## Some properties of Laplace transformation

Definition: 
$$\mathcal{L}{f(t)} = \overline{F}(s) = \int_{0}^{\infty} f(t)e^{-st}dt$$
  
Linearity:  $\mathcal{L}{af(t) + bg(t)} = a\overline{F}(s) + b\overline{G}(s)$   
 $\mathcal{L}{f(t)g(t)} \neq \overline{F}(s)\overline{G}(s)$ 

Derivative :  $\mathcal{L}{f'(t)} = s\overline{F}(s) - f(0)$ 

Integral: 
$$\mathcal{L}\left\{\int_{0}^{t} f(u)du\right\} = \frac{1}{s}\overline{F}(s)$$

Convolution:  $\mathcal{L}^{-1}\left\{\overline{F}(s)\overline{G}(s)\right\} = \int_{0}^{t} f(u)g(t-u)du$ 

Translation:  $\mathcal{L}\left\{e^{at}f(t)\right\} = \overline{F}(s-a)$ 

## **Time-dependent aka Transient methods**

Time-dependent methods are particularly suitable for the determination of the reaction kinetics, and they provide often a clear qualitative picture of what is going on in an electrochemical cell. We concentrate, again, on the trace-ion case in the absence of migration and convection, leaving diffusion as the only mechanism of mass transport.

Since we are considering time-dependent behavior, we have to solve Fick's 2<sup>nd</sup> law (in 1-D):

$$\frac{\partial c_k}{\partial t} = D_k \frac{\partial^2 c_k}{\partial x^2}$$
(7.1)

An initial condition is needed which is

$$c_k(x,t=0) = c_k^b \tag{7.2}$$

The bulk concentration  $c_k^b$  can naturally be also zero if the species 'k' is initially absent. "Far from the electrode" (a few  $\mu$ m) concentration naturally is  $c_k^b$ :

$$c_k(x \to \infty, t) = c_k^b \tag{7.3}$$

The current boundary condition completes the problem set-up, as usual:

$$\frac{i(t)}{nF} = \pm D_k \left(\frac{\partial c_k}{\partial x}\right)_{x=0}$$
(7.4)

In an electrochemical problem, in order to obtain the dependence of current on potential, we need, in addition, the values of the surface concentrations of the electroactive species. We have two choices: for a very fast (reversible) reaction we use Nernst equation and for a slow (quasi-reversible) reaction Butler-Volmer equation:

$$\frac{c_{\rm o}^{\rm s}}{c_{\rm R}^{\rm s}} = \theta = \exp\left[\frac{nF}{RT}\left(E - E^{0'}\right)\right] \quad \text{reversible} \quad (7.5) \qquad \qquad \frac{i(t)}{nF} = k_{ox}c_{\rm R}^{\rm s} - k_{red}c_{\rm o}^{\rm s} \qquad \text{quasi-reversible} \quad (7.6)$$

Note that *E* can be time dependent, making also  $k_{ox}$  and  $k_{red}$  potential dependent; they assume the forms of Butler-Volmer kinetics.

How these boundary conditions are used depends on the transient methods used, but the general solution of the Fick's law proceeds applying **Laplace transform**. Eq. (7.1) is transformed using the derivative property of Laplace transform

$$s\overline{c}_{k} - c_{k}^{b} = D_{k} \frac{d^{2}\overline{c}_{k}}{dx^{2}} \iff \frac{d^{2}\overline{c}_{k}}{dx^{2}} - \frac{s}{D_{k}}\overline{c}_{k} = -\frac{c_{k}^{b}}{D_{k}}$$
(7.7)

where the upper bar denotes a variable in the Laplace domain and *s* is the Laplace variable (do not confuse it with the superscript '*s*' denoting for surface). The general solution of this differential equation is

$$\overline{c}_{k}(x,s) = \frac{c_{k}^{b}}{s} + A(s)e^{-\lambda_{k}x} + B(s)e^{\lambda_{k}x}, \quad \lambda_{k} = \sqrt{\frac{s}{D_{k}}}$$
(7.8)

A(s) and B(s) are coefficients to be determined from the boundary conditions. In order to have  $\overline{c}_k$  limited when  $x \to \infty$  (boundary condition, eq. (7.3)), it is required that B(s) = 0. A(s) is found using the current boundary condition (7.4):

$$\frac{\bar{i}(s)}{nF} = \pm D_k \left(\frac{d\bar{c}_k}{dx}\right)_{x=0} = \mp \sqrt{sD_k} A(s) \implies A(s) = \mp \frac{\bar{i}(s)}{nF\sqrt{sD_k}}$$

$$\overline{c}_k^s = \frac{c_k^b}{s} \mp \frac{\bar{i}(s)}{nF\sqrt{sD_k}}$$
(7.9)

where the upper minus sign corresponds to an oxidation and the lower plus sign to a reduction reaction. It has to be emphasized that eq. (7.9) is completely general, it applies to any time-dependent electrochemical method (for a trace-ion), as well as to the impedance method where a periodical perturbation signal is applied. The concentration profile in Laplace domain is

$$\overline{c}_{k}^{s}(x,s) = \frac{c_{k}^{b}}{s} \mp \frac{\overline{i}(s)}{nF\sqrt{sD_{k}}} e^{-\sqrt{s/D_{k}}x}$$
(7.10)

In order to proceed, a particular time-dependent method must be defined. We start with the simplest possible method, a potential step.

