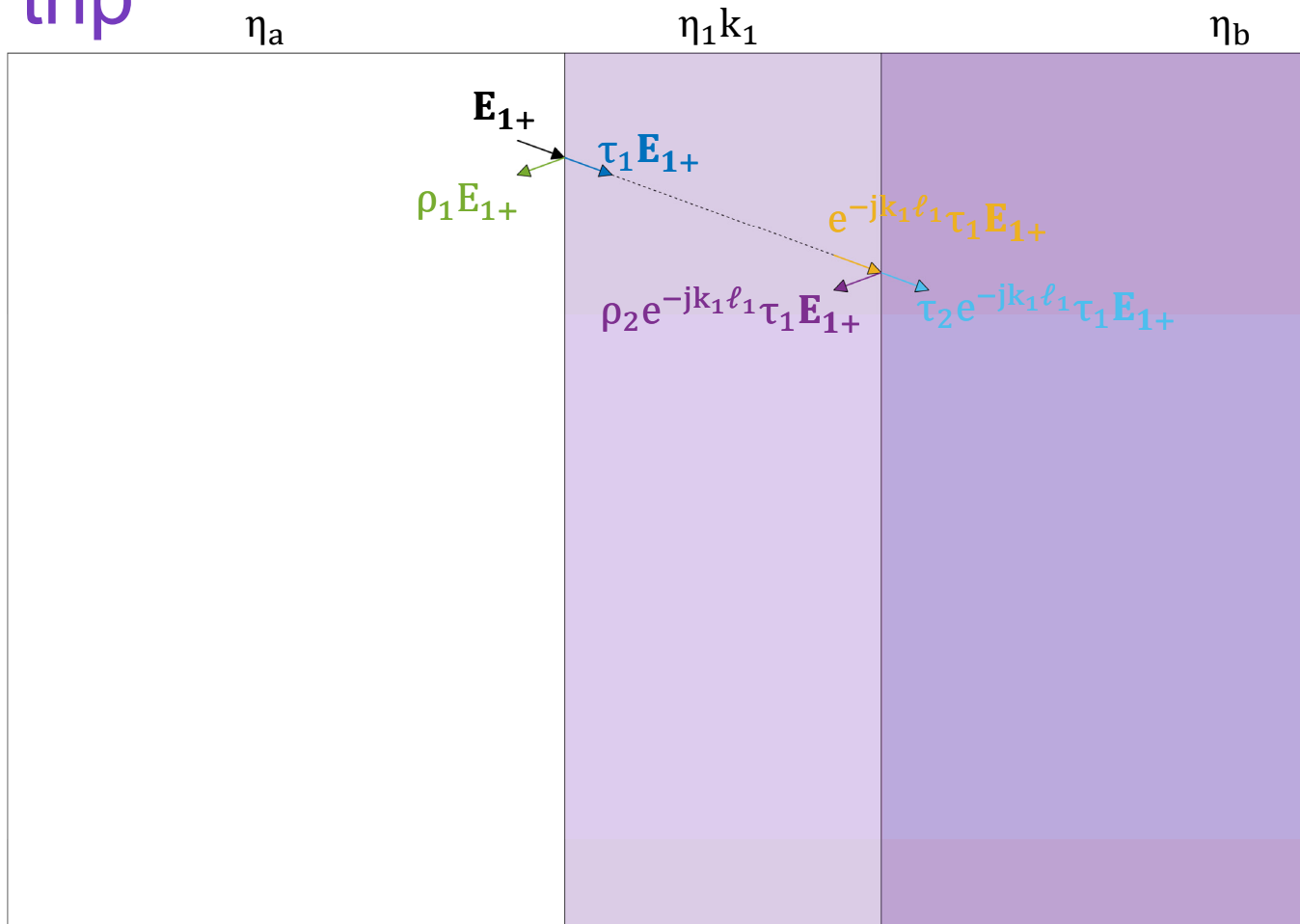
Round trip



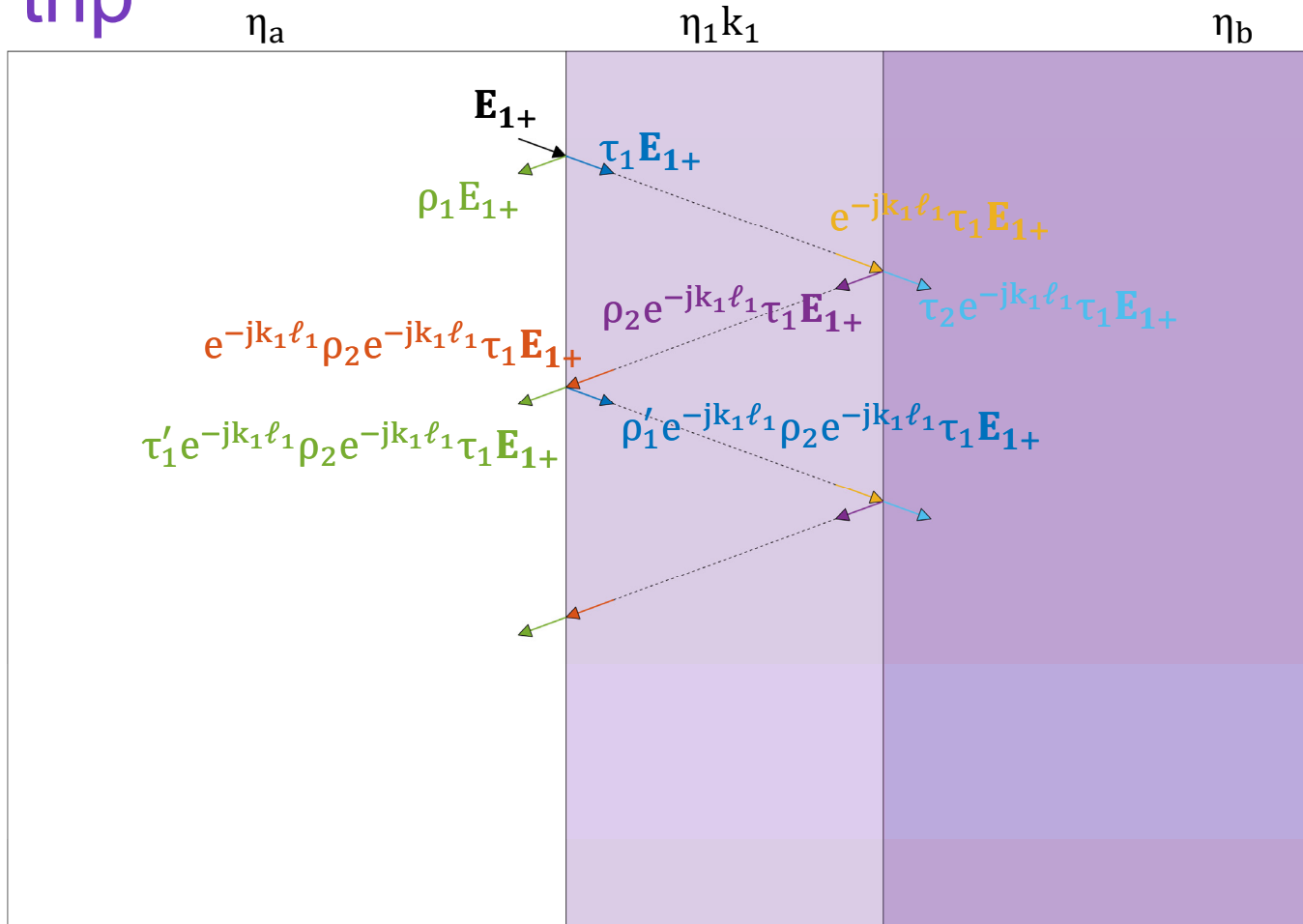
Round trip



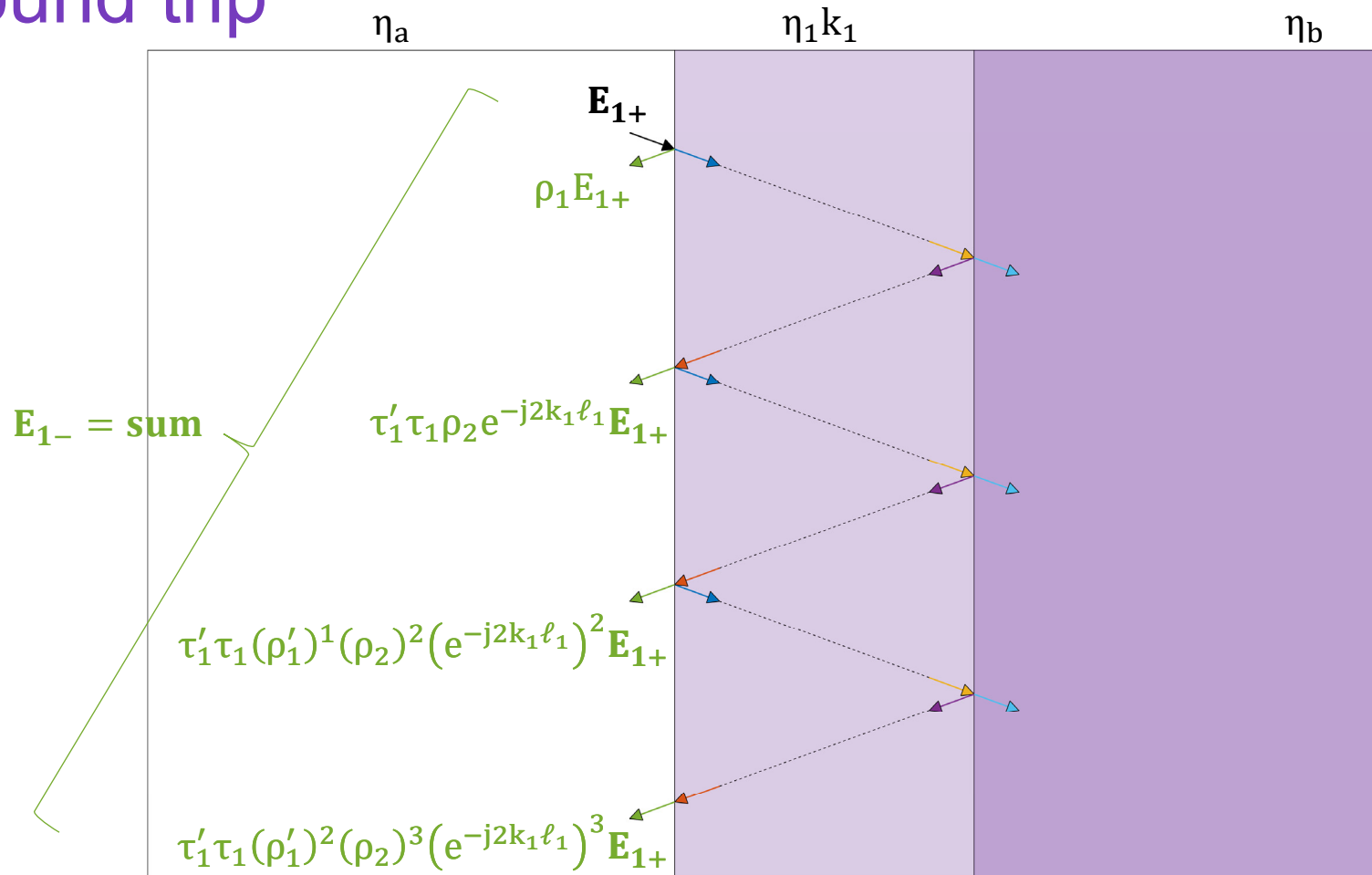
Round trip



Round trip



Round trip



Round trip

$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \rho_1 + \sum_{N=1}^{\infty} \tau'_1 \tau_1 (\rho'_1)^{N-1} (\rho_2)^N (e^{-j2k_1 \ell_1})^N$$

$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \sum_{N=1}^{\infty} (\rho'_1)^N (\rho_2)^N (e^{-j2k_1 \ell_1})^N$$

$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \left[-1 + \sum_{N=0}^{\infty} (\rho'_1 \rho_2 e^{-j2k_1 \ell_1})^N \right]$$

Geometric Series

$$\sum_{N=0}^{\infty} (x)^N = \frac{1}{1-x} \quad \text{when } |x| < 1$$

