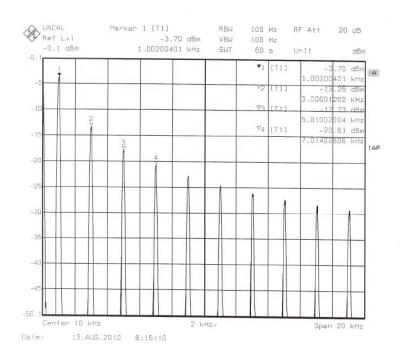
ELEC-A7200 Signals and Systems

Professor Riku Jäntti Fall 2022





Lecture 4 Fourier series

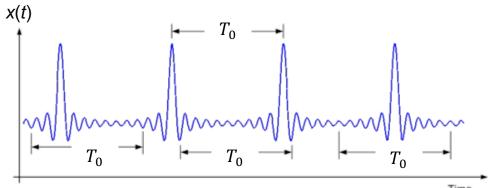
Contents

- Periodic signals
- Fourier series
 - Exponential Fourier series
 - Trigonometric Fourier series
- Amplitude and phase spectrum
- Parseval's theorem & Power spectrum
- Gibbs phenomena



Periodic signals

A signal x(t) is periodic with period T_0 if $x(t+T_0) = x(t)$

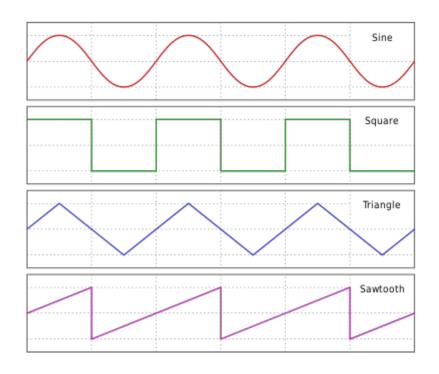


- The fundamental frequency of the signal is 1/ T_0
- Periodic signals are power signals

$$P = \frac{1}{T_0} \int_{T_0} |x(t)|^2 dt = \frac{1}{T_0} \int_{t_0}^{t_0 + T_0} |x(t)|^2 dt \, \forall t_0$$



Common periodic signals

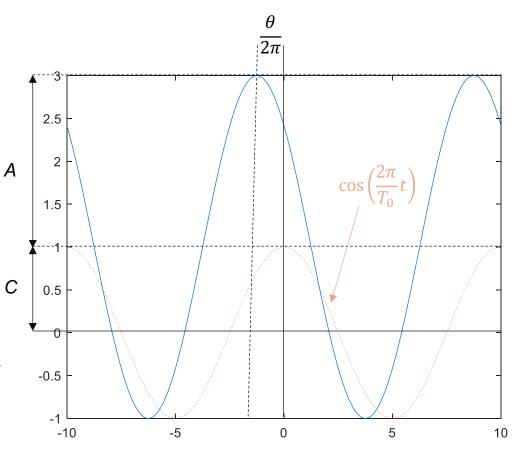




Sinusoids

$$x(t) = A\cos\left(\frac{2\pi}{T_0}t + \theta\right) + C$$

- Period T_0
- Frequency $\frac{1}{T_0}$
- Amplitude A
- Phase θ
- Offset (direct current DC component) C





Phase lead and lag

Phase lag = Delayed sinusoid

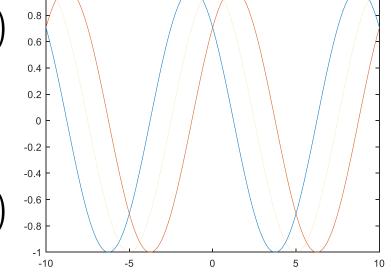
$$x(t) = A \cos\left(\frac{2\pi}{T_0}(t - \tau)\right) = A \cos\left(\frac{2\pi}{T_0}t + \theta\right)$$

$$\Rightarrow \theta = -\frac{2\pi}{T_0}\tau \text{ Phase lag}$$

Phase lead = Early signal

$$x(t) = A \cos\left(\frac{2\pi}{T_0}(t+\tau)\right) = A \cos\left(\frac{2\pi}{T_0}t+\theta\right)^{\frac{-0.4}{-0.6}}$$

$$\Rightarrow \theta = +\frac{2\pi}{T_0}\tau \text{ Phase lead}$$





Jean Baptiste Joseph Fourier

Jean Baptiste Joseph Fourier (1768-1830)

French mathematician, Egyptologist and administrator, who exerted strong influence on mathematical physics through his Théorie analytique de la chaleur (1822; The Analytical Theory of Heat).

