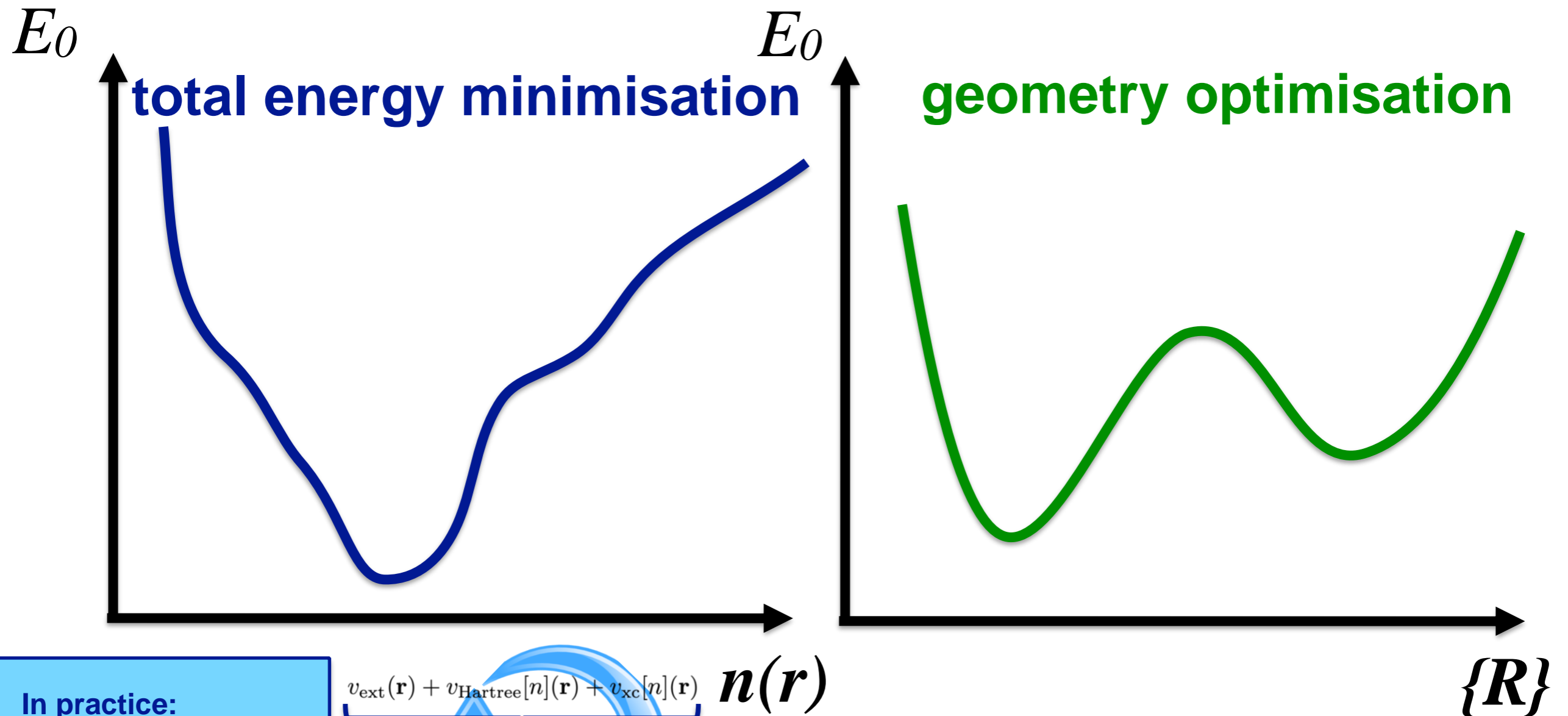


Density-Functional Theory for Practitioners - Tutorial 4

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Recap - two minimisations in DFT structure search

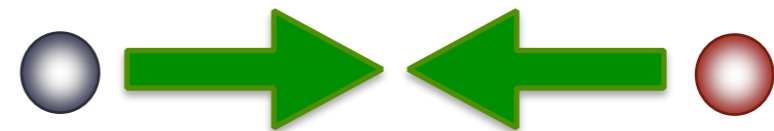


In practice:
 1. start with trial density
 2. iterate to self-consistency

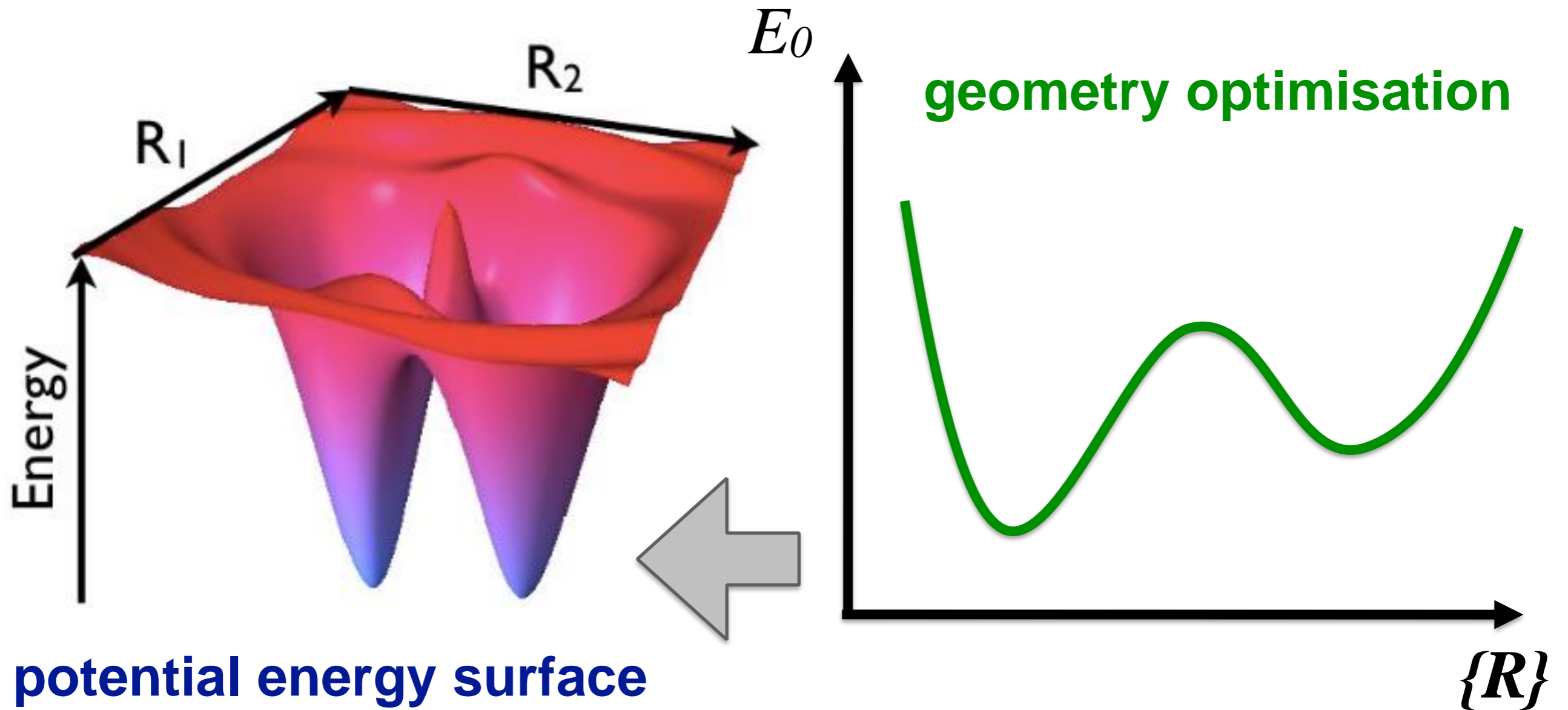
$$\underbrace{v_{\text{ext}}(\mathbf{r}) + v_{\text{Hartree}}[n](\mathbf{r}) + v_{\text{xc}}[n](\mathbf{r})}_{\text{potential}} n(\mathbf{r})$$

$$\left[-\frac{\nabla^2}{2} + v_{\text{aux}}[n](\mathbf{r}) \right] \phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$$