$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \left[-1 + \frac{1}{1 - \rho'_1 \rho_2 e^{-j2k_1 \ell_1}} \right]$$

Round trip

$$\Gamma_1 = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \left[-1 + \frac{1}{1 - \rho'_1 \rho_2 e^{-j2k_1 \ell_1}} \right]$$

$$\Gamma_1 = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \left[\frac{-1 + \rho'_1 \rho_2 e^{-j2k_1 \ell_1} + 1}{1 - \rho'_1 \rho_2 e^{-j2k_1 \ell_1}} \right]$$

Matching Matrices from Lecture 13

$$\Gamma_1 = \rho_1 + \frac{\tau'_1 \tau_1}{\rho'_1} \left[\frac{\rho'_1 \rho_2 e^{-j2k_1 \ell_1}}{1 - \rho'_1 \rho_2 e^{-j2k_1 \ell_1}} \right] \longrightarrow \begin{aligned} \tau'_1 \tau_1 &= 1 - \rho_1^2 \\ \rho_1 &= -\rho'_1 \end{aligned}$$

$$\Gamma_1 = \rho_1 + (1 - \rho_1^2) \left[\frac{\rho_2 e^{-j2k_1 \ell_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}} \right]$$

Round trip

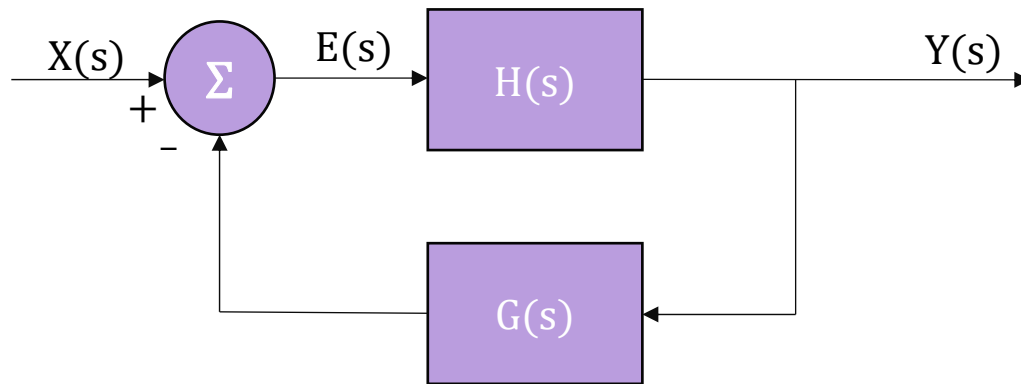
$$\Gamma_1 = \rho_1 + (1 - \rho_1^2) \left[\frac{\rho_2 e^{-j2k_1 \ell_1}}{1 - \rho_1 \rho_2 e^{-j2k_1 \ell_1}} \right]$$

$$\Gamma_1 = \frac{\rho_1(1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}) + (1 - \rho_1^2) \rho_2 e^{-j2k_1 \ell_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}}$$

$$\Gamma_1 = \frac{\rho_1 + \rho_2 e^{-j2k_1 \ell_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}}$$