Fourier was the first person to study the Earth's temperature from a mathematical perspective. In that study, he discovered the *greenhouse effect*.



https://www.britannica.com/biograph y/Joseph-Baron-Fourier

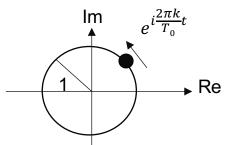


Signal representation in orthonormal basis: Fourier-series

Periodic signal $x(t) = x(t + T_0)$

Orthonormal basis

$$\phi_k(t) = \frac{1}{\sqrt{T_0}} e^{i\frac{2\pi k}{T_0}t}$$
 Complex signal
$$k = \dots -2, -1, 0, 1, 2, \dots$$



Phasor rotating counter clockwise with frequency $\frac{k}{T_0}$

Signal represented in the orthonormal basis

$$x(t) = \sum_{k=-\infty}^{\infty} \langle x(t), \phi_k(t) \rangle \phi_k(t) = \sum_{k=-\infty}^{\infty} \int_{T_0} x(t) \frac{1}{\sqrt{T_0}} e^{-i\frac{2\pi k}{T_0}t} dt \frac{1}{\sqrt{T_0}} e^{i\frac{2\pi k}{T_0}t} = \sum_{k=-\infty}^{\infty} X(k) e^{i\frac{2\pi k}{T_0}t}$$

$$X(k) = \frac{1}{T_0} \int_{T_0} x(t) e^{-i\frac{2\pi k}{T_0}t} dt$$

Coefficients of Exponential Fourier Series



Fourier basis

Basis functions of the exponential Fourier-series: $\phi_k(t) = \frac{1}{\sqrt{T_0}} e^{i\frac{2\pi k}{T_0}t}$

Proof that the basis is orthonormal:

$$\begin{split} \langle \phi_{k}(t), \phi_{l}(t) \rangle &= \frac{1}{T_{0}} \int_{T_{0}} e^{i\frac{2\pi(k-l)}{T_{0}}t} \, dt = \frac{1}{T_{0}} \int_{-T_{0}/2}^{T_{0}/2} e^{i\frac{2\pi(k-l)}{T_{0}}t} \, dt = \frac{1}{i2\pi(k-l)} \left(e^{i\pi(k-l)} - e^{-i\pi(k-l)} \right) \\ &= \frac{1}{\pi(k-l)} \frac{1}{i2} \left(e^{i\pi(k-l)} - e^{-i\pi(k-l)} \right) = \frac{1}{\pi(k-l)} \sin(k-l) = \sin(k-l) = \begin{cases} 0, k \neq l \\ 1, k = l \end{cases} \\ \frac{|\sin(x)| = \frac{1}{2i} \left(e^{ix} - e^{-ix} \right)|}{|\sin(x)|} \left[\frac{\sin(x)}{\pi x} \right] \end{split}$$

Sinc function

Sinc function

$$\operatorname{sinc}(x) = \frac{\sin(\pi x)}{\pi x}$$

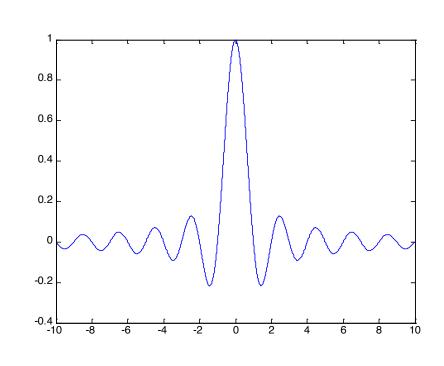
Zeros

$$\operatorname{sinc}(x) = \frac{\sin(\pi x)}{\pi x} = 0$$
$$x = \pm 1, \pm 2, \pm, 3, \dots$$

Limit

$$\lim_{x \to 0} \frac{\sin(\pi x)}{\pi x} = \lim_{x \to 0} \frac{\frac{\partial}{\partial x} \sin(\pi x)}{\frac{\partial}{\partial x} \pi x} = \lim_{x \to 0} \frac{\pi \cos(\pi x)}{\pi} = \frac{\pi}{\pi} = 1$$