## **Potential step**

In the potential step experiment, electrode potential is stepped from the region where practically no current is flowing to a value where oxidation/reduction reaction occurs. Because potential remains constant,  $\theta$  in eq. (7.5) is constant and we have

$$c_o^s = \theta c_R^s \implies \overline{c}_o^s = \theta \overline{c}_R^s$$
  
Let's consider the case where 'O' is initially absent. Now:  $\overline{c}_R^s = \frac{c_R^b}{s} - \frac{\overline{i}(s)}{nF_o/sD_o}$  and  $\overline{c}_o^s = \frac{\overline{i}(s)}{nF_o/sD_o}$ . Applying the above equation,

$$\bar{i}(s) = \frac{nF\sqrt{D_o}\theta c_R^b}{1+\xi\theta}\frac{1}{\sqrt{s}} \quad , \quad \xi = \sqrt{\frac{D_o}{D_R}}$$
(7.11)

Current in the time domain is obtained from an inverse transformation which can be found tabulated. The inverse transformation of  $s^{-1/2}$  is  $(\pi t)^{-1/2}$ . Thus,

$$i(t) = \frac{nF\sqrt{D_o}\Theta c_{\rm R}^b}{1+\xi\Theta} \frac{1}{\sqrt{\pi t}}$$
(7.12)

If the step is done in so anodic potential that  $c_{R}^{s}$  goes to zero ( $\xi \theta \gg 1$ ), current is immediately obtained from eq. (7.9) as

$$\bar{i}_{d}(s) = \frac{nF\sqrt{D_{R}}c_{R}^{b}}{\sqrt{s}} \implies i_{d}(t) = \frac{nF\sqrt{D_{R}}c_{R}^{b}}{\sqrt{\pi t}}$$
Cottrell equation (7.13)

Inserting eq. (7.11) into eq. (7.10), the concentration profile is obtained in Laplace domain as

$$\overline{c}_{R}(x,s) = \frac{c_{R}^{b}}{s} - \frac{\xi \theta c_{R}^{b}}{(1+\xi \theta)} \cdot \frac{e^{-\sqrt{s/D_{R}}x}}{s}$$

The inverse transform is found in the table as

$$\frac{c_{\rm R}(x,t)}{c_{\rm R}^b} = 1 - \frac{\xi\theta}{(1+\xi\theta)} \operatorname{erfc}\left[\frac{x}{2\sqrt{D_{\rm R}t}}\right] = \operatorname{erf}\left[\frac{x}{2\sqrt{D_{\rm R}t}}\right] \quad ; \quad \xi\theta >> 1$$



erf(x) is the error function and erfc(x) its complement.

Aside: Concentration profile after a potential step to a value where surface concentration goes to zero;  $D = 10^{-5}$  cm<sup>2</sup>/s. In a general case, surface concentrations jump at the value where eq. (7.5) is valid.

$$m erf(x) = rac{2}{\sqrt{\pi}} \int_{0}^{x} e^{-u^{2}} du \ ; \ erfc(x) = 1 - erf(x)$$

## **Current step**

Another easy technique to analyze is to apply constant current density *i* in the cell; Laplace transform of 1 is 1/s. Hence,

 $\overline{c}_k^s = \frac{c_k^b}{s} \mp \frac{i}{nF\sqrt{D_k}} \frac{1}{s^{3/2}}$ 

Assuming again that "O" is initially absent, the inverse transforms are found easily as

$$c_{\rm R}^{s}(t) = c_{\rm R}^{b} - \frac{2i\sqrt{t}}{nF\sqrt{\pi D_{\rm R}}}$$
 and  $c_{\rm O}^{s}(t) = \frac{2i\sqrt{t}}{nF\sqrt{\pi D_{\rm O}}}$ 

The moment when  $c_{s}^{s}$  goes to zero is known as the **transition time**,  $\tau$ . It is give by **Sand equation**:



Using the transition time, the surface concentrations can be written in a compact form as

$$c_{\scriptscriptstyle R}^{s} = c_{\scriptscriptstyle R}^{b} \left( 1 - \sqrt{t/\tau} \right)$$
 and  $c_{\scriptscriptstyle O}^{s} = c_{\scriptscriptstyle R}^{b} \xi^{-1} \sqrt{t/\tau}$ 

where  $\boldsymbol{\xi}$  is defined by eq. (7.11).

Inserting these surface concentrations in Nernst equation,

$$E = E^{0'} - \frac{RT}{nF} \ln \left[ \xi \left( \sqrt{\frac{\tau}{t}} - 1 \right) \right]$$

Plot aside shows how potential changes as the function of time. At  $t = \tau$  potential jumps abruptly which gives the experimental means of determining  $\tau$ .

The concentration profile is calculated as follows:

$$\overline{c}_{R}^{s} = \frac{c_{R}^{b}}{s} - \frac{i}{nF\sqrt{D_{R}}} \frac{1}{s^{3/2}} \exp\left(-\sqrt{\frac{s}{D_{R}}}x\right)$$

$$c_{R}(x,t) = c_{R}^{b} - \frac{i}{nF\sqrt{D_{R}}} \left[2\sqrt{\frac{t}{\pi}} \exp\left(-\frac{x^{2}}{4D_{R}t}\right) - \frac{x}{\sqrt{D_{R}}} \operatorname{erfc}\left(\frac{x}{2\sqrt{D_{R}t}}\right)\right]$$

$$\frac{c_{R}(x,t)}{c_{R}^{b}} = 1 - \sqrt{\frac{t}{\tau}} \exp\left(-\frac{x^{2}}{4D_{R}t}\right) + \frac{\sqrt{\pi x}}{2\sqrt{D_{R}\tau}} \operatorname{erfc}\left(\frac{x}{2\sqrt{D_{R}t}}\right)$$

Aside simulated concentration profiles with  $D = 10^{-5}$  cm<sup>2</sup>/s and  $\tau = 4$  s.

Note the difference between potential and current steps: in the former surface concentration and in the latter the concentration gradient at the surface is constant.