- Reflection from a single slab can be computed from the infinite sum of plane waves bouncing around the slab
- Amplitude of waves decays geometrically
- Interference at surface dictated by
 - $k_1 \ell_1$
 - Phase contributions of ρ_1, ρ_2

Comparisons to feedback control theory



$$\begin{aligned} E(s) &= X(s) - G(s)Y(s) \\ Y(s) &= H(s)E(s) \\ Y(s) &= H(s)(X(s) - G(s)Y(s)) \end{aligned}$$

$$\frac{Y(s)}{X(s)} = \frac{H(s)}{1 + G(s)H(s)}$$

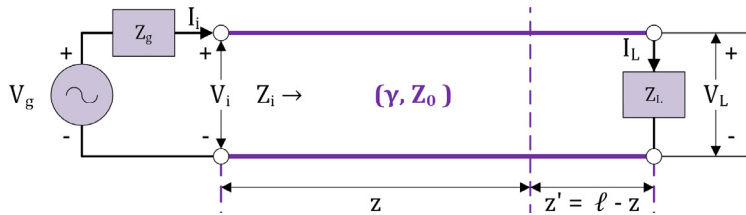
From slide 11

$$\Gamma_1 = \rho_1 + (1 - \rho_1^2) \left[\frac{\rho_2 e^{-j2k_1 \ell_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}} \right] \rightarrow \begin{aligned} H(s) &\sim \rho_2 e^{-j2k_1 \ell_1} \\ G(s) &\sim \rho_1 \end{aligned}$$

Feedback term

- Can you design the layers to act like a filter? **YES!**
- Can you use filter/linear system theory to design the layers? **YES!**

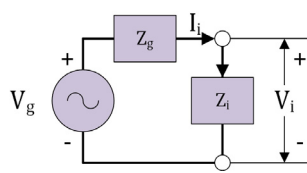
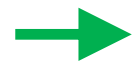
Comparisons to transmission line theory



Reflection Coefficient of dielectric slab

$$\Gamma_1 = \frac{\rho_1 + \rho_2 e^{-j2k_1 l_1}}{1 + \rho_1 \rho_2 e^{-j2k_1 l_1}}$$

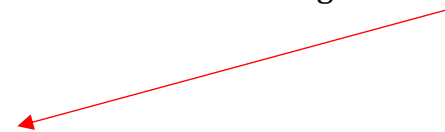
➤ Physical thickness and propagation constant dependent



Voltage and current phasors, arb. Line

$$V(z') = V_g \frac{Z_0}{Z_0 + Z_g} e^{\gamma(z'-\ell)} \frac{1 + \Gamma e^{-2\gamma z'}}{1 - \Gamma_g \Gamma e^{-2\gamma \ell}}$$

$$I(z') = V_g \frac{1}{Z_0 + Z_g} e^{\gamma(z'-\ell)} \frac{1 + \Gamma e^{-2\gamma z'}}{1 - \Gamma_g \Gamma e^{-2\gamma \ell}}$$



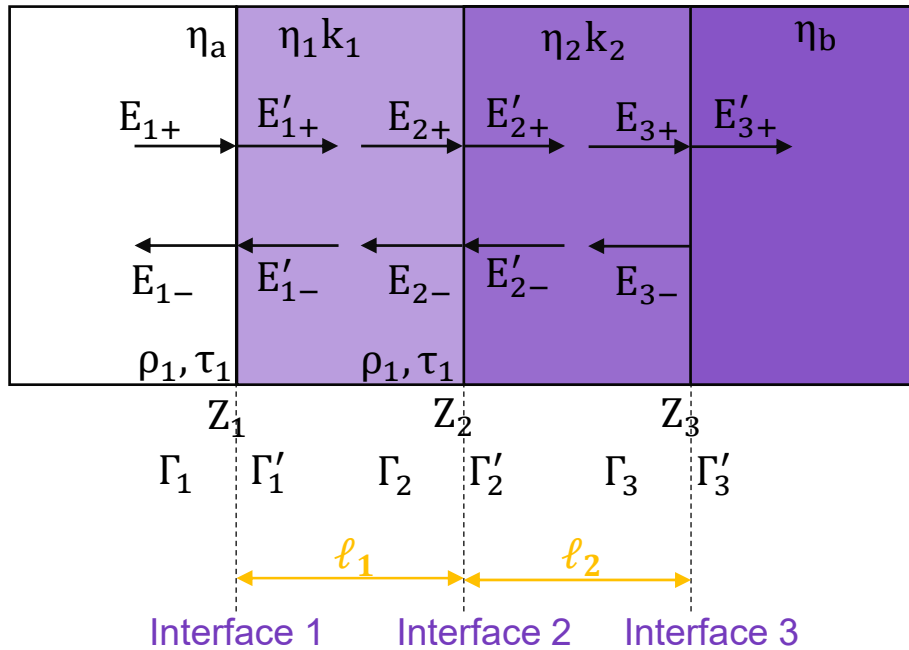
$$Z_i = Z_0 \frac{Z_L + Z_0 \tanh(\gamma \ell)}{Z_0 + Z_L \tanh(\gamma \ell)}$$