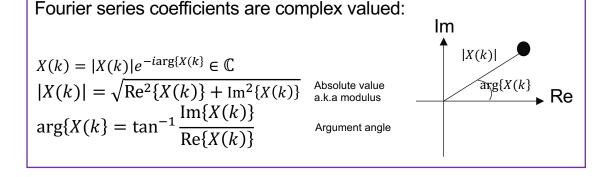


Signal representation in orthonormal basis: Fourier series

Exponential Fourier series representation of a periodic signal $x(t) = x(t + T_0)$

$$x(t) = \sum_{k=-\infty}^{\infty} X(k) e^{i\frac{2\pi k}{T_0}t}$$

$$X(k) = \frac{1}{T_0} \int_{T_0} x(t) e^{-i\frac{2\pi k}{T_0}t} dt$$





Fourier series properties

Let x(t) and y(t) be periodic signals with the same period T_0

- **Linearity** $z(t)=x(t)+y(t) \Rightarrow Z(k)=X(k)+Y(k)$
- **Komplex konjugation** $y(t) = x^*(t) \Rightarrow Y(k) = X^*(-k)$
- Time shifting $y(t)=x(t-t_0) \Rightarrow Y(t) \Rightarrow X(k)e^{-i\frac{2\pi k}{T_0}t_0}$
- Frequency shifting $y(t) = x(t)e^{i\frac{2\pi l}{T_0}t} \Rightarrow Y(k) = X(k-l)$
- Differentation $y(t) = \frac{d}{dt}x(t) \Rightarrow Y(k) = i\frac{2\pi k}{T_0}X(k)$
- Integration $y(t) = \int_{-\infty}^{t} x(\tau) d\tau \Rightarrow \frac{1}{i\frac{2\pi k}{T_0}} X(k)$



Exponential Fourier-series representation of a sinusoidal signal

Sinusoidal:

$$x(t) = C + A\cos\left(\frac{2\pi}{T_0}t + \theta\right) = \frac{A}{2}e^{-i\theta}e^{-i\frac{2\pi}{T_0}t} + C + \frac{A}{2}e^{i\theta}e^{i\frac{2\pi}{T_0}t}$$

$$x(t) = C + A\cos\left(\frac{2\pi}{T_0}t + \theta\right) = \frac{A}{2}e^{-i\theta}e^{-i\frac{2\pi}{T_0}t} + C + \frac{A}{2}e^{i\theta}e^{i\frac{2\pi}{T_0}t}$$

The modulus of the Fourier series coefficient is amplitude divided by 2

The argument angle of the Fourier series coefficient gives the phase of the sinusoidal

$$\cos(x) = \frac{1}{2}e^{ix} + \frac{1}{2}e^{-ix}$$

Trigonometrical Fourier Series

For real periodic signal $x(t) = x(t + T_0) \in \mathbb{R}$ we have

$$X(-k) = \frac{1}{T_0} \int_{T_0} x(t) e^{i\frac{2\pi k}{T_0}t} dt = X^*(k)$$
 Complex conjugate

By noting that $cos(x) = \frac{1}{2}(e^{-ix} + e^{ix})$, we can write

$$x(t) = X(0) + \sum_{k=1}^{\infty} 2|X(k)| \cos\left(2\pi \frac{k}{T_0}t + \arg\{X(k)\}\right)$$

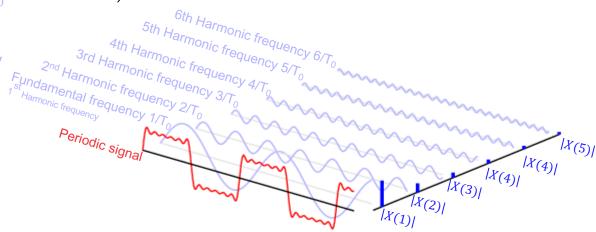


Trigonometrical Fourier Series

Periodic signal can be expressed as a sum of sinusoids with different frequencies and phases

$$x(t) = X(0) + \sum_{k=1}^{\infty} 2|X(k)| \cos\left(2\pi \frac{k}{T_0} t + \arg\{X(k)\}\right)$$

Harmonic frequency is an integer multiple of the fundamental frequency





Fourier sine and cosine series

Fourier series coefficients

$$X(k) = \frac{1}{T_0} \int_{T_0} x(t) e^{-i\frac{2\pi k}{T_0}t} dt = \frac{1}{T_0} \int_{T_0} x(t) \cos\left(\frac{2\pi k}{T_0}t\right) dt - i\frac{1}{T_0} \int_{T_0} x(t) \sin\left(\frac{2\pi k}{T_0}t\right) dt$$

$$a_k$$

$$b_k$$

If the periodic signal is real, then $X(k)=X^*(-k)$

$$x(t) = \sum_{k=-\infty}^{\infty} X(k) e^{i\frac{2\pi k}{T_o}t} = \sum_{k=-\infty}^{\infty} (a_k - ib_k) e^{i\frac{2\pi k}{T_o}t} = a_0 + \sum_{k=1}^{\infty} \left[\underbrace{(a_k + ib_k)e^{-i\frac{2\pi k}{T_o}t}} + \underbrace{(a_k - ib_k)e^{+i\frac{2\pi k}{T_o}t}} \right]$$

$$= a_0 + \sum_{k=1}^{\infty} 2a_k \cos\left(\frac{2\pi k}{T_o}\right) + \sum_{k=1}^{\infty} 2b_k \sin\left(\frac{2\pi k}{T_o}\right)$$
Aalto University
School of Electrical
$$\cos(x) = \frac{e^{ix} + e^{-ix}}{2} \quad \sin(x) = \frac{e^{ix} - e^{-ix}}{2i} \quad \frac{1}{i} = -i$$