$$n(\mathbf{r}) = \sum_i^{N_e} |\phi_i(\mathbf{r})|^2$$

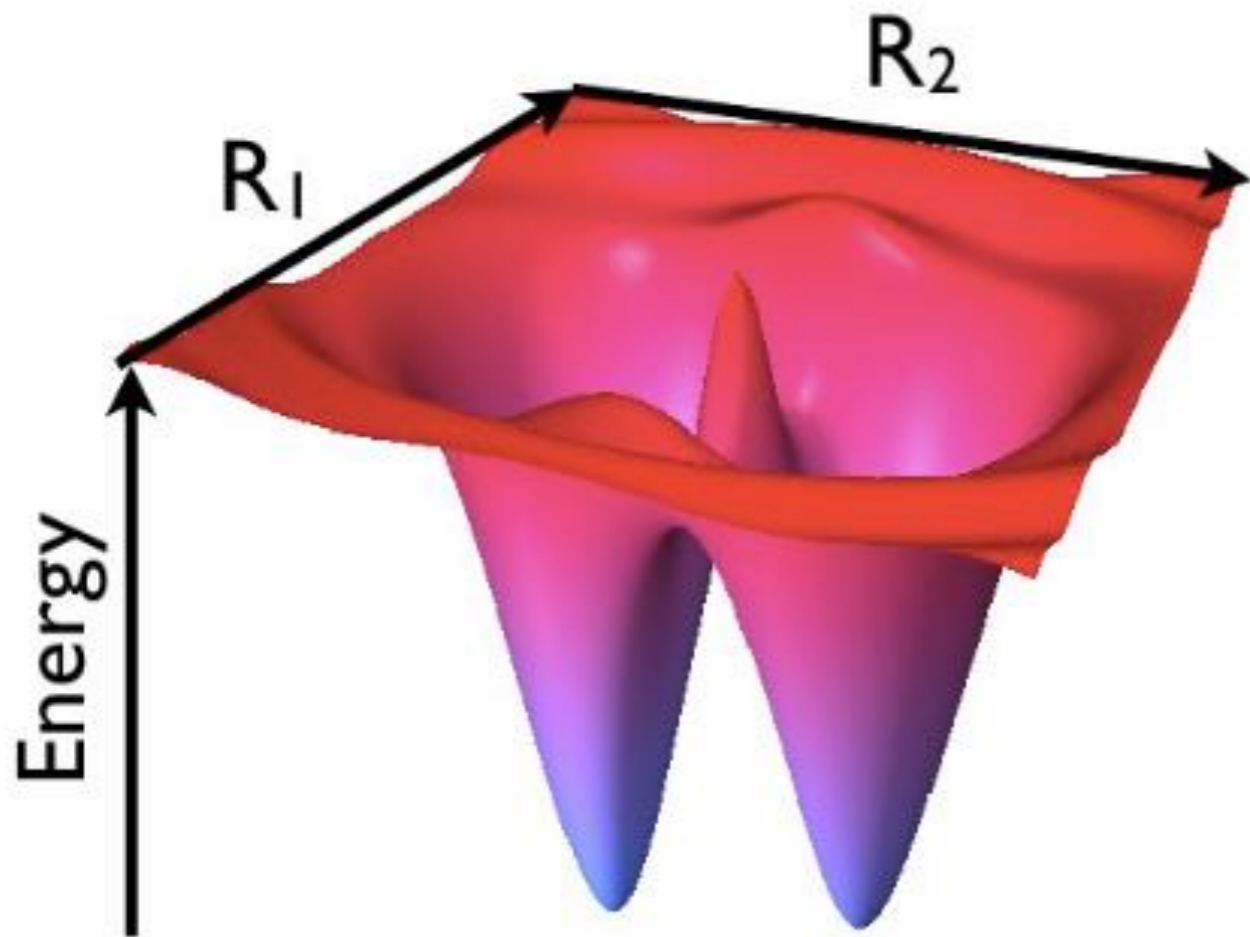


DFT potential energy surface



potential energy surface

DFT molecular dynamics (MD)

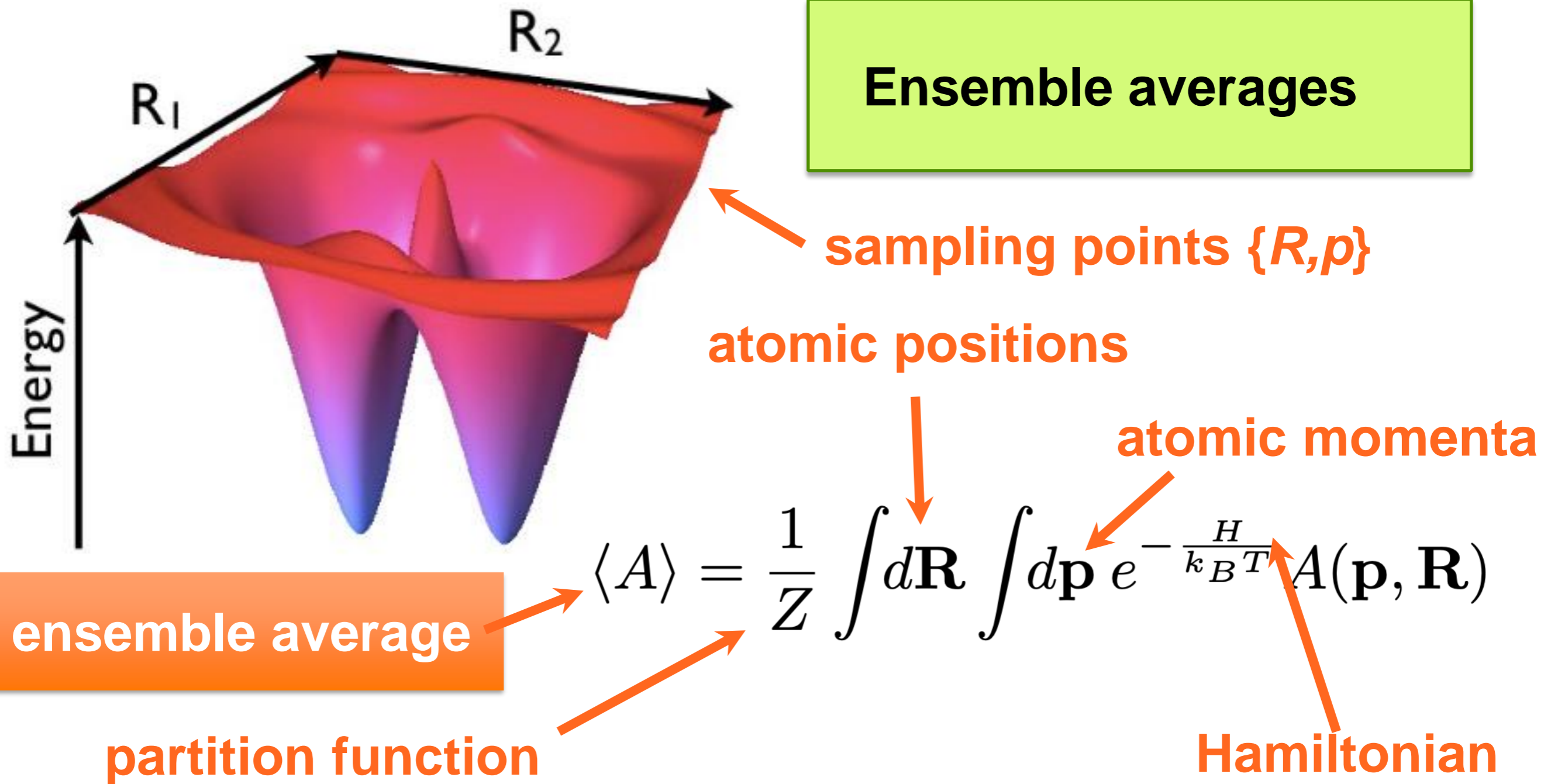


Molecular dynamics
or how to move on a
potential energy
surface.