Normal Incidence, M layers

2...M layers

Lets let $E'_{2-} \neq 0$



$$\begin{bmatrix} E_{1+} \\ E_{1-} \end{bmatrix} = \frac{1}{\tau_1} \begin{bmatrix} 1 & \rho_1 \\ \rho_1 & 1 \end{bmatrix} \begin{bmatrix} e^{jk_1 \ell_1} & 0 \\ 0 & e^{-jk_1 \ell_1} \end{bmatrix} \frac{1}{\tau_2} \begin{bmatrix} 1 & \rho_2 \\ \rho_2 & 1 \end{bmatrix} \begin{bmatrix} E'_{2+} \\ E'_{2-} \end{bmatrix}$$

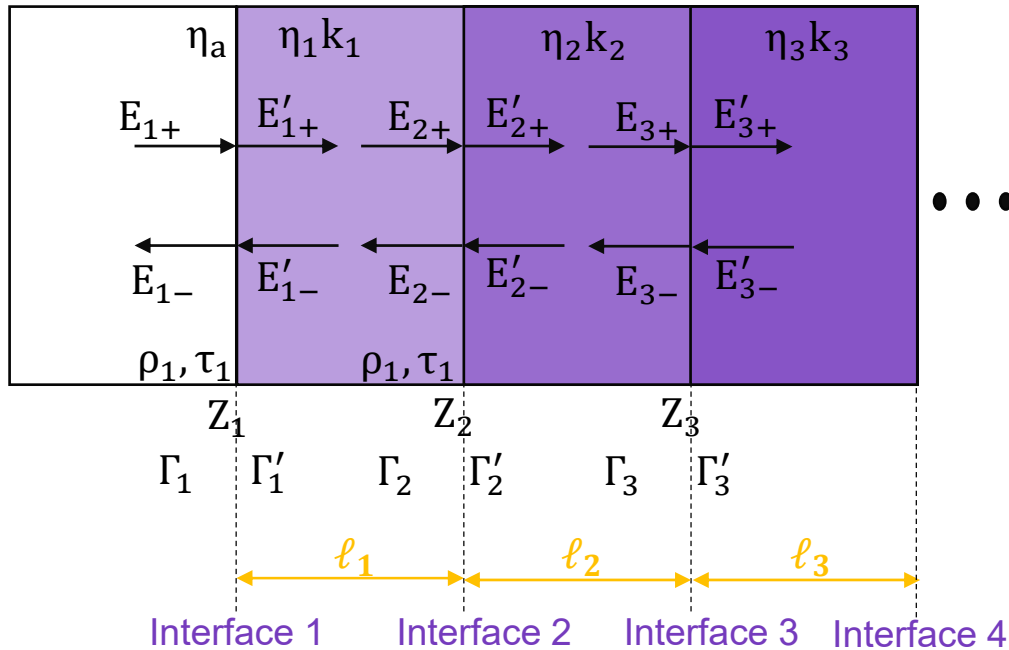
$$E_{1+} = \frac{e^{jk_1 \ell_1}}{\tau_1 \tau_2} (1 + \rho_1 \rho_2 e^{-j2k_1 \ell_1}) E'_{2+}$$

$$E_{1-} = \frac{e^{jk_1 \ell_1}}{\tau_1 \tau_2} (\rho_1 + \rho_2 e^{-j2k_1 \ell_1}) E'_{2+}$$



$$\Gamma_1 = \frac{E_{1-}}{E_{1+}} = \frac{\rho_1 + \frac{E_{2-}}{E_{2+}} e^{-j2k_1 \ell_1}}{1 + \rho_1 \frac{E_{2-}}{E_{2+}} e^{-j2k_1 \ell_1}} = \frac{\rho_1 + \Gamma_2 e^{-j2k_1 \ell_1}}{1 + \rho_1 \Gamma_2 e^{-j2k_1 \ell_1}}$$

2...M layers



M slabs
 M+1 interfaces
 $m = 1 \dots M+1$

$$\Gamma_m = \frac{\rho_m + \Gamma_{m+1} e^{-j2k_m \ell_m}}{1 + \rho_m \Gamma_{m+1} e^{-j2k_m \ell_m}}$$