Fourier sine and cosine series

Fourier sine and cosine series for a real signal

$$x(t) = \alpha_0 + \sum_{k=1}^{\infty} \alpha_k \cos\left(\frac{2\pi k}{T_0}\right) + \sum_{k=1}^{\infty} \beta_k \sin\left(\frac{2\pi k}{T_0}\right)$$

Cosine series coefficients

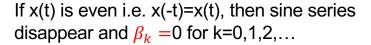
$$\alpha_0 = \frac{1}{T_0} \int_{T_0} x(t) dt$$

$$\alpha_k = \frac{2}{T_0} \int_{T_0} x(t) \cos\left(\frac{2\pi k}{T_0}t\right) dt$$

Sine series coefficients

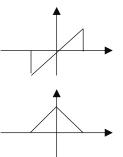
$$\beta_k = \frac{2}{T_0} \int_{T_0} x(t) \sin\left(\frac{2\pi k}{T_0}t\right) dt$$

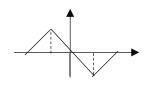
If x(t) is odd i.e. x(-t)=-x(t), then cosine series disappear and α_k =0 for k=0,1,2,...



If x(t) is half-wave symmetric $x(t)=-x(t+T_0/2)$, then even idex terms dissappear i.e.

$$\alpha_{2k} = \beta_{2k} = 0$$
 for k=0,1,2,...

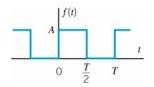






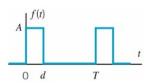
Fourier series repesentations of periodic signals

f(t)-A/2 is odd => Sine series



Square wave:
$$\omega_0 = \frac{2\pi}{T}$$

$$f(t) = \frac{A}{2} + \frac{4A}{\pi} \sum_{n=1}^{\infty} \frac{\sin((2n-1)\omega_0 t)}{2n-1}$$



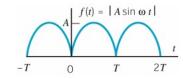
Pulse wave:
$$\omega_0 = \frac{2\pi}{T}$$

$$f(t) = \frac{Ad}{2} + \frac{2Ad}{\pi} \sum_{n=1}^{\infty} \frac{\sin\left(\frac{n\pi d}{T}\right)}{\frac{n\pi d}{T}} \cos\left(n\omega_0 t\right)$$

Half wave rectified sine wave: $\omega_0 = \frac{2\pi}{T}$

$$f(t) = \frac{A}{\pi} + \frac{A}{2}\sin\omega_0 t - \frac{2A}{\pi} \sum_{n=1}^{\infty} \frac{\cos(2n\omega_0 t)}{4n^2 - 1}$$

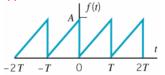
f(t) is even => Cosine series



Full wave rectified sine wave: $\omega_0 = \frac{2\pi}{T}$

$$f(t) = \frac{2A}{\pi} - \frac{4A}{\pi} \sum_{n=1}^{\infty} \frac{\cos(n\omega_0 t)}{4n^2 - 1}$$

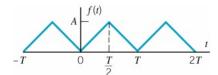
f(t)-A/2 is odd => Sine series



Sawtooth wave:
$$\omega_0 = \frac{2\pi}{T}$$

$$f(t) = \frac{A}{2} + \frac{A}{\pi} \sum_{n=0}^{\infty} \frac{\sin(n\omega_0 t)}{n}$$

f(t) is even => Cosine series



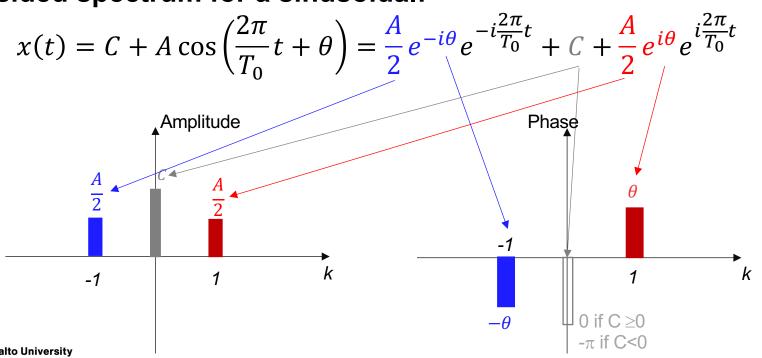
Triangle wave:
$$\omega_0 = \frac{2\pi}{T}$$

$$f(t) = \frac{A}{2} - \frac{4A}{\pi^2} \sum_{n=1}^{\infty} \frac{\cos((2n-1)\omega_0 t)}{(2n-1)^2}$$