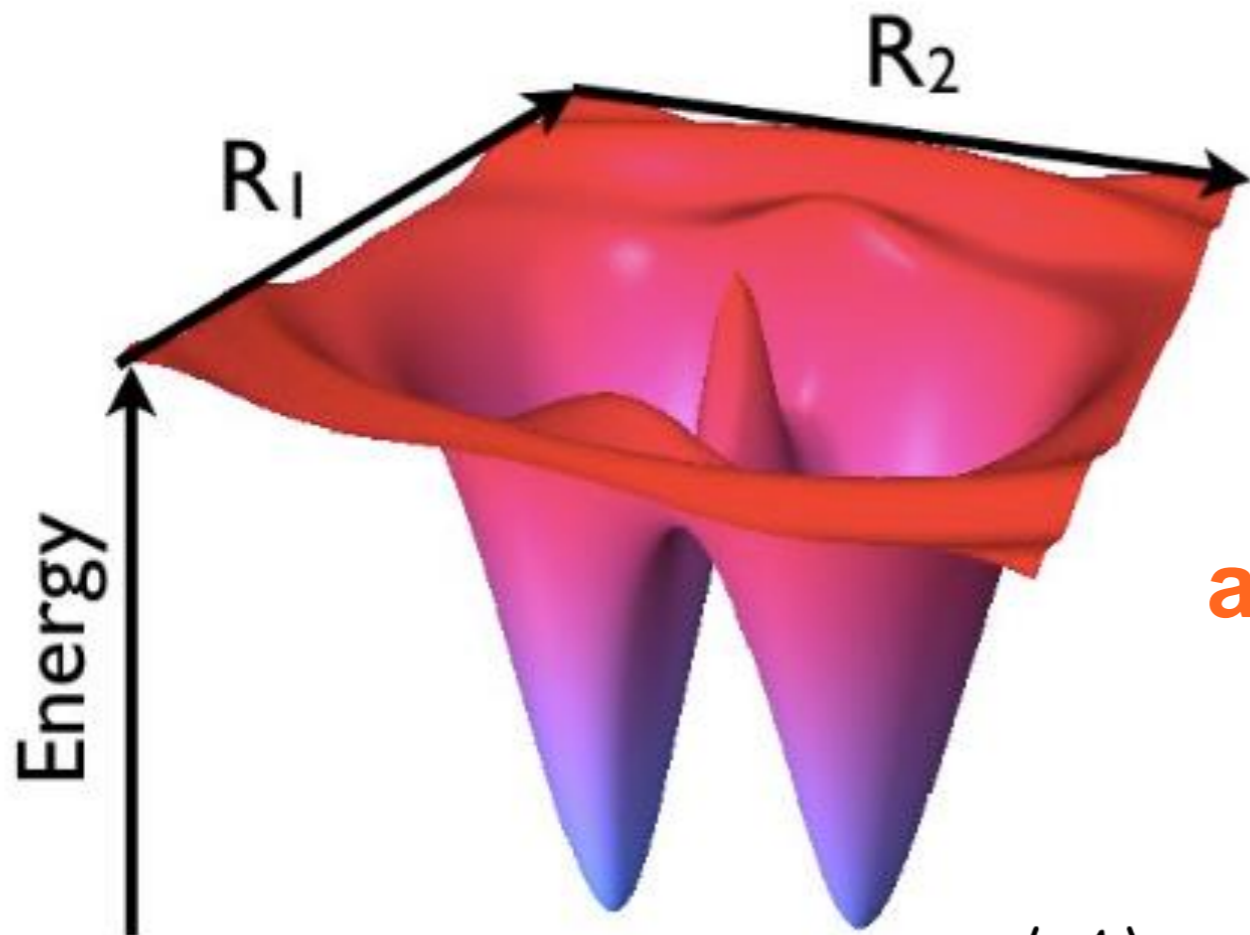
potential energy surface



Statistical sampling



Statistical sampling



Ensemble averages

atomic positions

atomic momenta

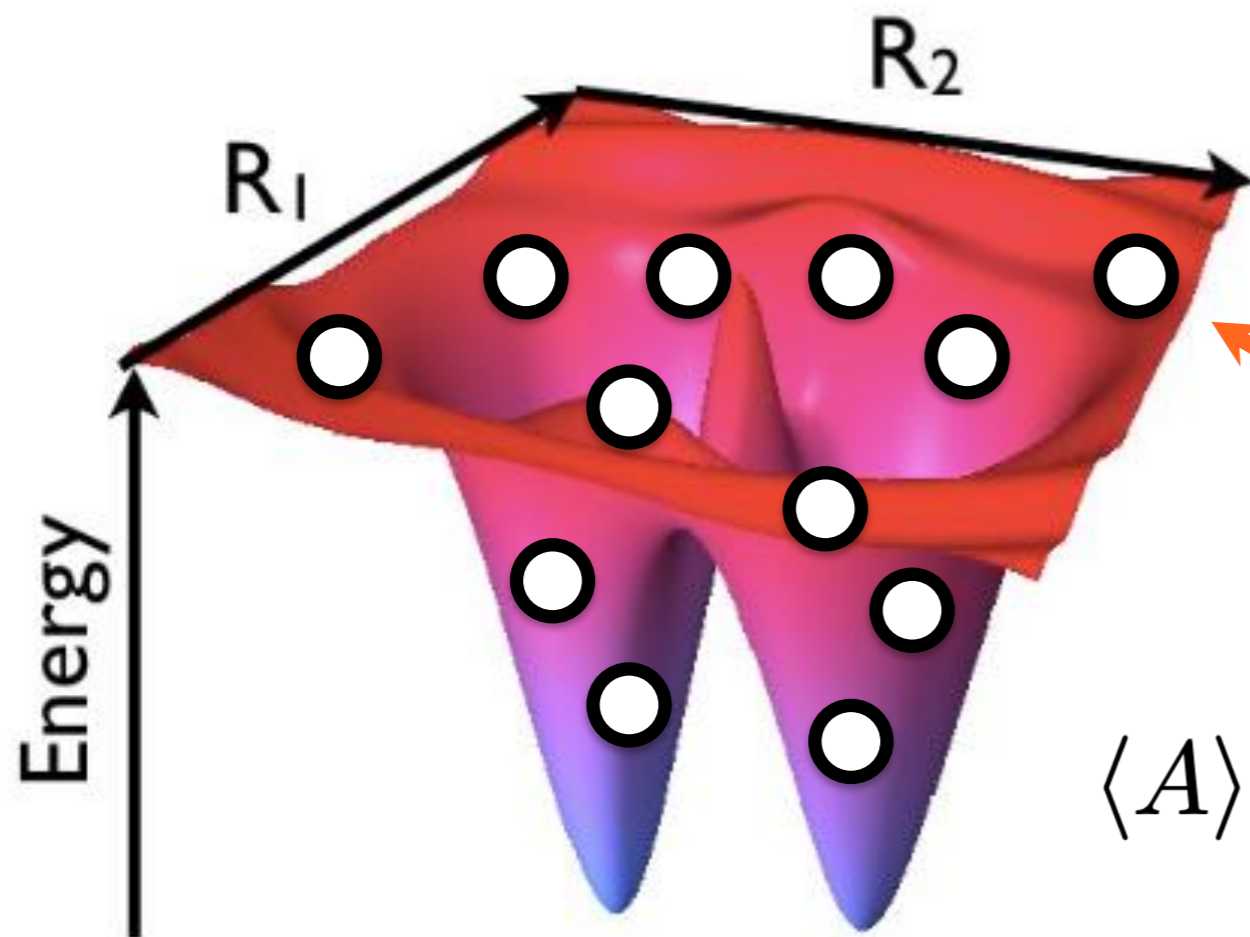
ensemble average

$$\langle A \rangle = \frac{1}{Z} \int d\mathbf{R} \int d\mathbf{p} e^{-\frac{H}{k_B T}} A(\mathbf{p}, \mathbf{R})$$

partition function

$$Z = \int d\mathbf{R} \int d\mathbf{p} e^{-\frac{H}{k_B T}}$$

Statistical sampling



Ensemble averages

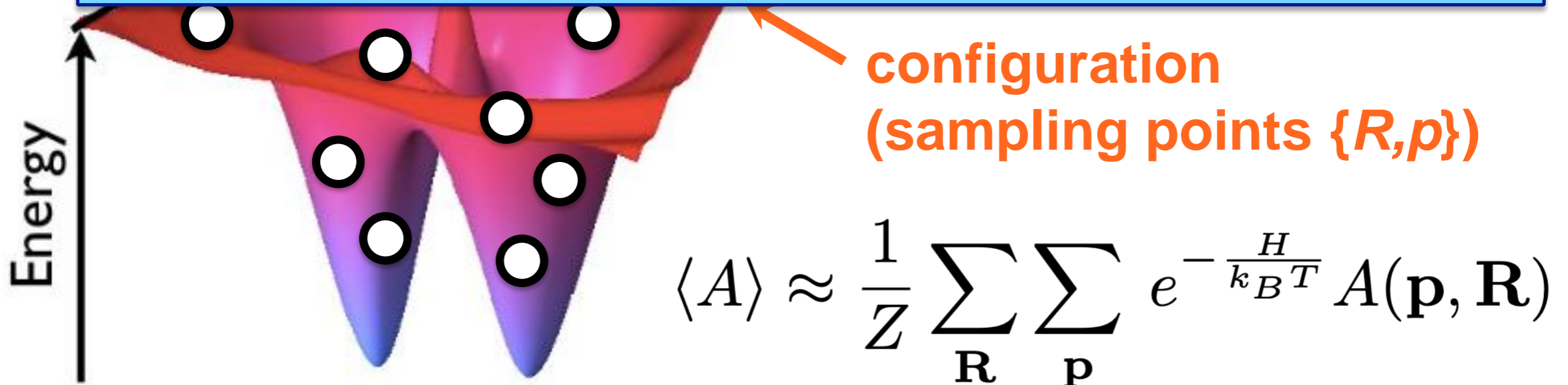
configuration
(sampling points $\{R,p\}$)

$$\langle A \rangle \approx \frac{1}{Z} \sum_{\mathbf{R}} \sum_{\mathbf{p}} e^{-\frac{H}{k_B T}} A(\mathbf{p}, \mathbf{R})$$

Compute by summing over configurations.