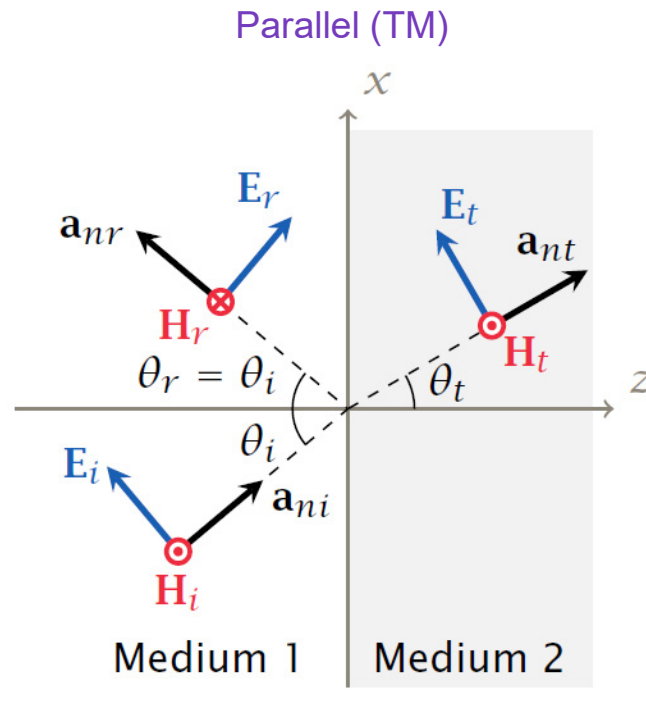
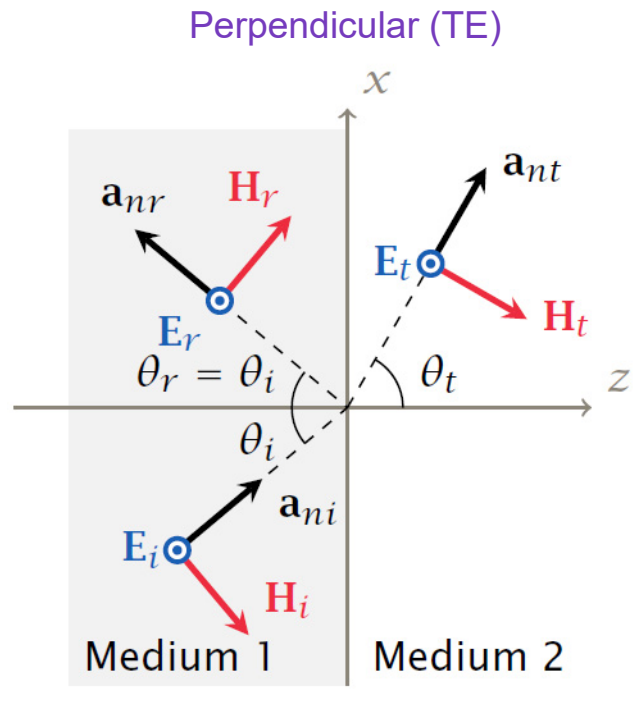
$\Gamma_1 =$ goal coefficient

$$\Gamma_{M+1} = \rho_{m+1}$$

- Start from the second half space and work your way back to the first
- Second half space only has a forward propagating wave thus ending the recursion

Oblique incidence (Dielectrics)

Review of oblique angle and polarization



Wave impedance

$$\eta_T = \begin{cases} \eta \cos(\theta) & \parallel, \text{TM} \\ \frac{\eta}{\cos(\theta)} & \perp, \text{TE} \end{cases}$$

Refractive index

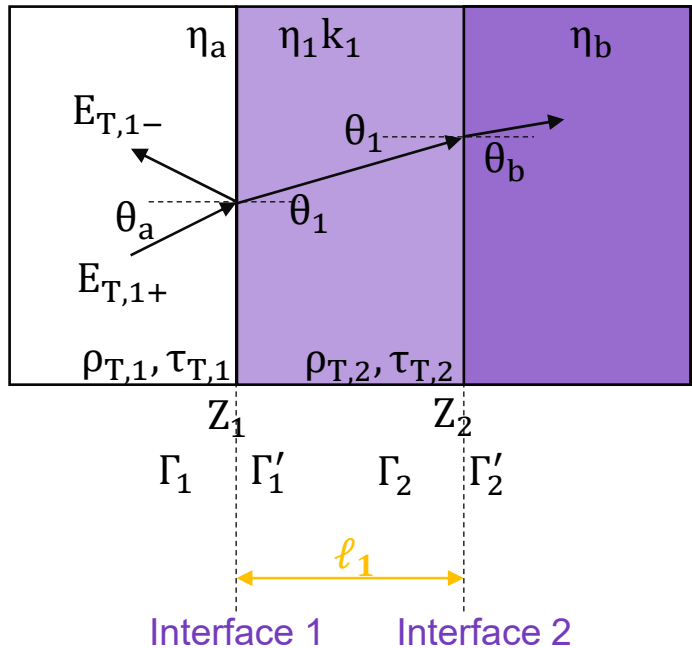
$$n_T = \begin{cases} \frac{n}{\cos(\theta)} & \parallel, \text{TM} \\ n \cos(\theta) & \perp, \text{TE} \end{cases}$$

Snell's law

$$\sin(\theta_i) = \sin(\theta_t)$$

- Obtained by continuity of transverse component
- Results matching coefficients indicated by the **T** subscript

Oblique Incidence: Transverse components



$$\eta_T = \begin{cases} \eta \cos(\theta) & \parallel, \text{TM} \\ \frac{\eta}{\cos(\theta)} & \perp, \text{TE} \end{cases} \quad \rightarrow \quad \eta'_T = \begin{cases} \eta' \cos(\theta') & \parallel, \text{TM} \\ \frac{\eta'}{\cos(\theta')} & \perp, \text{TE} \end{cases}$$