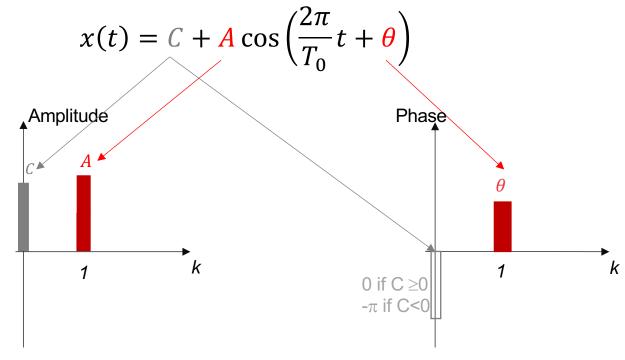


$$\cos(x) = \frac{1}{2}e^{ix} + \frac{1}{2}e^{-ix}$$

Two-sided spectrum for a sinusoidal:



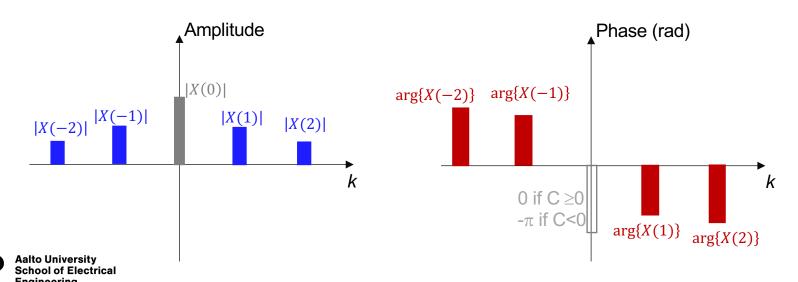
One-sided spectrum for a sinusoidal:





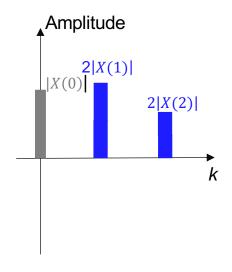
Two-sided amplitude and phase spectrum

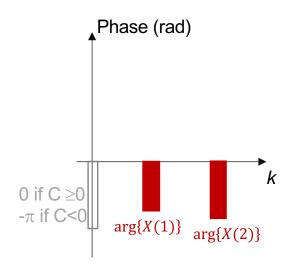
$$x(t) = \sum_{k=-\infty}^{\infty} |X(k)| e^{-i\arg\{X(k)\}} e^{i\frac{2\pi k}{T_0}t}$$



One-sided amplitude and phase spectrum (real signal)

$$x(t) = X(0) + \sum_{k=1}^{\infty} 2|X(k)| \cos\left(2\pi \frac{k}{T_0} t + \arg\{X(k)\}\right)$$







Parseval's theorem

Parsevals' theorem:

$$P = \frac{1}{T_0} \int_{T_0} |x(t)|^2 dt = \sum_{k = -\infty}^{\infty} |X(k)|^2$$

Squared modulus of the Fourier series coefficients $|X(k)|^2$ describe how the signal power is distributed among the harmonic frequencies.

Power spectrum

$$\cos(x) = \frac{1}{2}e^{ix} + \frac{1}{2}e^{-ix}$$

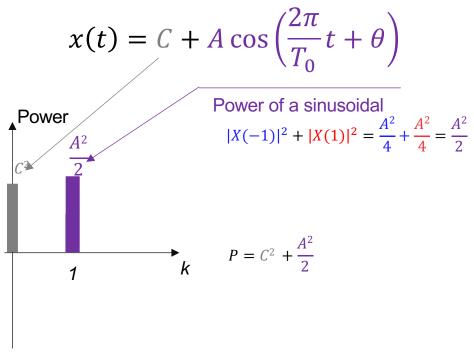
Two-sided power spectrum of a sinusoidal

$$x(t) = C + A\cos\left(\frac{2\pi}{T_0}t + \theta\right) = \frac{A}{2}e^{-i\theta}e^{-i\frac{2\pi}{T_0}t} + C + \frac{A}{2}e^{i\theta}e^{i\frac{2\pi}{T_0}t}$$
Power
$$x(-1) \qquad x(0) \qquad x(1)$$

$$P = \sum_{k=-1}^{1} |x(k)|^2 = \frac{A^2}{4} + C^2 + \frac{A^2}{4} = \frac{A^2}{2} + C^2$$
to University