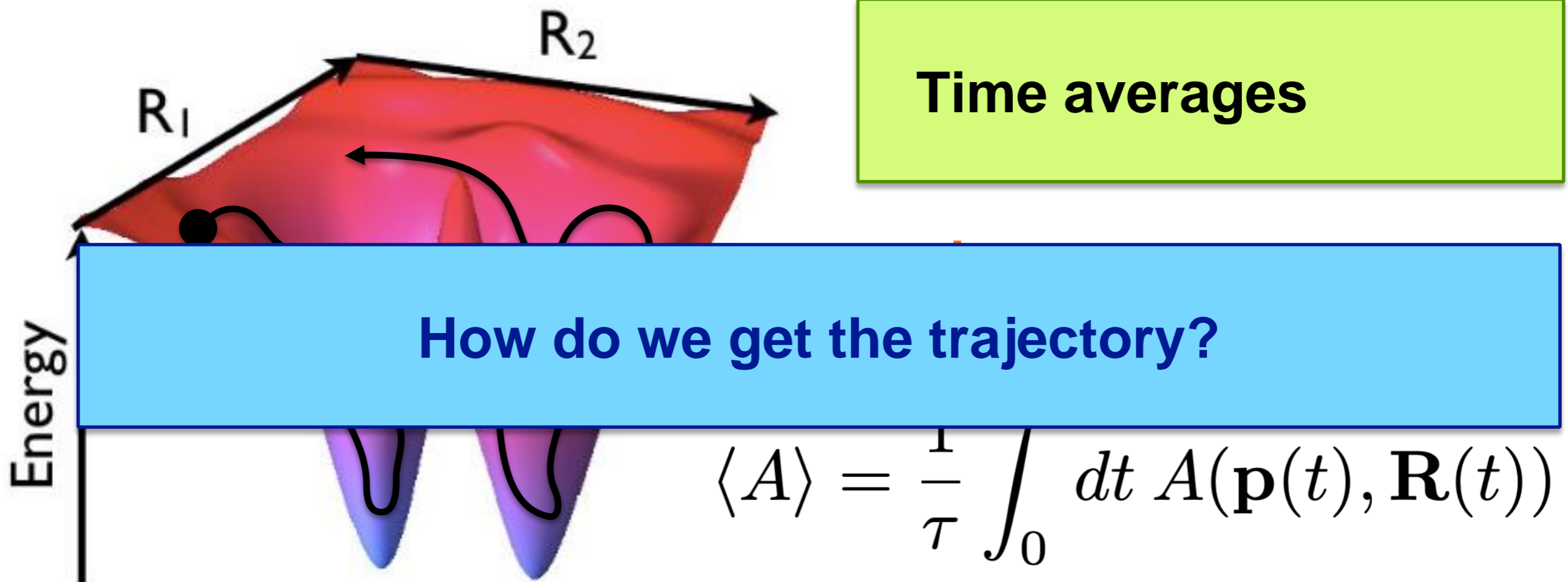
Statistical sampling

Is there a different way we can proceed?



- We might need a lot of configurations.
- We do not necessarily know how to find them.

Statistical sampling



Ergodic hypothesis:

ensemble averages equal time averages

Remember lecture 1

We could propagate the Schrödinger equation:

$$\mathcal{H}\Psi(t) = i\frac{\partial}{\partial t}\Psi(t)$$

**This would treat the nuclei quantum mechanically,
but that is maybe not what we want.**

Back to classical nuclei and DFT.

The full Hamiltonian

$$H(\text{●●●} \dots \text{●●●} \dots)$$

$$\mathcal{H} = \underbrace{-\sum_{I=1}^{N_n} \frac{\nabla_{\mathbf{R}_I}^2}{2M_I}}_{T^n} + \underbrace{\sum_{I=1}^{N_n} \sum_{J>I}^{N_n} \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|}}_{V^{nn}}$$

$$\underbrace{-\sum_{i=1}^{N_e} \frac{\nabla_{\mathbf{r}_i}^2}{2}}_{T^e} + \underbrace{\sum_{i=1}^{N_e} \sum_{j>i}^{N_e} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}}_{V^{ee}} + \underbrace{\sum_{i=1}^{N_e} \sum_{J=1}^{N_n} \frac{-Z_J}{|\mathbf{r}_i - \mathbf{R}_J|}}_{V^{ne}}$$

The full Hamiltonian

$$H(\text{●} \text{●} \text{●} \dots \text{●}^- \text{●}^- \text{●}^- \dots)$$

$$\mathcal{H} = \underbrace{-\sum_{I=1}^{N_n} \frac{\nabla_{\mathbf{R}_I}^2}{2M_I}}_{T^n} + \underbrace{\sum_{I=1}^{N_n} \sum_{J>I}^{N_n} \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|}}_{V^{nn}}$$

$$+ E_{\text{DFT}}^e(\{\mathbf{R}\})$$

We have solved the electronic problem.

The full Hamiltonian

$$H(\text{●} \text{●} \text{●} \dots \text{●}^- \text{●}^- \text{●}^- \dots)$$

$$\mathcal{H} = \underbrace{-\sum_{I=1}^{N_n} \frac{\nabla_{\mathbf{R}_I}^2}{2M_I}}_{T^n} + \underbrace{\sum_{I=1}^{N_n} \sum_{J>I}^{N_n} \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|}}_{V^{nn}}$$

$$+ E_{\text{DFT}}^e(\{\mathbf{R}\})$$

potential energy surface

$$V(\{\mathbf{R}\})$$

Classical Hamiltonian for the nuclei

$$H(\text{●} \text{●} \text{●} \dots)$$

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{\mathbf{p}_I^2}{2M_I} + V(\{\mathbf{R}\})$$

classical Hamiltonian

Molecular dynamics

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{\mathbf{p}_I^2}{2M_I} + V(\{\mathbf{R}\})$$

classical Hamiltonian

$$\mathbf{F}_I = -\frac{\partial \mathcal{H}}{\partial \mathbf{R}_I} = -\frac{\partial V(\{\mathbf{R}\})}{\partial \mathbf{R}_I}$$

Forces

$$\begin{aligned}\mathbf{F}_I &= M_I \ddot{\mathbf{R}}_I \\ &= M_I \mathbf{a}_I\end{aligned}$$