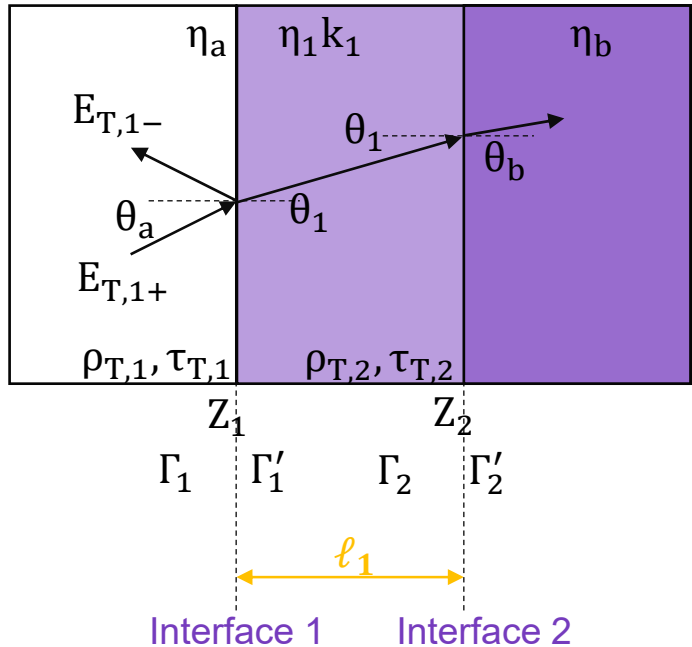
$$n_T = \begin{cases} \frac{n}{\cos(\theta)} & \parallel, \text{TM} \\ n \cos(\theta) & \perp, \text{TE} \end{cases} \quad \rightarrow \quad n'_T = \begin{cases} \frac{n'}{\cos(\theta')} & \parallel, \text{TM} \\ n' \cos(\theta') & \perp, \text{TE} \end{cases}$$

Medium \rightarrow Medium'



Medium 1 \rightarrow Medium 2

Oblique Incidence: Transverse components



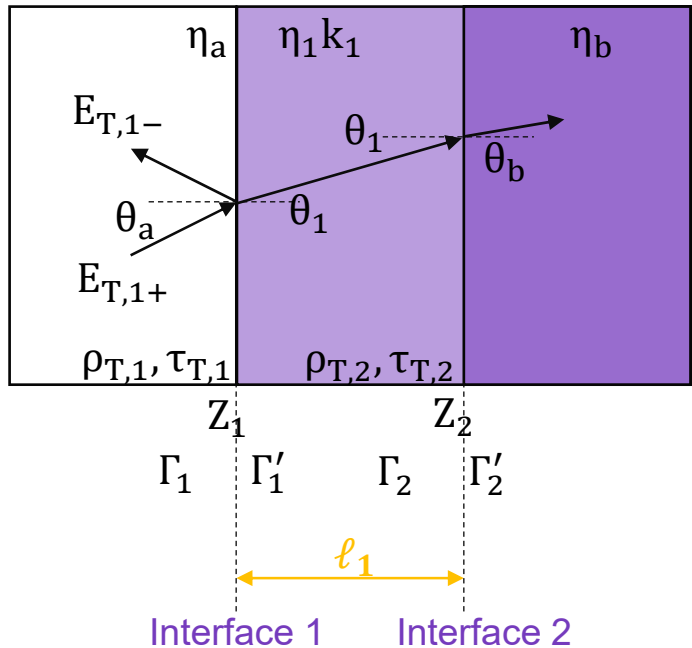
Fresnel's Coefficients

$$\rho_{T,1} = \frac{\eta_{T,1} - \eta_{T,a}}{\eta_1 + \eta_a} \quad \tau_{T,1} = 1 + \rho_{T,1}$$

$$\rho_{T,2} = \frac{\eta_{T,b} - \eta_{T,1}}{\eta_{T,b} + \eta_{T,1}} \quad \tau_{T,2} = 1 + \rho_{T,2}$$

- Same form as those computed with normal incidence

Oblique Incidence: Transverse components



Transverse field matching

Matching the z-component of the k-vector

$$k \rightarrow k_z$$

$$k_z = k \cos(\theta)$$

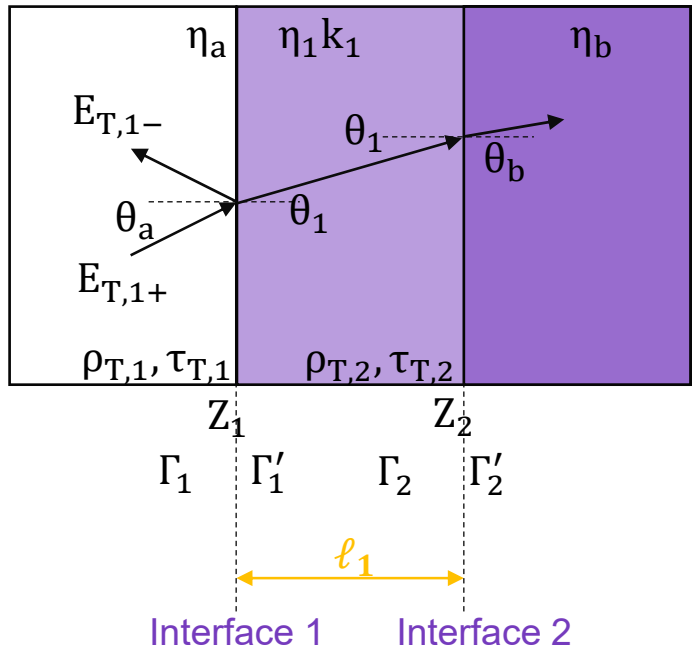
Propagation matrix

$$\begin{bmatrix} e^{jk_1 \ell_1} & 0 \\ 0 & e^{-jk_1 \ell_1} \end{bmatrix} \rightarrow \begin{bmatrix} e^{jk_{z,1} \ell_1} & 0 \\ 0 & e^{-jk_{z,1} \ell_1} \end{bmatrix}$$

$$\begin{bmatrix} e^{jk_{z,1} \ell_1} & 0 \\ 0 & e^{-jk_{z,1} \ell_1} \end{bmatrix} \rightarrow \begin{bmatrix} e^{jk_1 \ell_1 \cos(\theta)} & 0 \\ 0 & e^{-jk_1 \ell_1 \cos(\theta)} \end{bmatrix}$$