Power Spectrum

One-sided spectrum for a sinusoidal:

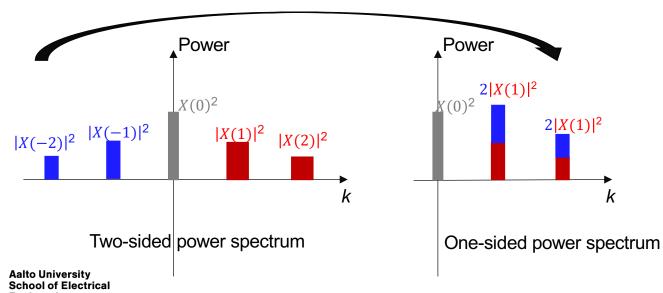


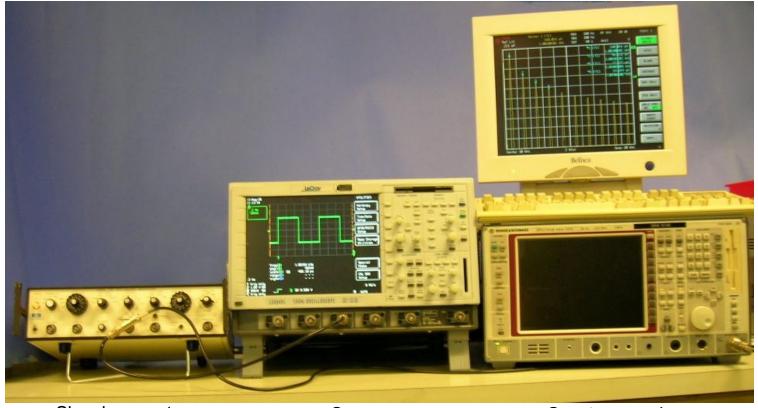


Power Spectrum

Fourier series of a real periodic signal

$$x(t) = \sum_{k=-\infty}^{\infty} |X(k)| e^{-i\arg\{X(k)\}} e^{i\frac{2\pi k}{T_0}t} = X(0) + \sum_{k=1}^{\infty} 2|X(k)| \cos\left(2\pi \frac{k}{T_0}t + \arg\{X(k)\}\right)$$





Signal generator

Scope

Spectrum analyzer

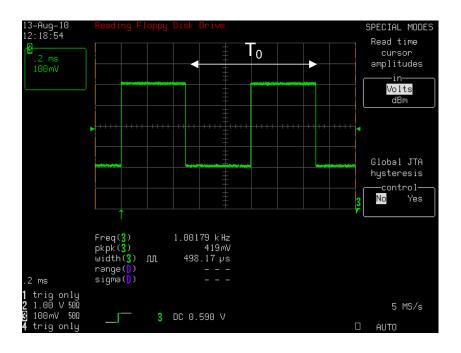


Square wave signal

$$x(t) = A\operatorname{sgn}\left(\sin\left(\frac{2\pi k}{T_0}t\right)\right)$$

$$\begin{split} X(k) &= \frac{1}{T_0} \int_{T_0} x(t) e^{-i\frac{2\pi k}{T_0}t} dt \\ &= -\frac{A}{T_0} \int_{-\frac{1}{2}T_0}^{0} e^{-i\frac{2\pi k}{T_0}t} dt + \frac{A}{T_0} \int_{0}^{\frac{1}{2}T_0} e^{-i\frac{2\pi k}{T_0}t} dt \\ &= -\frac{A}{-i2\pi k} \left(1 - e^{i\pi k}\right) + \frac{A}{-i2\pi k} (e^{-i\pi k} - 1) \\ &= \frac{A}{i2\pi k} \left(1 - e^{i\pi k} + 1 - e^{-i\pi k}\right) = -i\frac{A}{\pi k} 2(1 - \cos(\pi k)) \\ &= -i\frac{2A}{\pi k} \sin^2\left(\frac{\pi k}{2}\right) \\ &= -i\frac{2A}{\pi k} \sin^2\left(\frac{\pi k}{2}\right) \end{split}$$

$$\mathbf{A2} \quad \text{Aalto University School of Electrical Engineering} \quad = \begin{cases} -i\frac{2A}{\pi k} & k \text{ odd} \\ 0 & k \text{ even} \end{cases}$$



$$\cos(x) = \frac{1}{2} (e^{ix} + e^{-ix})$$
 $\sin^2(x) = \frac{1 - \cos(2x)}{2}$

Example: Square wave (again)

Square wave signal

$$x(t) = A\operatorname{sgn}\left(\sin\left(\frac{2\pi k}{T_0}t\right)\right)$$

is odd => cosine series vanish.