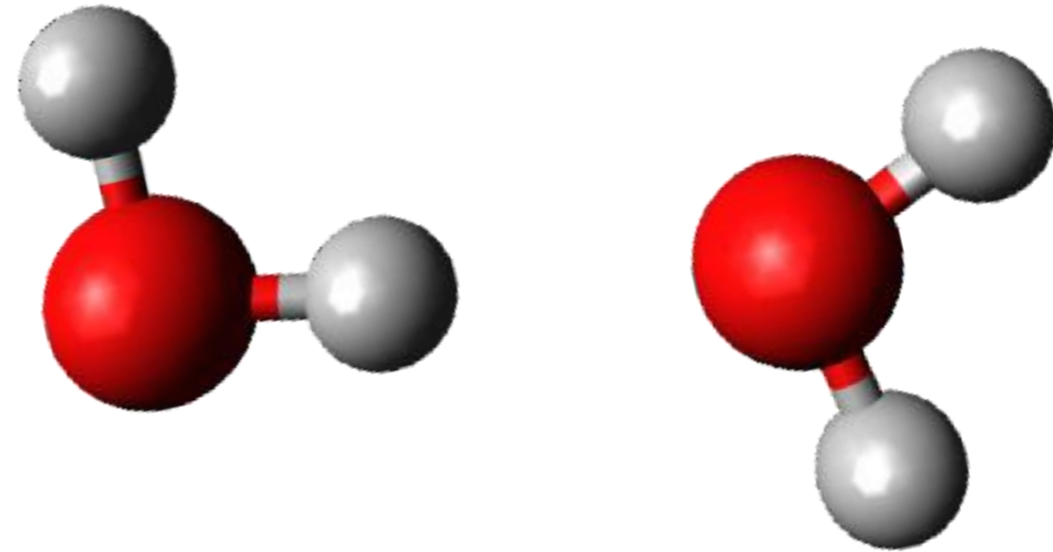
Newton's equation of motion



Molecular dynamics — moving atoms

$$\mathbf{F}_I = M_I \ddot{\mathbf{R}}_I$$

at time t



velocities: $\mathbf{v}_I(t) = \dot{\mathbf{R}}_I(t)$

positions: $\mathbf{R}_I(t)$

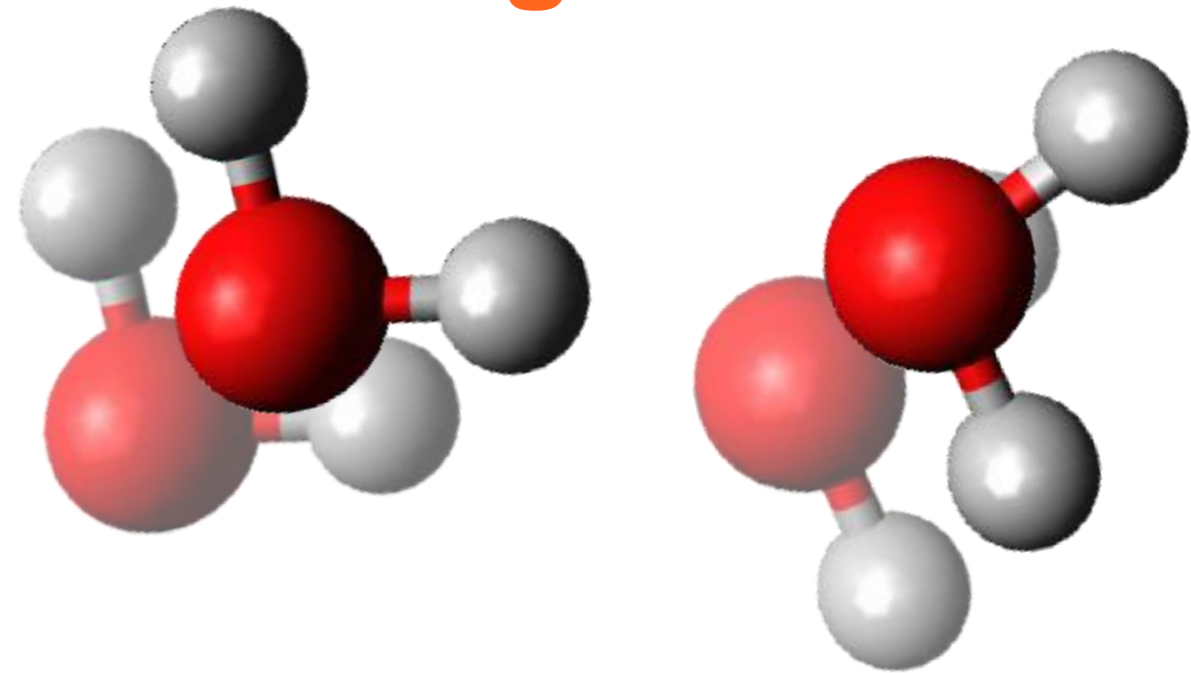
Molecular dynamics — moving atoms

$$\mathbf{F}_I = M_I \ddot{\mathbf{R}}_I$$

at time $t + \Delta t$

velocities: $\mathbf{v}_I(t + \Delta t)$

positions: $\mathbf{R}_I(t + \Delta t)$



move atoms
according to forces

Numeric integration — naive propagation

We start from Taylor expansion:

$$\mathbf{R}(t + \Delta t) = \mathbf{R}(t) + \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2 + \frac{1}{6}\dddot{\mathbf{R}}\Delta t^3 + \dots$$

Usual approximation: truncation

$$\mathbf{R}(t + \Delta t) = \mathbf{R}(t) + \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2$$

velocity v

acceleration a



Numeric integration — naive propagation

- Not time reversible!
- Does not conserve volume in phase space!
- Suffers from energy drift!

Usual approximation: truncation

$$\mathbf{R}(t + \Delta t) = \mathbf{R}(t) + \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2$$

velocity v

acceleration a



Numeric integration — Verlet algorithm

Taylor expansion for $+\Delta t$ and $-\Delta t$

$$\mathbf{R}(t + \Delta t) = \mathbf{R}(t) + \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2 + \frac{1}{6}\dddot{\mathbf{R}}\Delta t^3 + \dots$$
$$\mathbf{R}(t - \Delta t) = \mathbf{R}(t) - \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2 - \frac{1}{6}\dddot{\mathbf{R}}\Delta t^3 + \dots$$

The sum gives:

$$\mathbf{R}(t + \Delta t) + \mathbf{R}(t - \Delta t) = 2\mathbf{R}(t) + \ddot{\mathbf{R}}\Delta t^2 + \mathcal{O}(\Delta t^4)$$

Numeric integration — Verlet algorithm

Taylor expansion for $+\Delta t$ and $-\Delta t$

$$\mathbf{R}(t + \Delta t) = \mathbf{R}(t) + \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2 + \frac{1}{6}\dddot{\mathbf{R}}\Delta t^3 + \dots$$
$$\mathbf{R}(t - \Delta t) = \mathbf{R}(t) - \dot{\mathbf{R}}\Delta t + \frac{1}{2}\ddot{\mathbf{R}}\Delta t^2 - \frac{1}{6}\dddot{\mathbf{R}}\Delta t^3 + \dots$$

The sum gives, which is:

$$\mathbf{R}(t + \Delta t) = 2\mathbf{R}(t) - \mathbf{R}(t - \Delta t) + \ddot{\mathbf{R}}\Delta t^2$$



Numeric integration — Verlet algorithm

- Time reversible!
- Conserve volume in phase space!
- No energy drift!

The sum gives, which is:

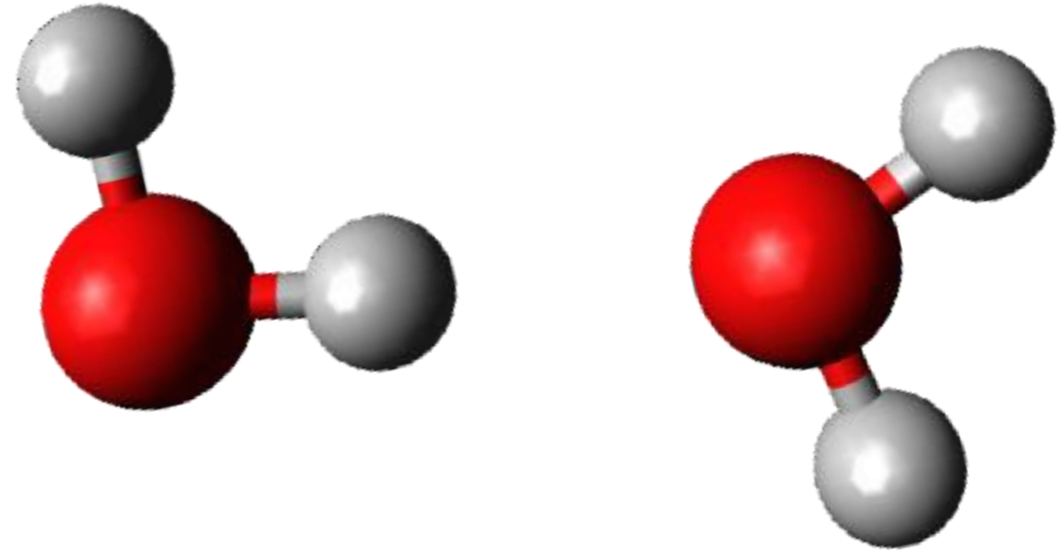
$$\mathbf{R}(t + \Delta t) = 2\mathbf{R}(t) - \mathbf{R}(t - \Delta t) + \ddot{\mathbf{R}}\Delta t^2$$