The transverse fields across the slab are related by $\cos(\theta)$

Oblique Incidence: Transverse components



$$\begin{bmatrix} E_{T,1+} \\ E_{T,1-} \end{bmatrix} = \frac{1}{\tau_1} \cdot \begin{bmatrix} 1 & \rho_{T,1} \\ \rho_{T,1} & 1 \end{bmatrix} \cdot \begin{bmatrix} e^{jk_1 \ell_1 \cos(\theta)} & 0 \\ 0 & e^{-jk_1 \ell_1 \cos(\theta)} \end{bmatrix} \\ \cdot \frac{1}{\tau_2} \cdot \begin{bmatrix} 1 & \rho_{T,2} \\ \rho_{T,2} & 1 \end{bmatrix} \cdot \begin{bmatrix} E'_{T,2+} \\ E'_{T,2-} \end{bmatrix}$$

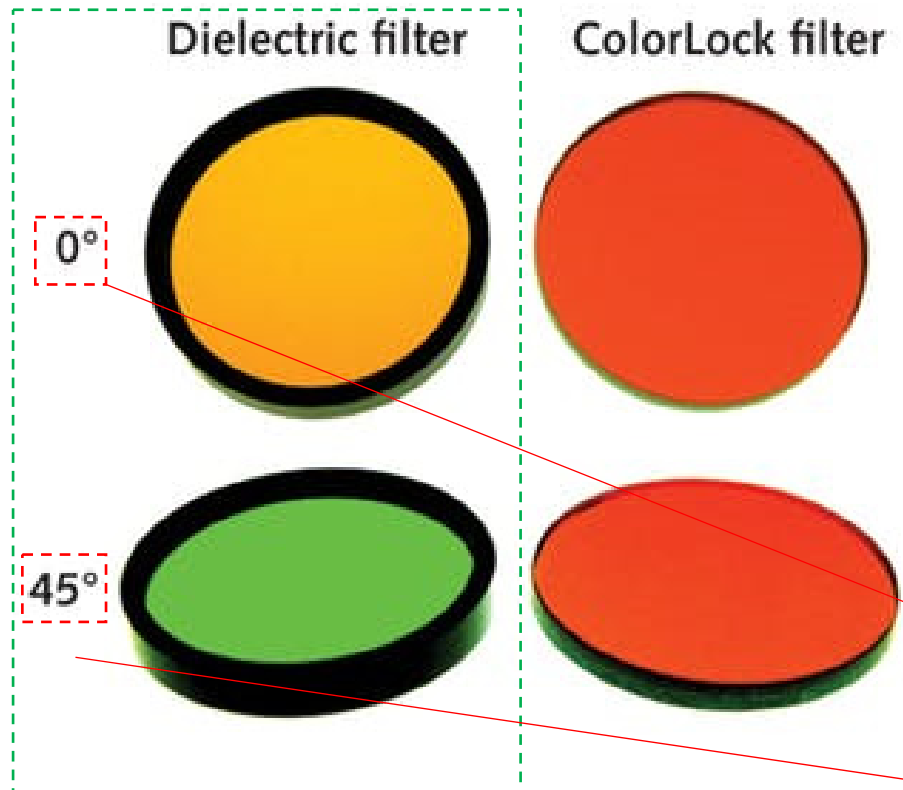
$$\Gamma_{T,m} = \frac{\rho_{T,m} + \Gamma_{T,m+1} e^{-j2\delta_m}}{1 + \rho_{T,m} \Gamma_{T,m+1} e^{-j2\delta_m}}$$

$$\delta_m = k_m \ell_m \cos(\theta)$$

$\Gamma_1 = \text{goal coefficient}$

$$\Gamma_{T,M+1} = \rho_{T,m+1}$$

Angle dependent behavior



- Transmission and reflection are defined by steady state constructive and destructive interference
- Interference controlled by optical path length within the dielectric
- Optical path is increased going from normal to oblique incidence angle

$$\Gamma_{T,m} = \frac{\rho_{T,m} + \Gamma_{T,m+1} e^{-j2\delta_m}}{1 + \rho_{T,m} \Gamma_{T,m+1} e^{-j2\delta_m}}$$

$$\delta_m = k_m \ell_m \cos(\theta)$$

Conclusions and next time

- This time we
 - Solve for the reflection via geometric series (ray tracing)
 - Extrapolate to multiple layers
 - Explore oblique incidence angle

- Stack of dielectric layers acts like a filter by modulating optical path lengths
 - Delayed and attenuated copies of the incident field

- This Thursday we'll do some numerical examples