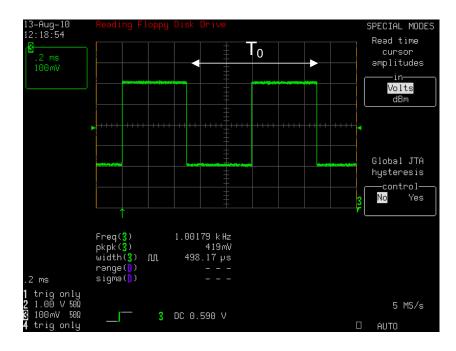
$$\beta_{k} = \frac{2}{T_{0}} \int_{T_{0}} x(t) \sin\left(\frac{2\pi k}{T_{0}}t\right) dt$$

$$-\frac{2A}{T_{0}} \int_{-\frac{T_{0}}{2}}^{0} \sin\left(\frac{2\pi k}{T_{0}}t\right) dt + \frac{2A}{T_{0}} \int_{0}^{\frac{T_{0}}{2}} \sin\left(\frac{2\pi k}{T_{0}}t\right) dt$$

$$= \frac{4A}{T_{0}} \int_{0}^{\frac{T_{0}}{2}} \sin\left(\frac{2\pi k}{T_{0}}t\right) dt = \frac{2A}{\pi k} (1 - \cos(\pi k)) = \frac{4A}{\pi k} \sin^{2}\left(\frac{\pi k}{2}\right)$$



$$= \begin{cases} \frac{4A}{\pi k} & k \text{ odd} \\ 0 & k \text{ even} \end{cases}$$



$$\int \sin(ax) dx = -\frac{1}{a}\cos(ax) + c$$
$$\sin^2(x) = \frac{1 - \cos(2x)}{2}$$

Example: Square wave (again)

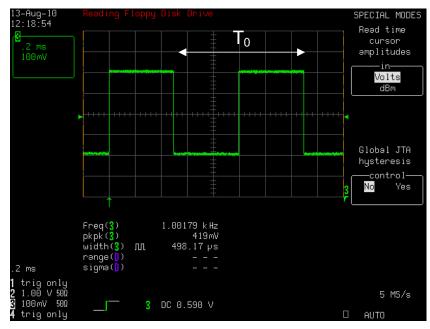
Fourier cosine series coefficients

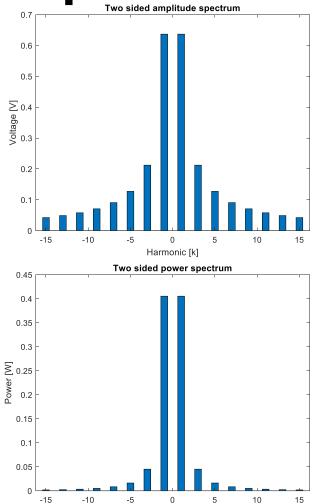
$$\alpha_k = \mathbf{0}$$

Fourier sine series coefficients

$$\beta_k = \frac{4A}{\pi k} \sin^2\left(\frac{\pi k}{2}\right)$$

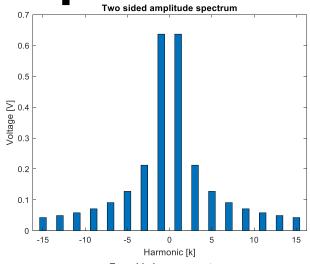
$$X(k) = \mathbf{a}_k - i\mathbf{b}_k = 0 - i\frac{\beta_k}{2}$$
$$= -i\frac{4A}{\pi k}\sin^2\left(\frac{\pi k}{2}\right)$$

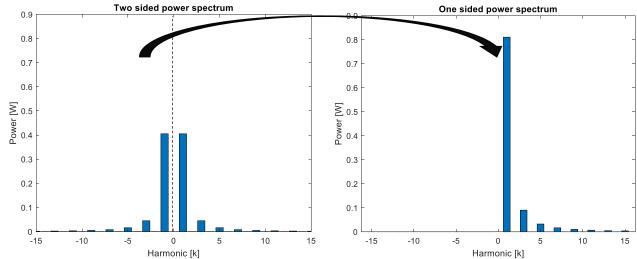




Harmonic [k]