Molecular dynamics in practice

Initial $v(0)$, $R(0)$

Converge DFT
calculation

Calculate forces



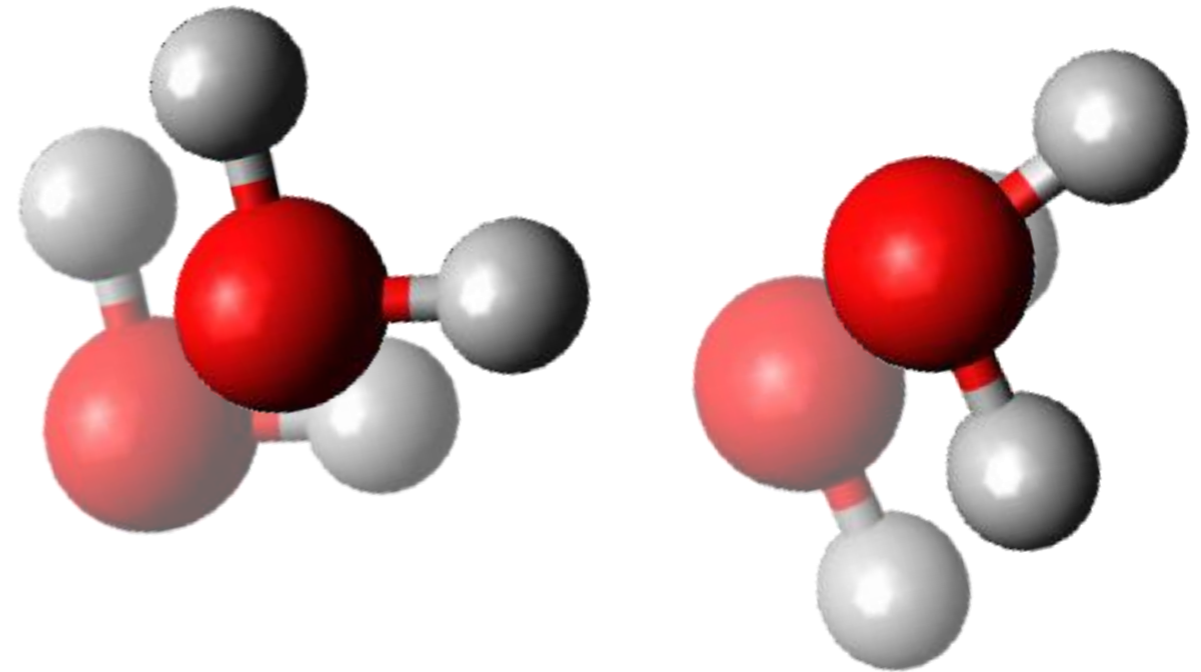
Molecular dynamics in practice

Initial $v(0), R(0)$

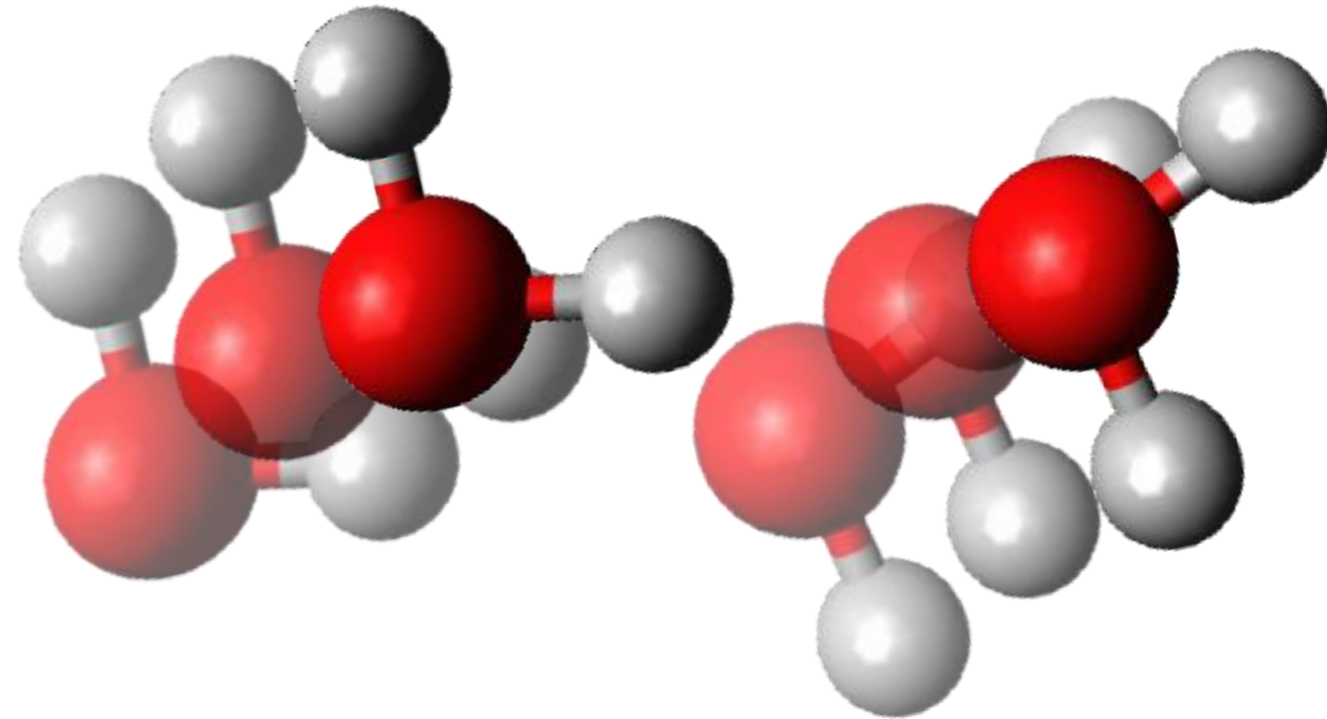
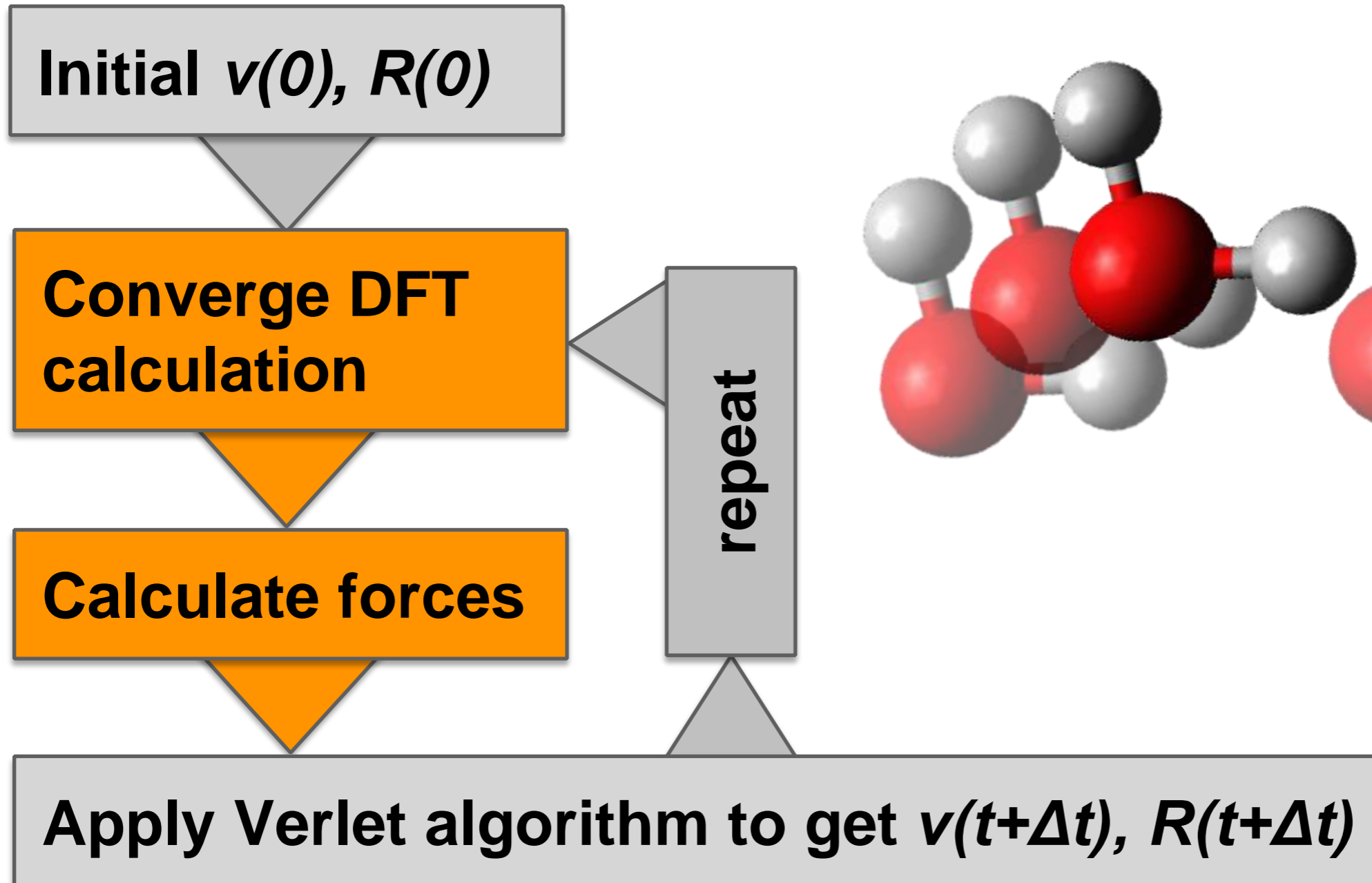
Converge DFT
calculation

