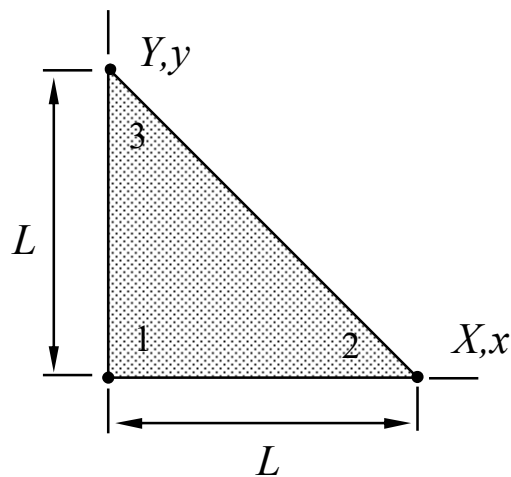


**LECTURE ASSIGNMENT 2.** In stationary thermo-elasticity, the variational densities of the thin slab mode of the plate model

$$\delta p_{\Omega}^{\text{int}} = - \begin{Bmatrix} \partial \delta \mathcal{G} / \partial x \\ \partial \delta \mathcal{G} / \partial y \end{Bmatrix}^T kt \begin{Bmatrix} \partial \mathcal{G} / \partial x \\ \partial \mathcal{G} / \partial y \end{Bmatrix}, \quad \delta p_{\Omega}^{\text{ext}} = \delta \mathcal{G}_s$$

represent the energy balance. Assuming that the thermal conductivity  $k$  and thickness  $t$  of the element are constants, derive the element contribution  $\delta P^{\text{int}}$  for the linear triangle element shown. Temperature at nodes 1 and 3 is known to be  $\mathcal{G}^{\circ}$  and the unknown nodal temperature is  $\mathcal{G}_2$ .



Name \_\_\_\_\_ Student number \_\_\_\_\_

- The linear shape functions can be deduced from the figure

$$N_2 = \frac{x}{L}, \quad N_3 = \frac{y}{L}, \quad N_1 = 1 - N_2 - N_3 = 1 - \frac{x}{L} - \frac{y}{L}$$

- Approximation to  $\mathcal{G}(x, y)$  and its variation  $\delta\mathcal{G}(x, y)$  (notice that the variation of a given quantity vanishes)

$$\mathcal{G} = \left(1 - \frac{x}{L}\right)\mathcal{G}^\circ + \frac{x}{L}\mathcal{G}_2, \quad \frac{\partial\mathcal{G}}{\partial x} = \frac{1}{L}(\mathcal{G}_2 - \mathcal{G}^\circ), \quad \frac{\partial\mathcal{G}}{\partial y} = 0$$

$$\delta\mathcal{G} = \frac{x}{L}\delta\mathcal{G}_2, \quad \frac{\partial\delta\mathcal{G}}{\partial x} = \frac{\delta\mathcal{G}_2}{L}, \quad \frac{\partial\delta\mathcal{G}}{\partial y} = 0$$

- When the approximation is substituted there, the variational density simplifies to

$$\delta p_\Omega^{\text{int}} = -\frac{\delta\mathcal{G}_2}{L}kt\frac{1}{L}(\mathcal{G}_2 - \mathcal{G}^\circ).$$

- Integration over the element gives

$$\delta P^{\text{int}} = \int_\Omega \delta p_\Omega^{\text{int}} dA = -\delta\mathcal{G}_2 \frac{kt}{2}(\mathcal{G}_2 - \mathcal{G}^\circ). \quad \leftarrow$$