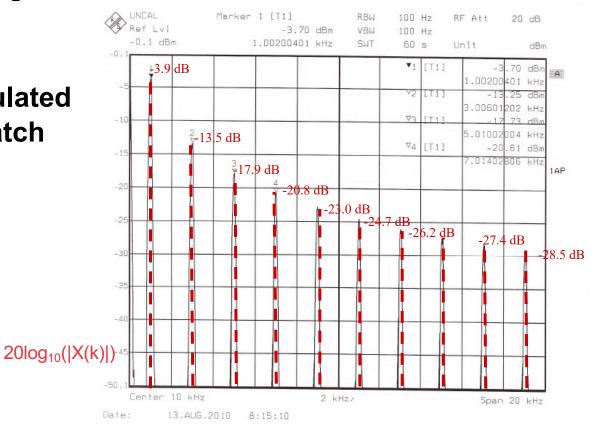








 Measured and calculated power spectrum match well.





Approximating a periodic with limited number of coefficients

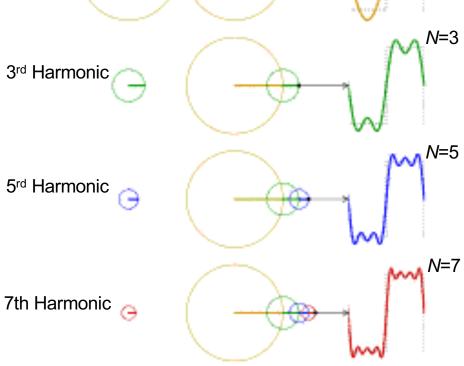


Fourier series

$$x(t) = \sum_{k=-\infty}^{\infty} X(k) e^{i\frac{2\pi k}{T_0}t}$$

Approximation

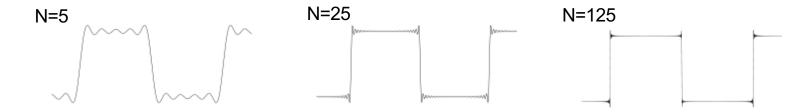
$$x_N(t) = \sum_{k=-N}^{N} X(k) e^{i\frac{2\pi k}{T_0}t}$$





Gibbs phenomena

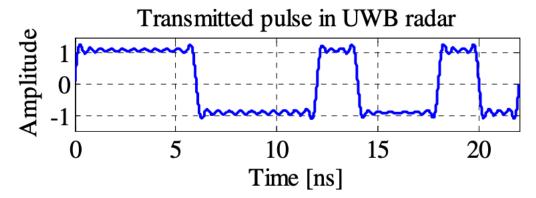
The Gibbs phenomenon involves both the fact that Fourier sums overshoot at a jump discontinuity, and that this overshoot does not die out as more terms are added to the sum.





Approximating a periodic with limited number of coefficients

Example: Generation of ultrawideband radar signal

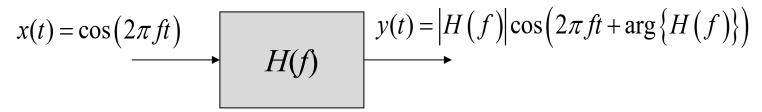


A UWB pulse signal is generated by Fourier series model using 40 harmonies. The pulse is produced by the Barker code with the length of 11, chip width of 2 ns and the total pulse duration 22 of ns



Filtering periodic signals

Given LTI system frequency transfer function H(f), we have



Hence for input
$$x(t) = X(0) + \sum_{k=1}^{\infty} 2|X(k)| \cos(2\pi \frac{k}{T_0}t + \arg\{X(k)\})$$

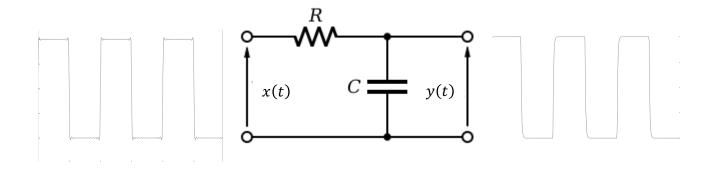
the output becomes

$$y(t) = H(0)X(0) + \sum_{k=1}^{\infty} 2\left| H\left(\frac{k}{T_0}\right) \right| |X(k)| \cos\left(2\pi \frac{k}{T_0}t + \arg\{X(k)\} + \arg\left\{H\left(\frac{k}{T_0}\right)\right\}\right)$$



Filtering periodic signals

 Example: Square wave passes through a lowpass RC-filter







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