Calculate forces

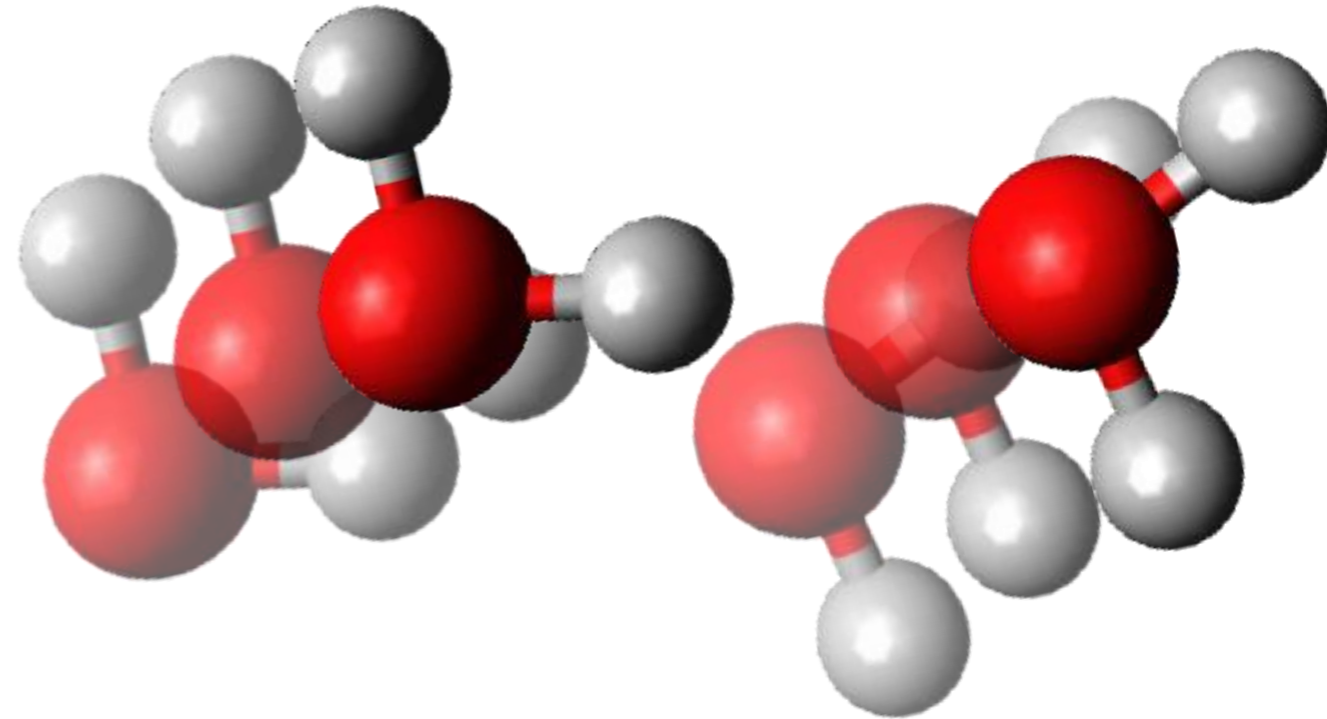
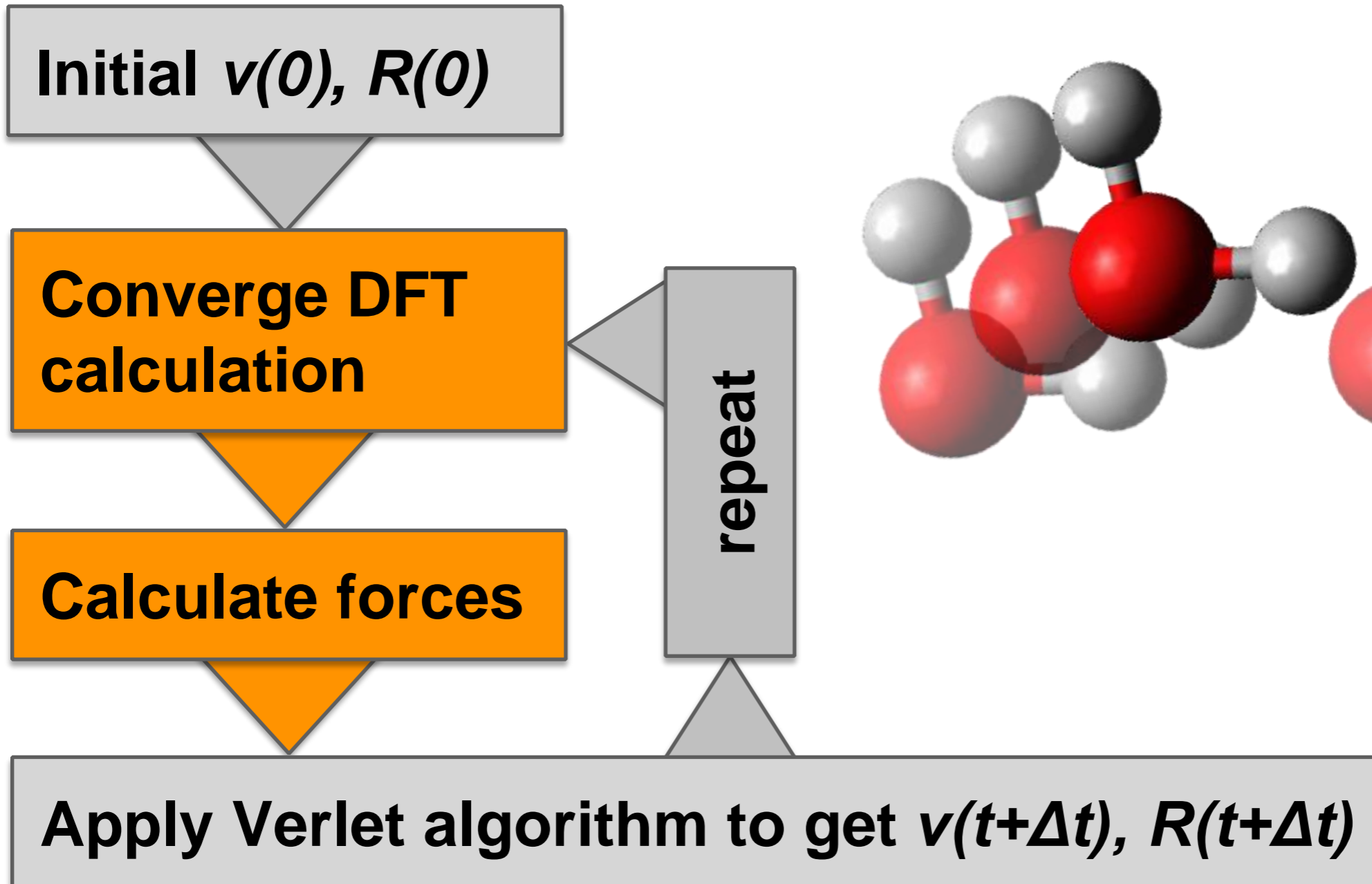
Apply Verlet algorithm to get $v(t+\Delta t), R(t+\Delta t)$



Molecular dynamics in practice



Questions – DFT MD in practice?



Thermodynamic ensembles

	conserved quantities
Micro-canonical ensemble (NVE)	# particles, volume, energy
Canonical ensemble (NVT)	# particles, volume, temperature
NPT or NPH ensemble (for phase transitions)	# particles, pressure, temperature or enthalpy
Grand-canonical ensemble (μVT) (for ad- or desorption)	chemical potential, volume, temperature

FHI-aims MD settings

FHI-aims technical detail:

settings for molecular dynamics calculations:

MD_run time ensemble : specify total simulation time and the MD ensemble

MD_time_step deltat : set the MD time step

sc_accuracy_forces : request an appropriate force accuracy

Thermodynamic ensembles

	conserved quantities
Micro-canonical ensemble (NVE)	# particles, volume, energy
Canonical ensemble (NVT)	# particles, volume, temperature
NPT or NPH ensemble (for phase transitions)	# particles, pressure, temperature or enthalpy
Grand-canonical ensemble (μ VT) (for ad- or desorption)	chemical potential, volume, temperature

typical MD ensemble (follows directly from Newton's equation of motion)

Thermodynamic ensembles

conserved quantities

Micro-canonical ensemble
(NVE)

particles, volume, energy

**Canonical ensemble
(NVT)**

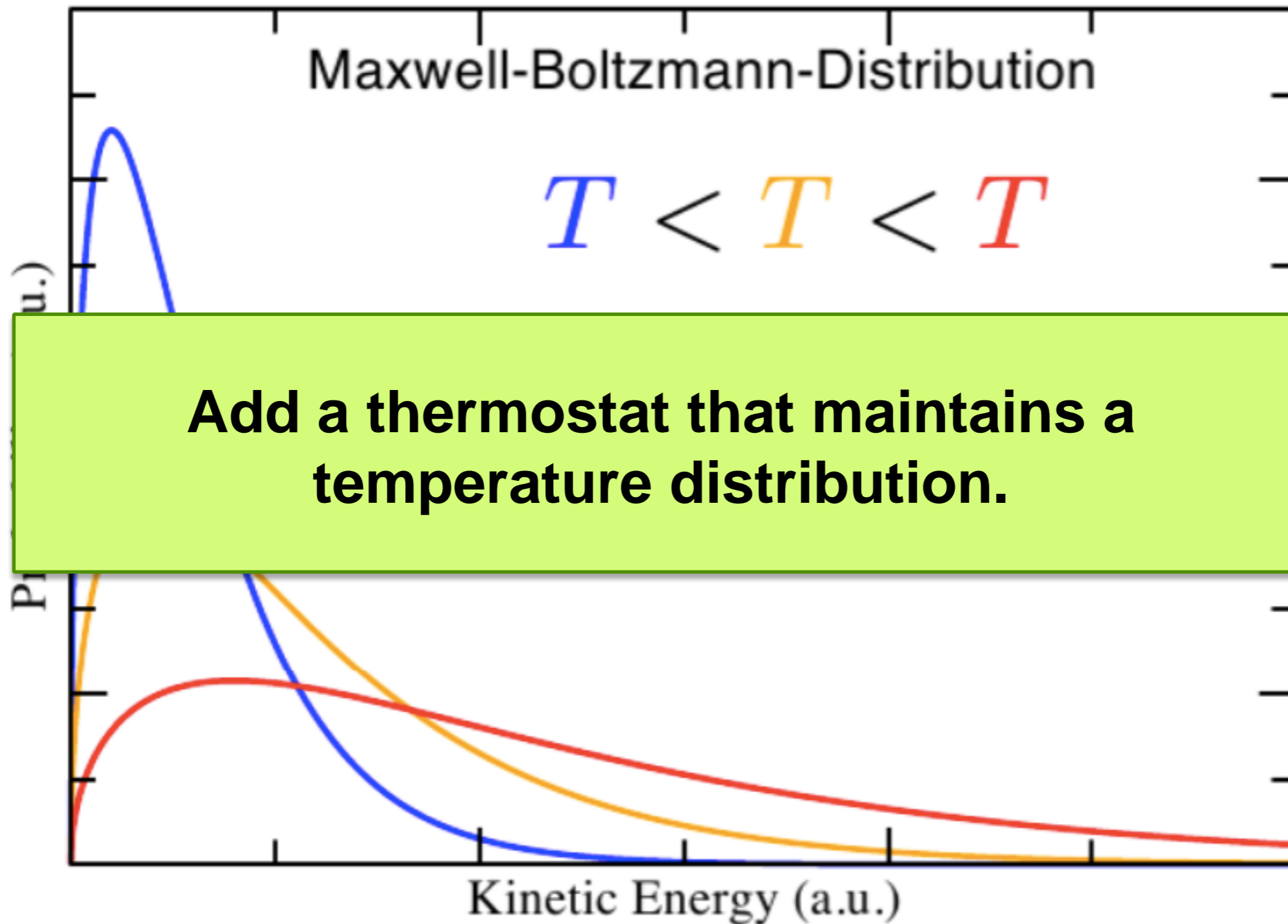
**# particles, volume,
temperature**

What about this ensemble?

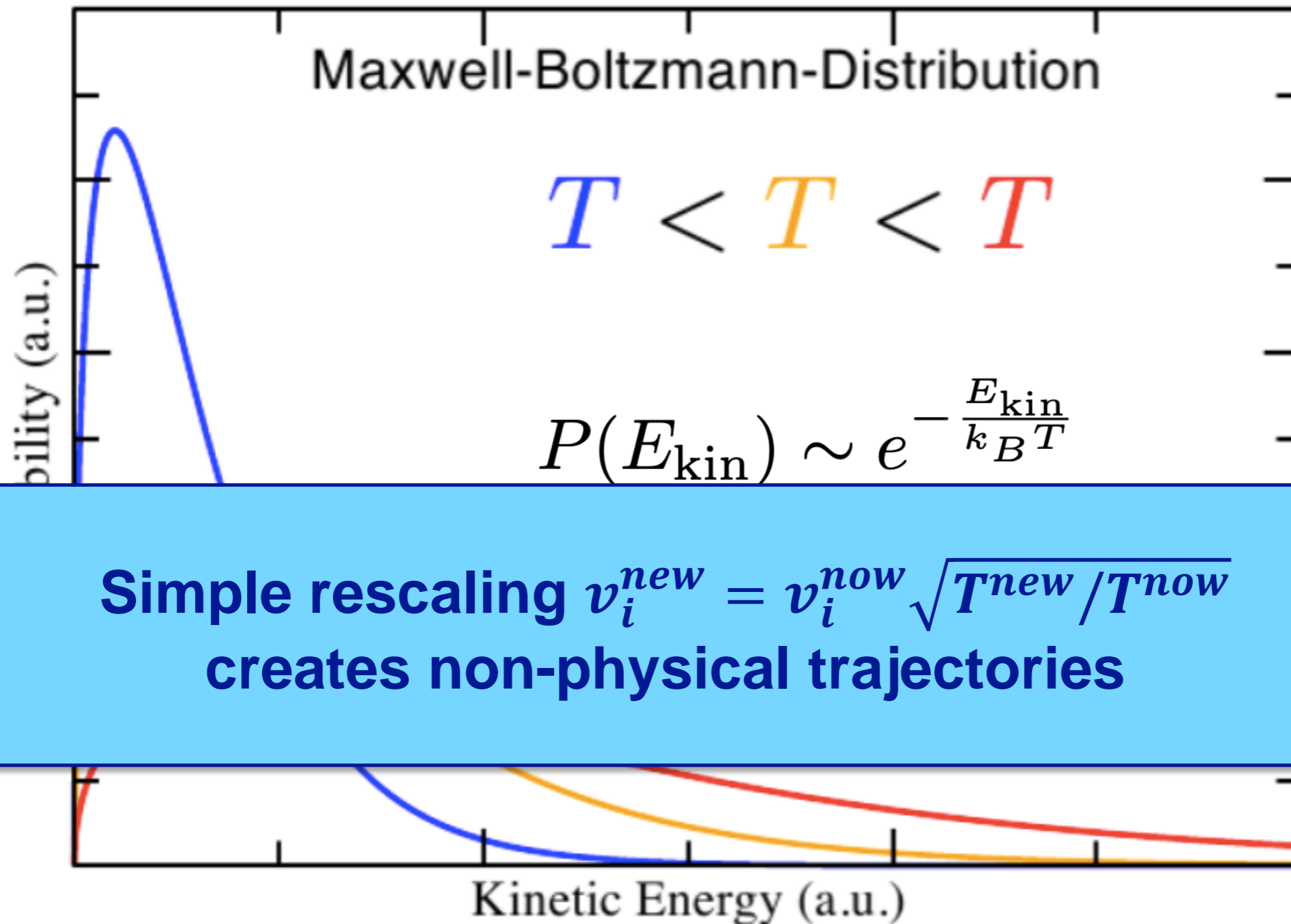
Need to maintain a constant temperature.



Temperature distributions



Temperature distributions



Nosé-Hoover thermostat

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{(\mathbf{p}'_I)^2}{2M_I} + V(\{\mathbf{R}\}) + \frac{p_\eta^2}{2Q} + 3Nk_B T \eta$$

original Hamiltonian

heat bath (in form of
fictitious oscillator)

One of possible thermostats giving appropriate trajectories (close to the wanted T)

Nosé-Hoover thermostat

fictitious mass

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{(\mathbf{p}'_I)^2}{2M_I} + V(\{\mathbf{R}\}) + \frac{p_\eta^2}{2Q} + 3Nk_B T \eta$$

original Hamiltonian

heat bath (in form of
fictitious oscillator)

$$\mathbf{p}'_I = \frac{\mathbf{p}_I}{s} \quad \text{scaled momenta}$$

Nosé-Hoover thermostat

fictitious mass

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{(\mathbf{p}'_I)^2}{2M_I} + V(\{\mathbf{R}\}) + \frac{p_\eta^2}{2Q} + 3Nk_B T \eta$$

original Hamiltonian

heat bath (in form of fictitious oscillator)

$$\dot{\mathbf{p}}'_I = \mathbf{F}_I - \frac{p_\eta}{Q} \mathbf{p}'_I$$

momenta are damped by fictitious mass

This is what maintains a constant temperature.

Questions?

fictitious mass

$$\mathcal{H} = \sum_{I=1}^{N_n} \frac{(\mathbf{p}'_I)^2}{2M_I} + V(\{\mathbf{R}\}) + \frac{p_\eta^2}{2Q} + 3Nk_B T \eta$$

original Hamiltonian

heat bath (in form of fictitious oscillator)

$$\dot{\mathbf{p}}'_I = \mathbf{F}_I - \frac{p_\eta}{Q} \mathbf{p}'_I$$

momenta are damped by fictitious mass

This is what maintains a constant temperature.

Computational Project

Preassignment for next Monday:

**Read *Planning a Computational Project* chapter
(on MyCourses->Lectures>Lecture 5)**

Questions?

Enjoy our “DFT MD & dispersion interactions” tutorial:

- Download from Mycourses
- `/work/courses/unix/PHYS/E0546/TUTORIALS`
- Ask for printouts

Interesting links related to the tutorial: More information on

Thermostats: https://nanohub.org/resources/7576/download/Martini_L4_TemperatureControl.pdf

https://www2.mpip-mainz.mpg.de/~andrienk/journal_club/thermostats.pdf

<http://www.courses.physics.helsinki.fi/fys/moldyn/lectures/L5.pdf>

And also FHI-aims manual.

Don't forget to put your name in the list of attendance!

